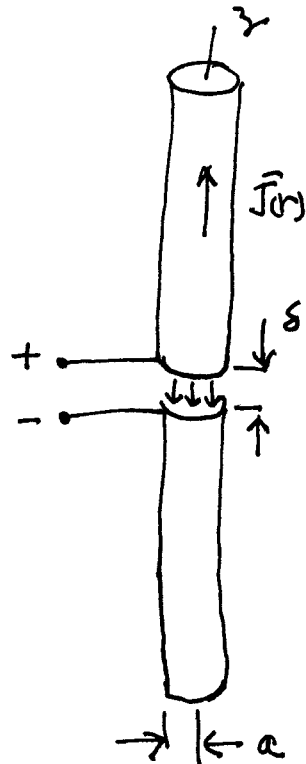


Last time

$$\underbrace{(\nabla_r \nabla_r \cdot + k^2) \int_{\text{vol}} \frac{\bar{J}(r') e^{-jkR}}{4\pi R} dr' = j\omega\epsilon_0 \bar{E}}_{\text{Pocklington's Eq.}}$$

Pocklington's Eq.

↓

$$\int_{-l/2}^{l/2} \frac{I(z') e^{-jkR}}{4\pi R} dz' = C \cos kz - j \frac{\omega\epsilon_0}{2k} \sin k|z|$$

$$R^2 = (z-z')^2 + a^2 \quad *$$

Hallen's Equation

* current on z-axis,
observation point on dipole, or v.v.

For solution to Hallen's Eq. we write

$$\int_0^{l/2} G(z, z') I(z') dz' = C \cos kz - j \frac{1}{2\eta} \sin k|z|$$

$$\text{where } G(z, z') = \frac{e^{-jkr}}{4\pi r} + \frac{e^{-jkr'}}{4\pi r'}$$

$$r^2 = a^2 + (z - z')^2 ; r'^2 = a^2 + (z + z')^2$$

Which we propose to solve approximately by
method of moments.

den general

$$\int_a^b G(z, z') f(z') dz' = g(z) \quad \text{F. H., first kind}$$

$$f(z) \approx \hat{f}(z) = c_1 f_1(z) + c_2 f_2(z) + \dots + c_N f_N(z)$$

$$\sum_{n=1}^N c_n \int_a^b \int_a^b G(z, z') f_n(z') w_m(z) dz' dz = \int_a^b g(z) w_m(z) dz$$

$\sum_{n=1}^N c_n a_{mn} = b_m$ /...

Rewriting $[a_{mn}][c_n] = [b_m]$,

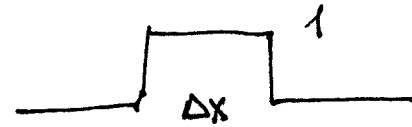
N equations in N unknowns

Galerkin's Method $w_m = f_m$

Point-Matching Method $w_m = \delta(z - z_m)$

Three Popular Basis Functions

Pulse



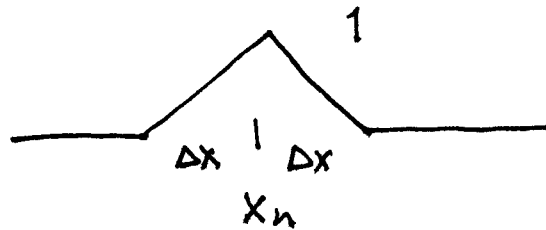
$$\begin{aligned} p_n(x) &= 1 & x \in \Delta x \\ &= 0 & x \notin \Delta x \end{aligned}$$

Is the simplest and most common.

But there are alternatives that are more effective.

/...

Triangle Basis Functions

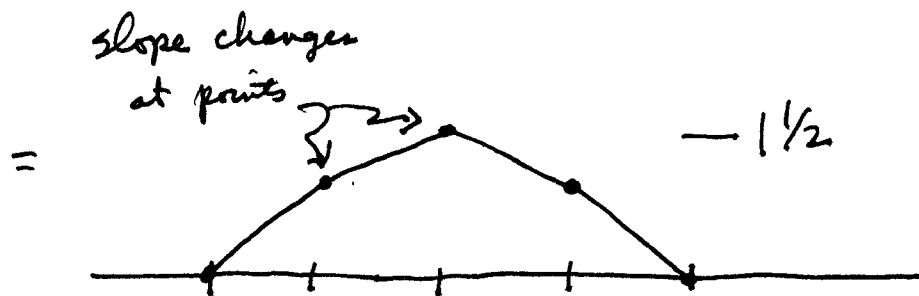
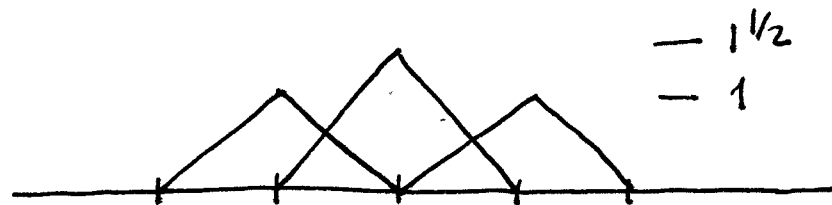


$$\begin{aligned} \hat{x}_n(x) &= \frac{x - x_n + \Delta x}{\Delta x} & x_n - \Delta x < x < x_n \\ &= -\frac{(x - x_n - \Delta x)}{\Delta x} & x_n < x < x_n + \Delta x \\ &= 0 & \text{otherwise} \end{aligned}$$

...

Triangle basis functions provide a piecewise-linear approximation to the referenced function

E.g.



Similarly, a sinusoidal basis function is illustrated in S&T, Section 7.5 (pp 323-ff).

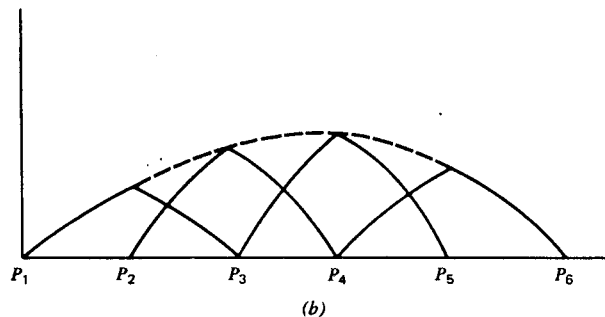
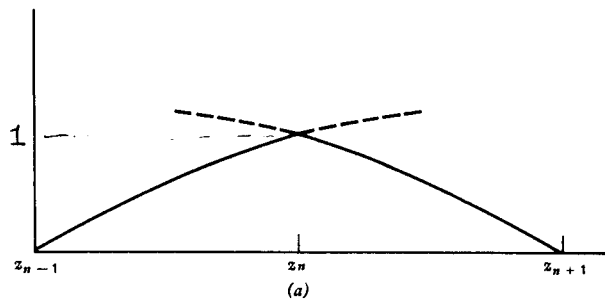


Figure 7-11 (a) Piecewise sinusoidal expansion function. (b) Set of overlapping piecewise sinusoidal expansion functions.

