

# Chapter 2

# Mathematical Formulation of the Superposition Principle

**Superposition** → add states together, get new states.

**Math quantity associated with states must also have this property.**

**Vectors have this property.**

**In real three dimensional space, three basis vectors are sufficient to describe any point in space.**

**Can combine three vectors to get new bases vectors, which are also O.K. under appropriate combination rules.**

**In Q. M., may have many more than three states.**

**Need one vector for each state.**

**May be finite or infinite number of vectors;**

**depends on finite or infinite number of states of the system.**

Dirac  $\longrightarrow$  Call Q.M. vectors  
ket vectors or just  
kets.

Symbol  $\longrightarrow$   $| \rangle$  general ket

Particular ket, A  $\longrightarrow$   $|A\rangle$

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Kets can be multiplied by complex numbers and added.

$$|R\rangle = C_1 |A\rangle + C_2 |B\rangle$$

$C_1$  and  $C_2$  are **complex numbers**

Can add together any number of kets.  $|S\rangle = \sum_i C_i |L_i\rangle$  **Vector space**

If  $|x\rangle$  is continuous over some range of  $x$

$$|Q\rangle = \int C(x) |x\rangle dx$$

**Hilbert space**

## Dependent and Independent Kets

A ket that is expressible linearly in terms of other kets is **dependent** on them.

$$|R\rangle = C_1 |A\rangle + C_2 |B\rangle$$

$|R\rangle$  is dependent on  $|A\rangle$  and  $|B\rangle$

A set of kets is **independent** if no one of them is expressible linearly in terms of the others.

(In real space, three orthogonal basis vectors are independent.)

## Connection between theory and real world

**Assume:** Each state of a dynamical system at a particular time corresponds to a ket vector; the correspondence being such that if a state results from a superposition of other states, its corresponding ket vector is expressible linearly in terms of the corresponding ket vectors of the other states.

**Order of superposition doesn't matter.**

$$|R\rangle = C_1|A\rangle + C_2|B\rangle = C_2|B\rangle + C_1|A\rangle$$

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If  $|R\rangle = C_1|A\rangle + C_2|B\rangle$ ,

then,  $|A\rangle$  can be expressed in terms of  $|R\rangle$  and  $|B\rangle$ , etc.

$$|A\rangle = b_1|B\rangle + b_2|R\rangle$$

**A state is dependent on other states if its ket is dependent on the corresponding kets of the other states.**

**The superposition of a state with itself is the same state.**

$$C_1|A\rangle + C_2|A\rangle = (C_1 + C_2)|A\rangle \Rightarrow |A\rangle$$

unless  $(C_1 + C_2) = 0$  then no state.

**A ket corresponding to a state multiplied by any complex number (other than 0) gives a ket corresponding to the same state.**

**State corresponds to **direction** of ket vector.**

**The length is irrelevant.**

**The sign is irrelevant.**

**$|R\rangle \Rightarrow$  some state**

**$C|R\rangle = |R\rangle$  same state (Unless  $C = 0$ )**

**Ket**  $|A\rangle$             **number**  $\varphi$

$\varphi$  is a linear function of  $|A\rangle$

**If**

$|A\rangle \Rightarrow \varphi$  and  $|A'\rangle \Rightarrow \varphi'$

**then**

$|A\rangle + |A'\rangle \Rightarrow \varphi + \varphi'$

**and**

$C|A\rangle \Rightarrow C\varphi$

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**To get a number that is a linear function of a vector,  
take the scalar product with some other vector.**



**Other vector**  $\longrightarrow$  **Bra**

**symbol**  $\langle |$

**particular one**  $\langle \mathbf{B} |$

**Put a bra and a ket together**  $\langle bra | ket \rangle \Rightarrow$  **bracket**

**Any complete bracket**  $\langle | \rangle \Rightarrow$  **number**

**Incomplete bracket**  $\Rightarrow$  **vector**

$| \rangle$  **ket vector**

$\langle |$  **bra vector**

$\langle | \rangle \Rightarrow$  **scalar product, a complex number**

$$\langle \mathbf{B} | \{ | \mathbf{A} \rangle + | \mathbf{A}' \rangle \} = \langle \mathbf{B} | \mathbf{A} \rangle + \langle \mathbf{B} | \mathbf{A}' \rangle$$

**Scalar product with sum of kets.**

$$\langle \mathbf{B} | \{ C | \mathbf{A} \rangle \} = C \langle \mathbf{B} | \mathbf{A} \rangle$$