While the derivatives are still discontinuous at the "joints", the overall approximation to a general current distribution, e.g., is improved.

From S&T p. 324

Sample solutions for

$$\left(\frac{a}{L}\right)^{-1} = 1/4, 1/2, 3/4, 1$$

$$a/x = 0.0001, 0.01$$

shown for pulse function $f_n(z)$

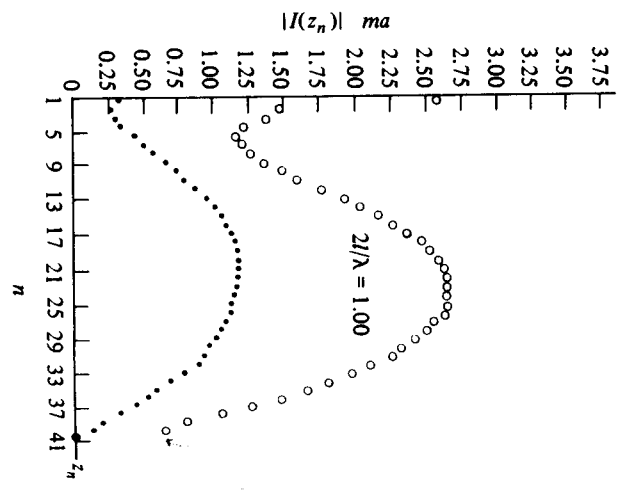
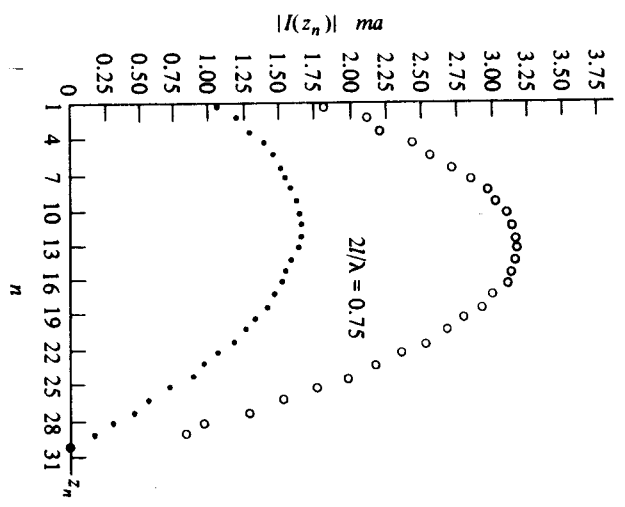
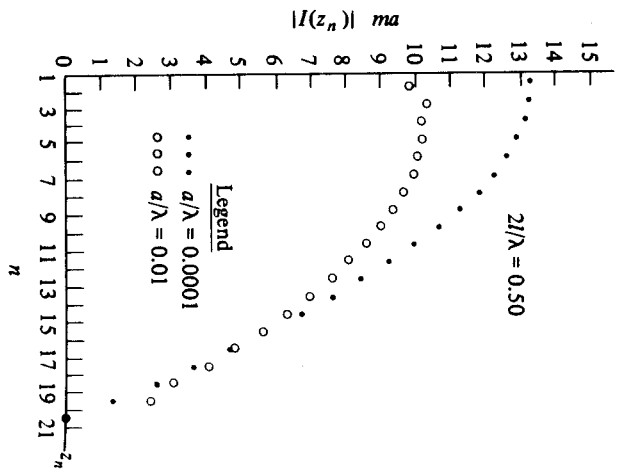
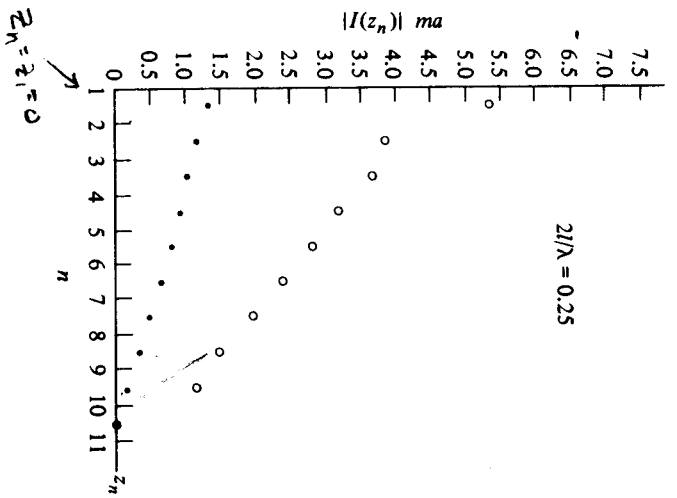


Fig. 7.6 The Magnitude of $I(z)$ for Center-Fed Cylindrical Dipoles of Various Lengths and Diameters; Pulse Function Solution

from Elliott

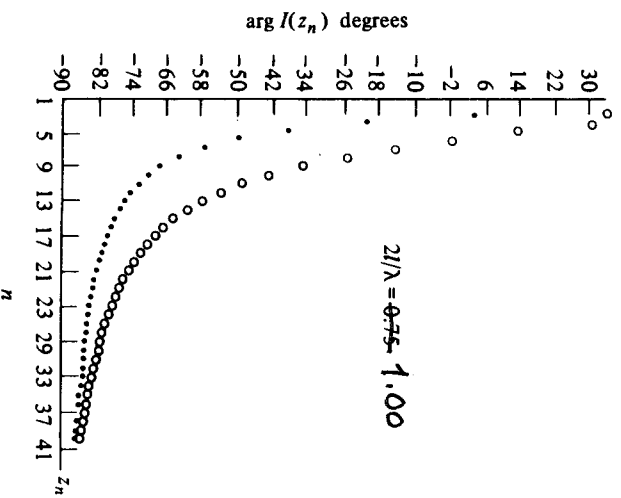
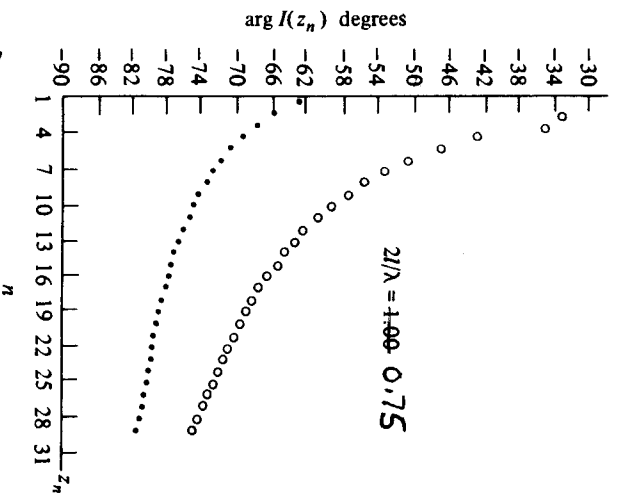
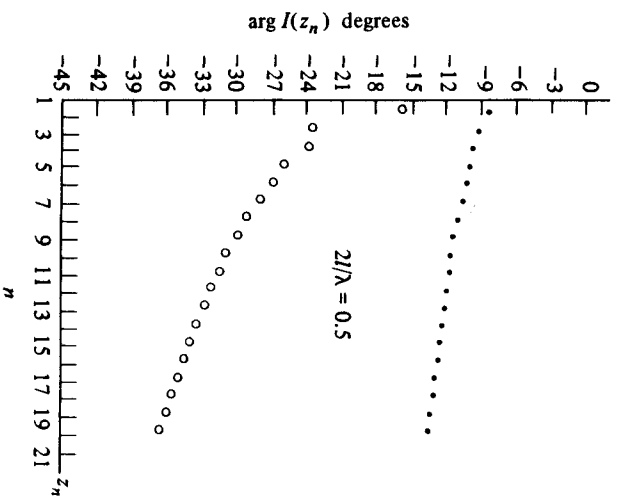
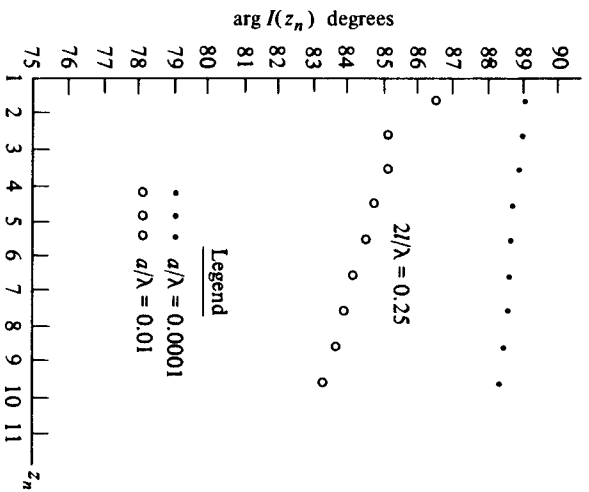


Fig. 7.7 The Phase of $I(z)$ for Center-Fed Cylindrical Dipoles of Various Lengths and Diameters; Pulse Function Solution

L256

Elvira

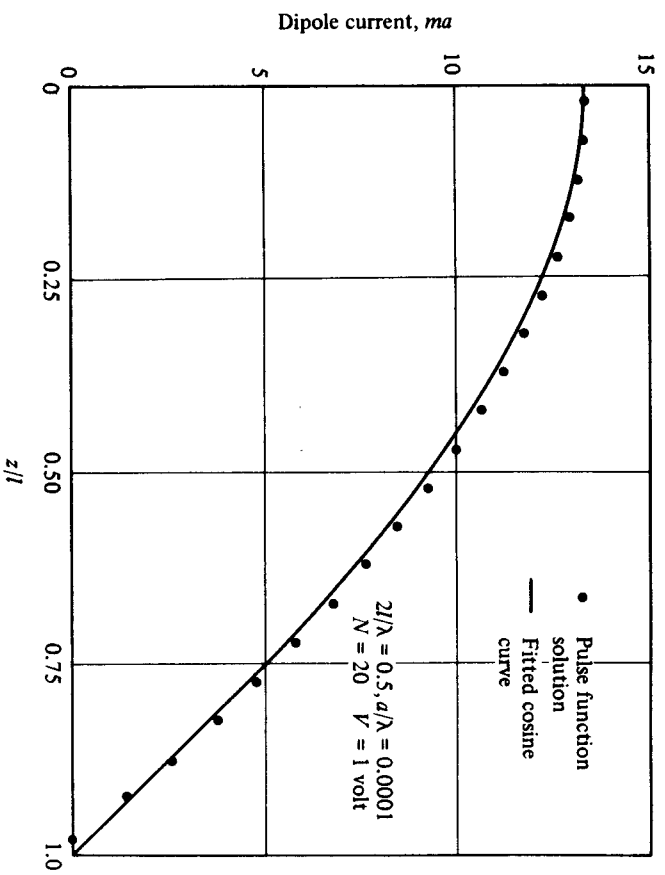


Fig. 7.8 The Current Distribution on a Half-Wavelength Long Center-Fed Cylindrical Dipole

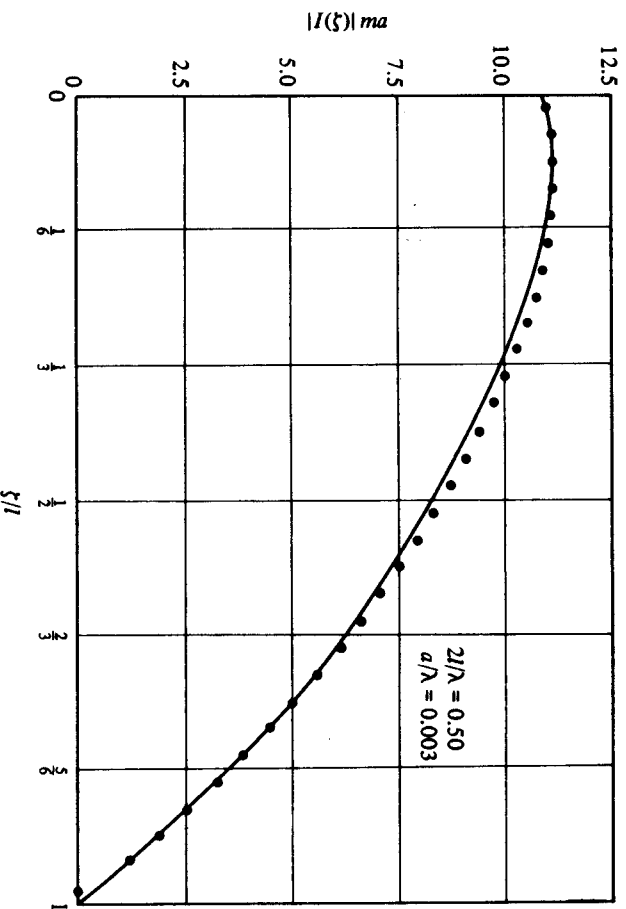


Fig. 7.9 Comparison of Theory and Experiment: Current Distribution on a Half-Wavelength Center-Fed Cylindrical Dipole (Solid Curve Experimental Results of T. Morita, *Proc. IRE*, vol. 38, pp. 898-904, © 1950 IEEE; Dots Computer Results Using Pulse Functions.)