**Scalar product with constant times ket.**

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**Bra completely defined when its scalar product with every ket is known.**

$$\text{If } \langle \mathbf{P} | \mathbf{A} \rangle = 0 \text{ for all } | \mathbf{A} \rangle \Rightarrow \langle \mathbf{P} | = 0$$

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**Addition of bras defined by scalar product with ket  $| \mathbf{A} \rangle$ .**

$$\{ \langle \mathbf{B} | + \langle \mathbf{B}' | \} | \mathbf{A} \rangle = \langle \mathbf{B} | \mathbf{A} \rangle + \langle \mathbf{B}' | \mathbf{A} \rangle$$

**Multiplication of bra by a constant defined by scalar product with ket  $| \mathbf{A} \rangle$ .**

$$\{ C \langle \mathbf{B} | \} | \mathbf{A} \rangle = C \langle \mathbf{B} | \mathbf{A} \rangle$$

**Assume: There is a one to one correspondence between kets and bras such that the bra corresponding to  $|A\rangle + |A'\rangle$  is the sum of the bras corresponding to  $|A\rangle$  and  $|A'\rangle$  and that the bra corresponding to  $C|A\rangle$  is  $\bar{C}$  or  $C^*$  times the bra corresponding to  $|A\rangle$ .**

**$\bar{C}$  or  $C^*$  is the complex conjugate of the complex number  $C$ .**

**Bra is the complex conjugate of Ket.**

$$\langle A| = \overline{|A\rangle}$$

**Theory symmetrical in kets and bras. State of a system can be specified by the direction of a bra as well as a ket.**

Since

$$\langle A | = \overline{|A\rangle} \quad \text{and}$$

$$\langle B | = \overline{|B\rangle}$$

then

$$\langle B | A \rangle = \overline{\langle A | B \rangle}$$

The complex conjugate of a bracket is the bracket in “reverse order.”

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For  $|B\rangle = |A\rangle$

$$\langle A | A \rangle = \overline{\langle A | A \rangle}$$

$\Rightarrow \langle A | A \rangle$  is a real number because  
a number equal to its complex conjugate is real.

Assume  $\langle A | A \rangle > 0$       Length of a vector is positive unless  $|A\rangle = 0$

$$\left(\langle A | A \rangle\right)^{1/2} = \text{length of vector}$$

## Scalar Product

**Real space vectors** – scalar product gives a **real number**;  
**symmetrical with respect to interchange of order.**

$$\vec{A} \cdot \vec{B} = \vec{B} \cdot \vec{A}$$

**Kets and Bras** – scalar product gives **complex number**.

$$\langle A | B \rangle = \overline{\langle B | A \rangle}$$

**Interchange of order gives complex conjugate number.**

## Orthogonality

Bra and Ket orthogonal if  $\langle B | A \rangle = 0$

scalar product equals zero.

2 Bras or 2 Kets orthogonal;

$|B\rangle$  and  $|A\rangle$  orthogonal if

$$\langle B | A \rangle = 0$$

$\langle B |$  and  $\langle A |$  orthogonal if

$$\langle B | A \rangle = 0$$

Two states of a dynamical system are orthogonal if the vectors representing them are orthogonal.

## Normalization of Bras and Kets

Length of  $\langle A|$  or  $|A\rangle$  is

$$\left(\langle A|A\rangle\right)^{1/2}$$

**Length doesn't matter, only direction.**

May be convenient to have

$$\left(\langle A|A\rangle\right)^{1/2} = 1$$

**If this is done  $\Rightarrow$  States are normalized.**

## Normalized ket still not completely specified

A ket is a vector – direction

Normalized ket – direction and length

But, can multiply by  $e^{i\gamma}$   $\gamma$  a real number

$$\left(\langle A|A\rangle\right)^{1/2} = 1 \quad \text{normalized ket}$$

$$|A'\rangle = e^{i\gamma} |A\rangle \quad \text{multiply by } e^{i\gamma}$$

$$\langle A'| = e^{-i\gamma} \langle A| \quad \text{complex conjugate of } |A'\rangle$$

$$\left(\langle A'|A'\rangle\right)^{1/2} = \left[e^{-i\gamma} e^{i\gamma} \langle A|A\rangle\right]^{1/2} = 1 \quad \text{still normalized}$$

$e^{i\gamma}$  with  $\gamma$  real called a “phase factor.” Very important. See later.

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**State of system defined by direction of ket, not length.**



# Linear Operators

Kets and bras represent states of a dynamical system, s, p, d, etc states of H atom.

Need math relations to work with ket vectors to obtain observables.