Elliott

Obtained w/ sinusoidal basis functions

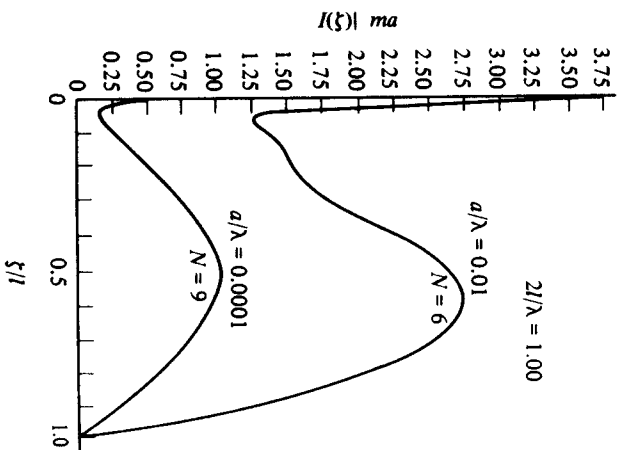
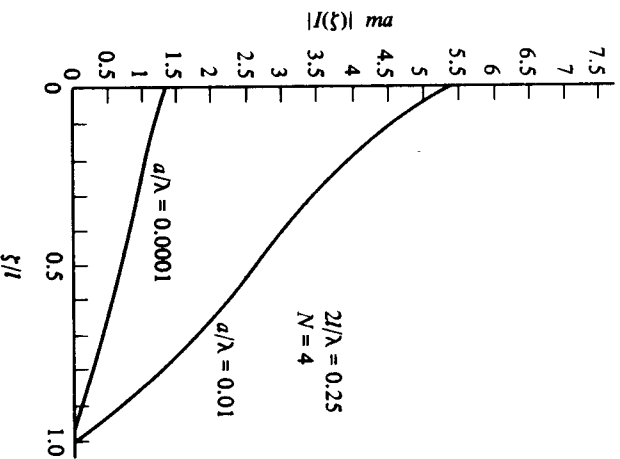
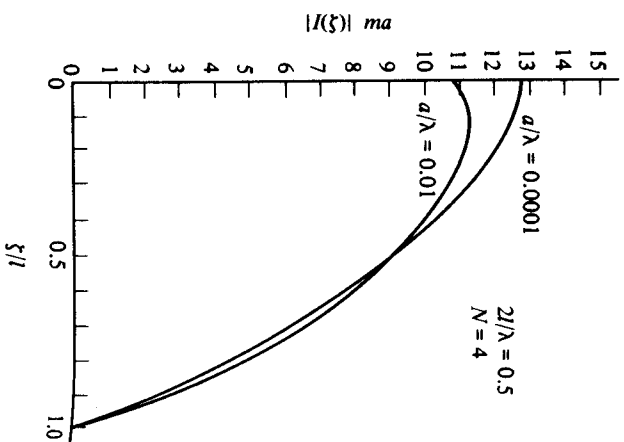
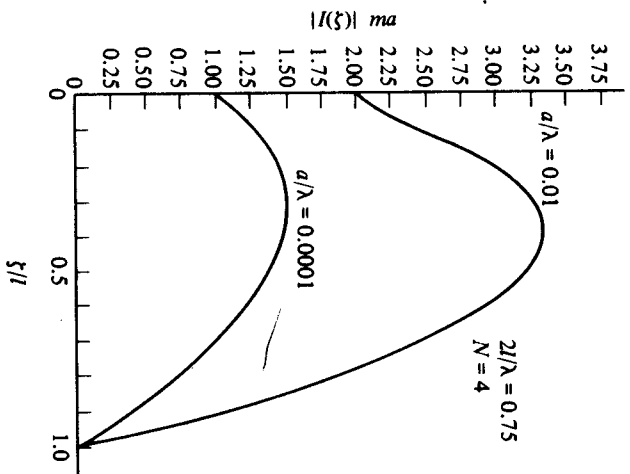


Fig. 7.10 The Magnitude of $I(\xi)$ for Center-Fed Cylindrical Dipoles of Various Lengths and Diameters: Solution Using Sinusoidal Basis Functions

Elleitt

TABLE 7.5 Comparison of the input impedance of a center-fed cylindrical dipole using different computational methods

Normalized Length $2l/\lambda$	Normalized Radius a/λ	Input Impedance in Ohms		
		Storer's Two-Term Approximation	First-Order Approximation to Hallén's Equation	King-Middleton Second-Order Approximation
0.25	0.01	$11.63 - j185$	$15.99 - j240$	$13.98 - j166$
0.25	0.0001	$12.93 - j723$	$14.67 - j756$	$12.90 - j811$
0.50	0.01	$101 + j32.82$	$87.34 + j35.68$	$92.51 + j38.30$
0.50	0.0001	$80.15 + j42.61$	$79.08 + j43.52$	$79.89 + j43.47$
0.75	0.01	$566 + j3.10$	$437 + j318$	$543 + j32.2$
0.75	0.0001	$521 + j1019$	$433 + j1018$	$540 + j1016$
1.00	0.01	$290 - j363$	$559 - j594$	$177 - j339$
1.00	0.0001	$2370 - j2129$	$3052 - j2626$	$2233 - j2150$

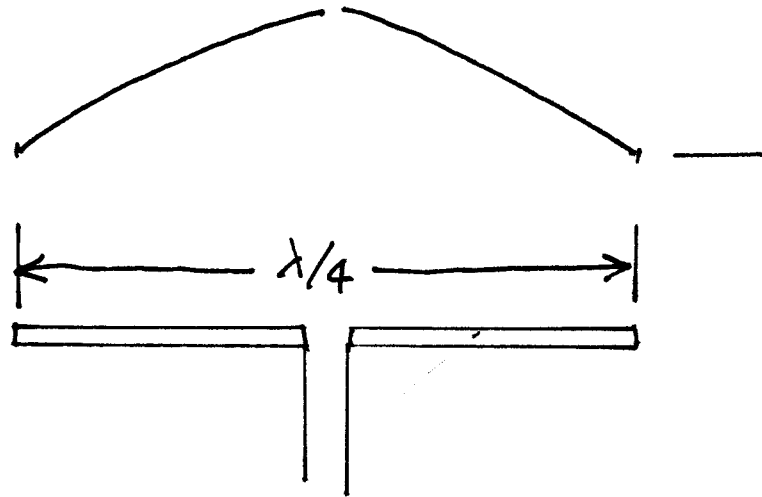
TABLE 7.1 $G_d(\theta)$ versus θ for idealized current distribution and for current distribution found using pulse functions

θ°	$G_d(\theta)$	
	$I(z) = I_m \text{sinc}(l - z)$	$I(z)$ Found by Using Pulse Functions
0	0.000	0.000 1.163°
6	0.082	0.081 1.141°
12	0.165	0.162 1.098°
18	0.249	0.244 1.029°
24	0.333	0.327 0.940°
30	0.418	0.411 0.833°
36	0.503	0.496 0.717°
42	0.587	0.580 0.596°
48	0.668	0.662 0.475°
54	0.746	0.740 0.361°
60	0.816	0.812 0.258°
66	0.879	0.876 0.169°
72	0.930	0.928 0.096°
78	0.968	0.967 0.043°
84	0.992	0.992 0.011°
90	1.000	1.000 0°

TABLE 7.2 Approximations to input impedance for various cylindrical dipoles using pulse functions or sinusoidal functions as basis functions

Normalized Length $2l/\lambda$	Normalized Radius a/λ	Number of Pulse Functions N	Self-Impedance in Ohms $(c)^{-1}$ Pulse Functions	Self-Impedance in Ohms Sinusoidal Basis Functions
0.25	0.01	10	* $11.3 - j1.86$	10.2 - j1.85
0.25	0.0001	10	12.9 - j7.37	12.5 - j7.39
0.50	0.01	20	97.3 + j27.8	90.2 + j22.2
0.50	0.0001	20	74.0 + j11.3	74.2 + j26.4
0.75	0.01	30	534 + j79.9	477 + j180
0.75	0.0001	30	424 + j82.7	403 + j88.2
1.00	0.01	40	178 - j344	40 - j255
1.00	0.0001	40	2724 - j1067	439 - j1445

Elvitt



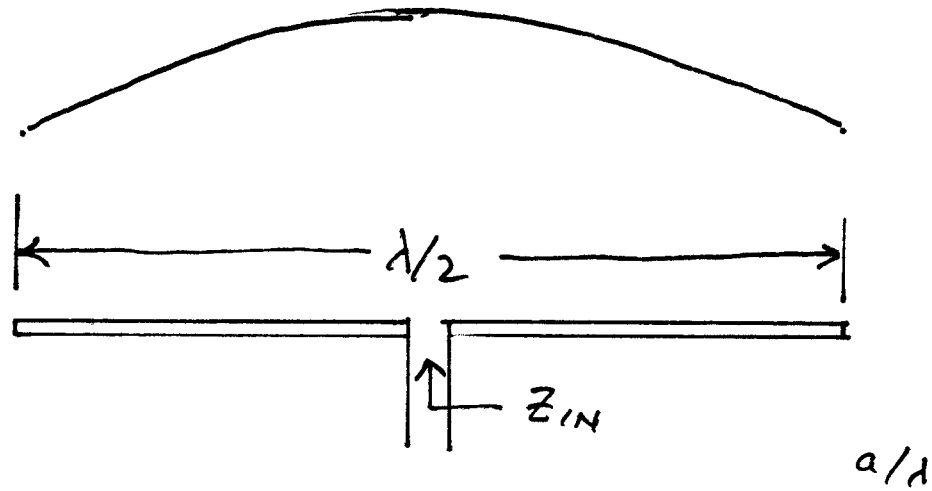
$$Z_{IN} = 666.7 \angle -89^\circ \approx 11.6 - j666.6$$

$$\frac{a/\lambda}{0.0001}$$

$$Z_{IN} = 181.3 \angle -87^\circ \approx 9.5 - j181$$

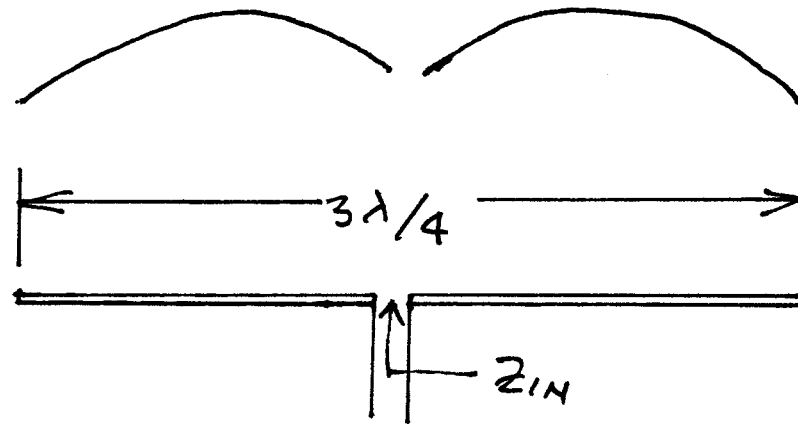
$$0.01$$

From Fig 7.6, 7.7.



$$Z_{IN} = 76.9 \angle +8^\circ = 76 + j10.7 \quad 0.0001$$

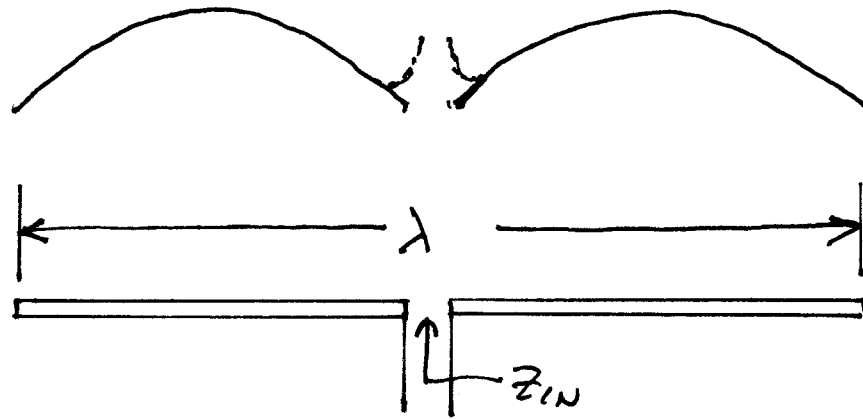
$$Z_{IN} = 105 \angle +23^\circ = 96 + j104 \quad 0.01$$



a/λ

$$Z_{IN} = 1000 \angle +62 = 464 + j880 \quad 0.0001$$

$$Z_{IN} = 571 \angle +32 = 484 + j302 \quad 0.01$$



$\frac{R}{\lambda}$

$$Z_{IN} = 4000 \angle -30^\circ = 3464 - j2000 \quad 0.0001$$

$$Z_{IN} = 666 \angle -34^\circ = 552 - j372 \quad 0.01$$

Table 5-1 Simple Formulas for the Input Resistance of Dipoles

Length L	Input resistance R_{in} (ohms)
$0 < L < \frac{\lambda}{4}$	$20\pi^2 \left(\frac{L}{\lambda}\right)^2$
$\frac{\lambda}{4} < L < \frac{\lambda}{2}$	$24.7 \left(\pi \frac{L}{\lambda}\right)^{2.4}$
$\frac{\lambda}{2} < L < 0.637\lambda$	$11.14 \left(\pi \frac{L}{\lambda}\right)^{4.17}$

Table 5-2 Wire Lengths Required To Produce a Resonant Half-Wave Dipole for a Wire Diameter of $2a$ and Length L

Length to diameter ratio, $L/2a$	Percent shortening required	Resonant length L	Dipole thickness class
5000	2	0.49λ	very thin
50	5	0.475λ	thin
10	9	0.455λ	thick

$L = 2l$

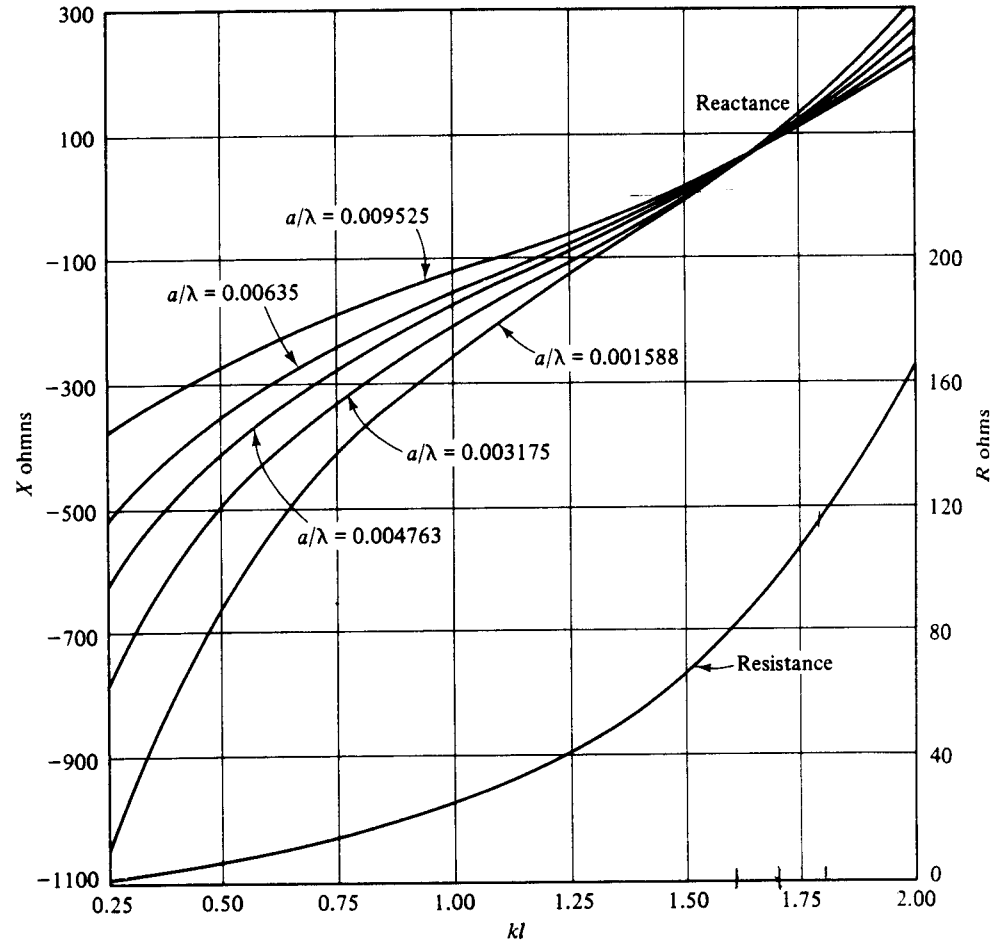


Fig. 7.11 The Resistance and Reactance of a Center-Fed Dipole versus kl and a/λ ; Values Computed by the Induced EMF Method Using Equation 7.65

Definitions: Figs & Tables, immediately following.

Figures
and Tables
beginning
with "7"
~~-----~~

$\frac{l}{2} \rightarrow l$, total length of Dipole is $2l$

$$S_2 = 2 \ln(2l/a) = 2 \ln \left[\frac{\text{length}}{\text{radius}} \right]$$

All cases center-fed.

$$kL = \frac{2\pi l}{\lambda}$$

$\frac{\lambda}{2}$ dipole	, $kL = 1.57$
λ dipole	, $kL = 3.14$

/...

<u>SZ</u>	Length / radius	radius / length
100	$e^{50} = 10^{21}$	≈ 0
20	$e^{10} = 22,000$	0.000045
10	$e^5 = 148$	0.0067
5	$e^{2.5} = 12.2$	0.08

Figs and Tables beginning with "5" are from S&T

+

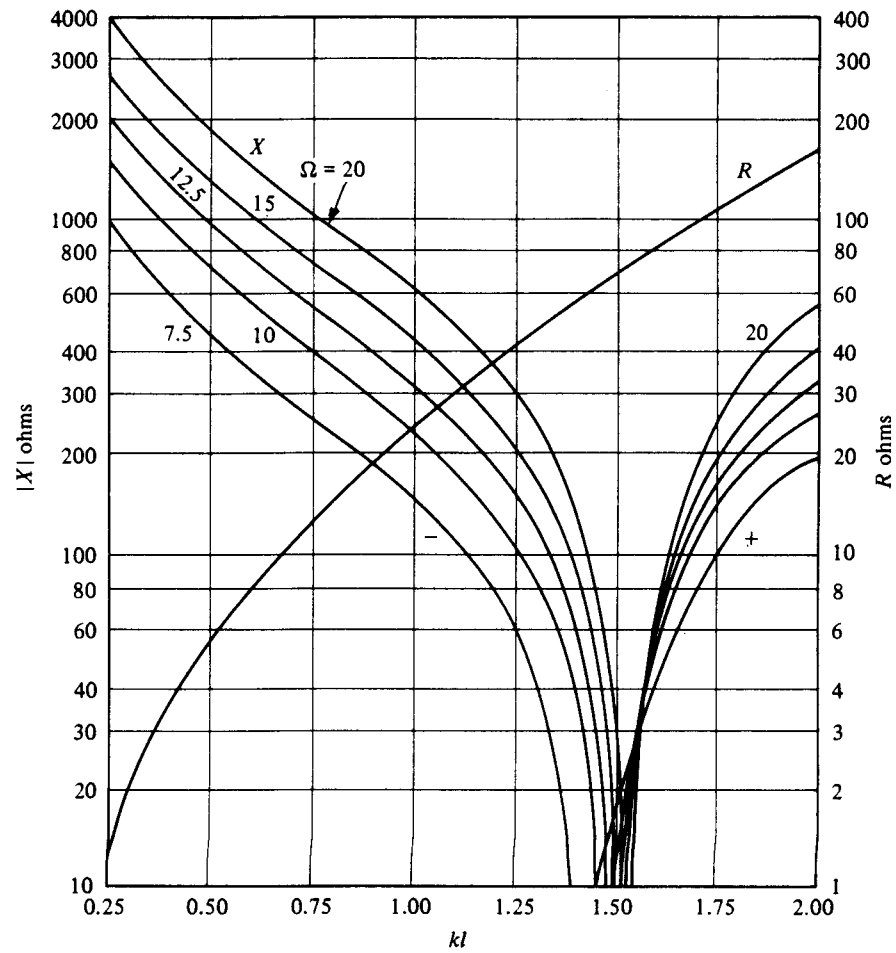


Fig. 7.12 The Resistance and Reactance of a Center-Fed Dipole versus kl and Ω ; Values Computed by the Induced EMF Method Using Equation 7.65

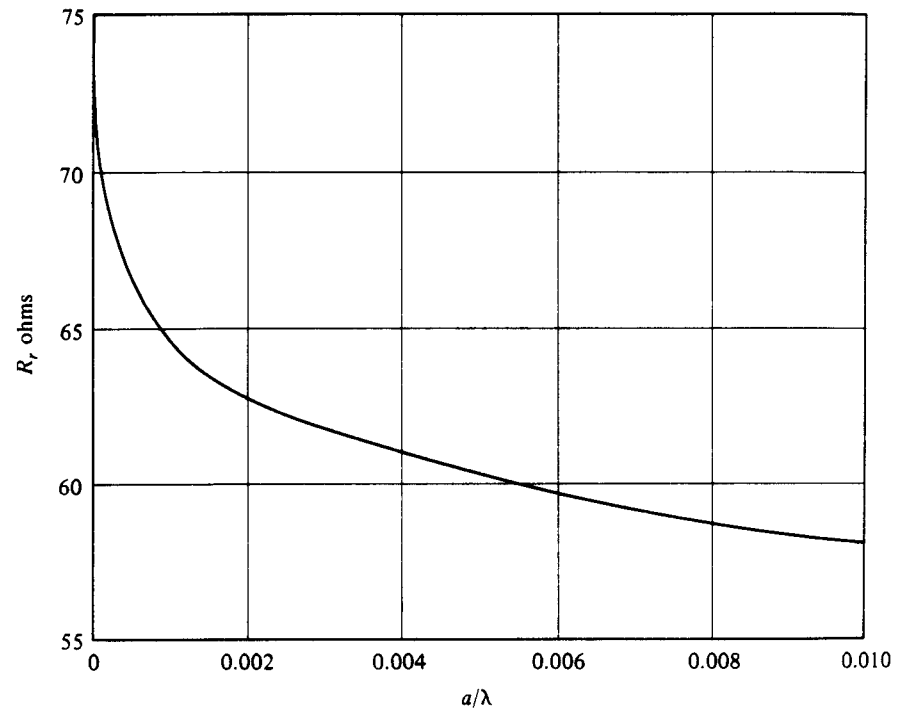


Fig. 7.14 Resonant Resistance versus Radius for Center-Fed Cylindrical Dipoles

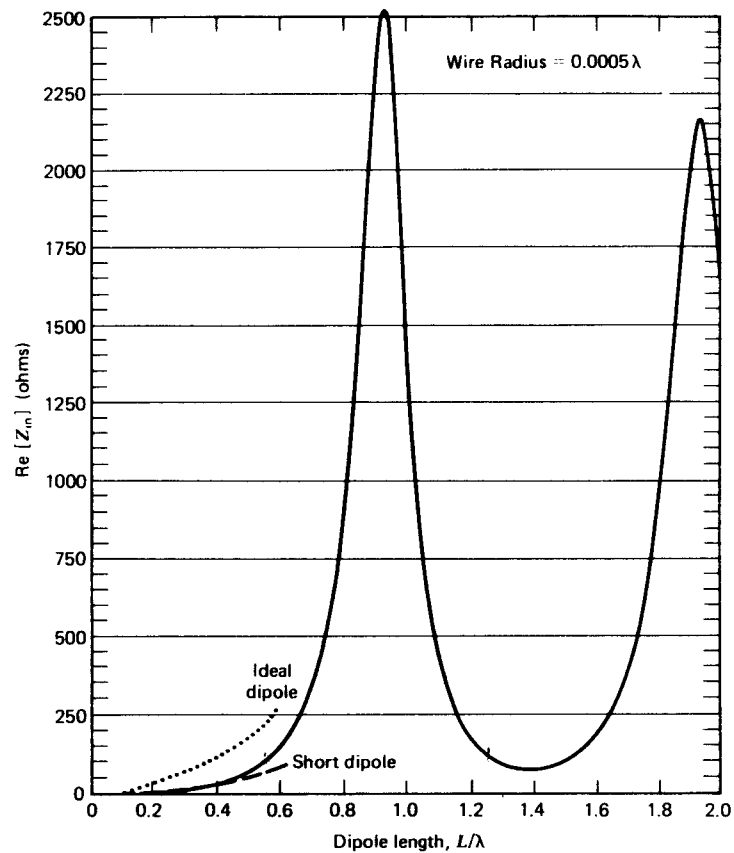


Figure 5-5 Calculated input resistance of a center-fed wire dipole of 0.0005λ radius as a function of length L (solid curve). Also shown is the input resistance $R_{in} = 80\pi^2(L/\lambda)^2$ of an ideal dipole with a uniform current distribution (dotted curve) and the input resistance $R_{in} = 20\pi^2(L/\lambda)^2$ of a short dipole with a triangular current distribution approximation (dashed curve).

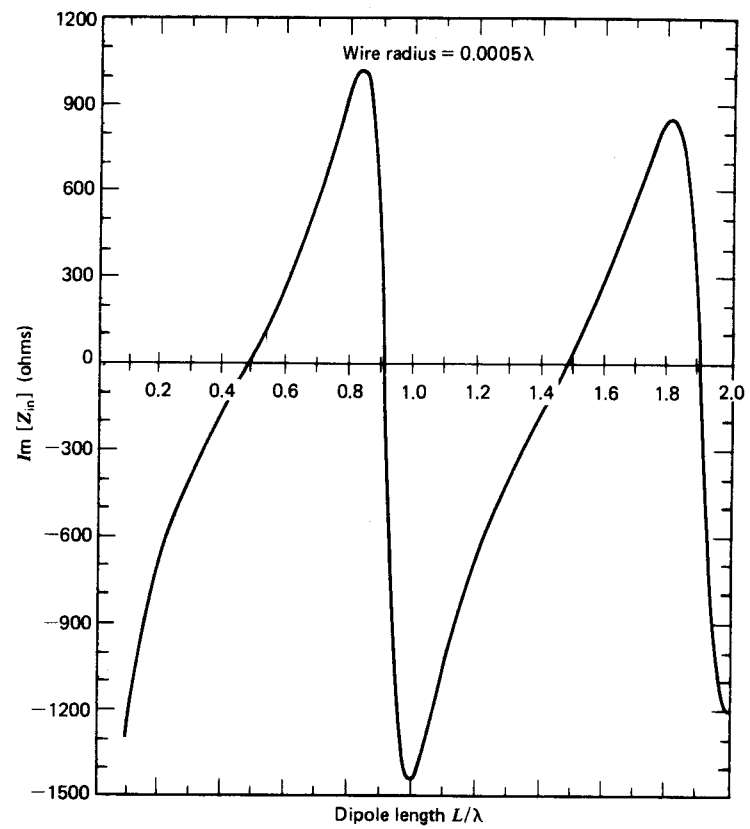


Figure 5-6 Calculated input reactance of center-fed wire dipole of radius 0.0005λ as a function of length L .

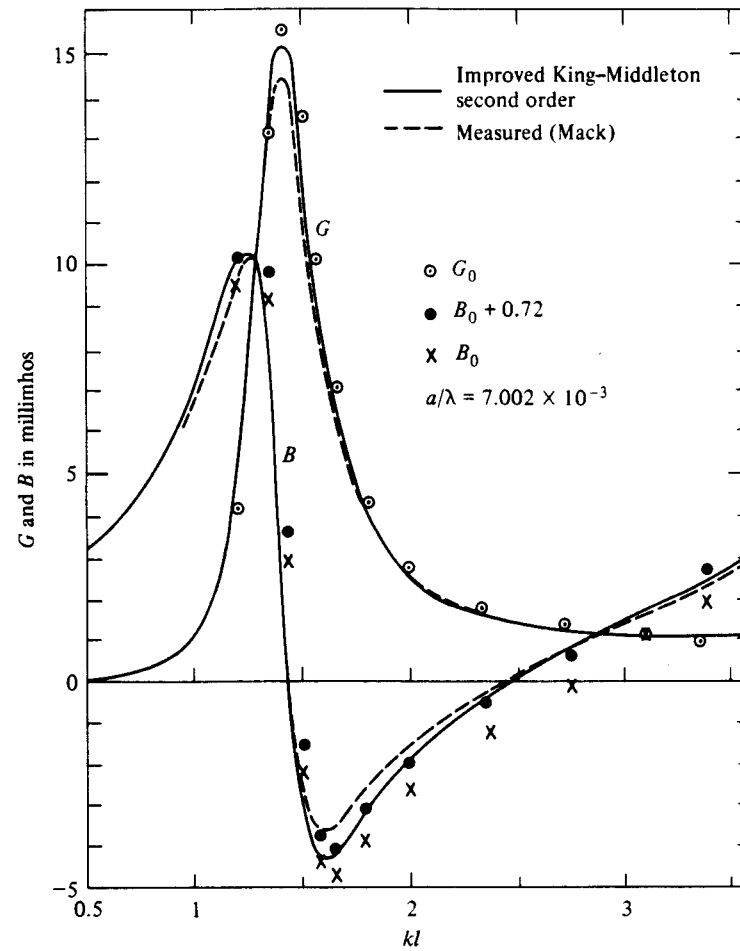


Fig. 7.17 A Comparison of the Improved King-Middleton Second-Order Admittance and the Measured Admittance of a Center-Fed Cylindrical Dipole (Measurements by Mack²⁸) (© 1967 IEEE. Reprinted from R. W. P. King, *IEEE Proceedings*, pp. 2-16, 1967.)

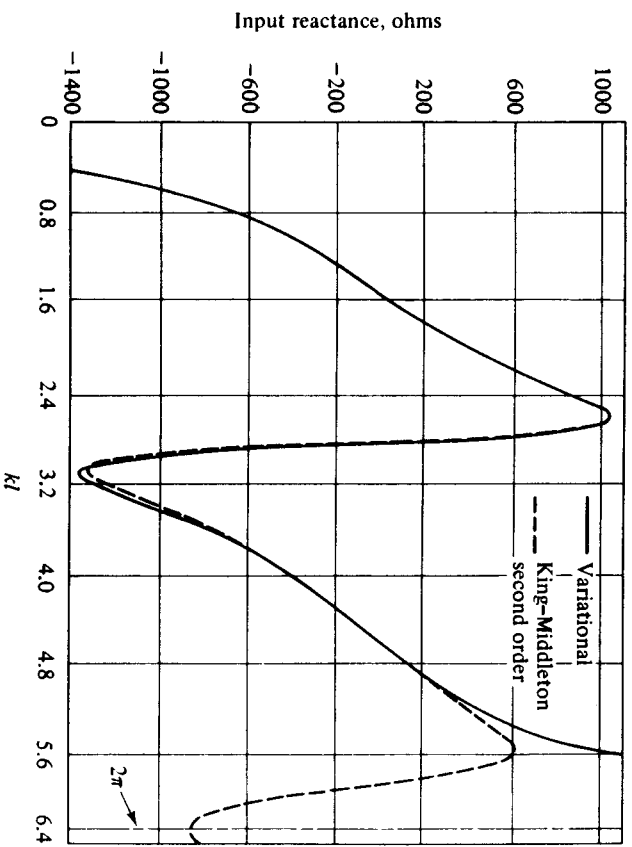
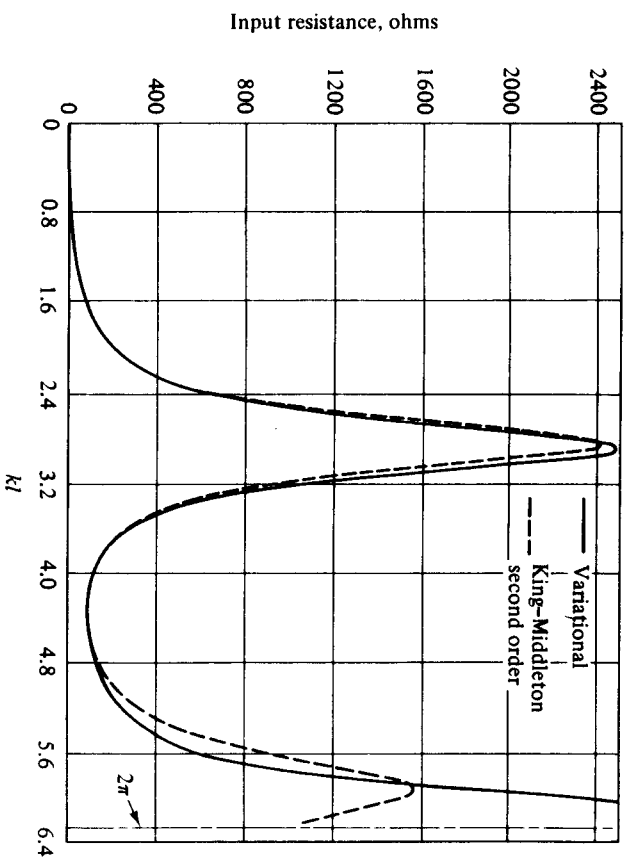


Fig. 7.16 A Comparison of Storer's Variational Solution and King-Middleton Second Order Values for the Input Impedance of a Center-Fed Cylindrical Dipole; $\Omega = 2l/\ln(2l/a) = 15$ (Reprinted from J. E. Storer, Cruft Laboratory Report No. 101, 1950, Courtesy of Harvard University.)