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Ket  $|F\rangle$  is a function of ket  $|A\rangle$

$$|F\rangle = \underline{\alpha}|A\rangle$$

$\underline{\alpha}$  is a linear operator. (**Underline operator to indicate it is an operator**)

Linear operators have the properties:

$$\underline{\alpha}(|A\rangle + |A'\rangle) = \underline{\alpha}|A\rangle + \underline{\alpha}|A'\rangle$$

$$\underline{\alpha}(c|A\rangle) = c \underline{\alpha}|A\rangle$$

**complex number**

**A linear operator is completely defined when its application to every ket is known.**

# Rules about linear operators

## Additive

$$\{\underline{\alpha} + \underline{\beta}\} |A\rangle = \underline{\alpha} |A\rangle + \underline{\beta} |A\rangle$$

## Multiplication - associative

$$\{\underline{\alpha}\underline{\beta}\} |A\rangle = \underline{\alpha}\{\underline{\beta} |A\rangle\} = \underline{\alpha}\underline{\beta} |A\rangle$$

Multiplication is **NOT** necessarily commutative

$$\underline{\alpha}\underline{\beta} |A\rangle \neq \underline{\beta}\underline{\alpha} |A\rangle \quad \text{in general}$$

**If**  $\underline{\gamma} \underline{\delta} = \underline{\delta} \underline{\gamma}$

**then**  $\underline{\gamma}$  and  $\underline{\delta}$  commute.

Only know what an operator does by operating on a ket.

$\underline{\gamma} \underline{\delta} = \underline{\delta} \underline{\gamma}$  really means

$$\underline{\gamma} \underline{\delta} |A\rangle = \underline{\delta} \underline{\gamma} |A\rangle$$

$\underline{\gamma} \underline{\delta} = \underline{\delta} \underline{\gamma}$  is a short hand. It does not mean the right and left hand sides of the equation are equal algebraically.

**It means**  $\underline{\gamma} \underline{\delta}$  and  $\underline{\delta} \underline{\gamma}$

**have the same result when applied to an arbitrary ket.**

**Put ket on right of linear operator.**

$$\underline{\alpha}|A\rangle = |F\rangle$$

**Bras and linear operators**

$$\langle Q| = \langle B|\underline{\gamma} \quad \text{Put bra on left of linear operator.}$$

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$$\langle B|\{\underline{\alpha}|A\rangle\} = \langle B|\underline{\alpha}|A\rangle \quad \text{This is a number. It is a closed bracket} \longrightarrow \text{number.}$$

**To see this consider**

$$\begin{aligned} \langle B|\underline{\alpha}|A\rangle &= \langle B|\{\underline{\alpha}|A\rangle\} \\ &= \langle B|Q\rangle = c \quad \text{a complex number} \end{aligned}$$

## Kets and Bras as linear operators

$\langle A|B\rangle$  is a number

but

$|A\rangle\langle B|$  is a linear operator.

not closed bracket

Operate on an arbitrary ket  $|P\rangle$

$$\begin{aligned} |A\rangle\langle B|P\rangle &= |A\rangle\varphi \\ &= \varphi|A\rangle \end{aligned}$$

a number

Operating on  $|P\rangle$  gave  $|A\rangle$ , a new ket. Therefore,  $|A\rangle\langle B|$  is an operator.

Can also operate on a bra  $\langle Q|$ .

$$\langle Q|A\rangle\langle B| = \theta\langle B|$$

new bra

Will use later; projection operators

## Have algebra involving bras, kets and linear operators.

1. associative law of multiplication holds
2. distributive law holds
3. commutative law does not hold
4.  $\langle \rangle \Rightarrow$  number
5.  $| \rangle$  or  $\langle | \Rightarrow$  vector

**Assume: The linear operators correspond to the dynamical variables of a physical system.**

# Dynamical Variables

coordinates

components of velocity

momentum

angular momentum

energy

etc.

**Linear Operators  $\Rightarrow$  Questions you can ask about a system.**

**For every experimental observable, there is a linear operator.**

**Q. M. dynamical variables not subject to an algebra in which the commutative law of multiplication holds.**

**Consequence of Superposition Principle.**

**(Will lead to the Uncertainty Principle.)**

## Conjugate relations

Already saw that  $\langle B|A\rangle = \overline{\langle A|B\rangle}$

Complex conjugate – reverse order

### With linear operators

$$\langle B|\underline{\bar{\alpha}}|P\rangle = \overline{\langle P|\underline{\alpha}|B\rangle}$$

$\underline{\bar{\alpha}}$   $\longrightarrow$  complex conjugate of the operator  $\underline{\alpha}$ .

$\underline{\bar{\bar{\alpha}}} = \underline{\alpha}$   $\longrightarrow$  complex conjugate of complex conjugate

If  $\underline{\bar{\alpha}} = \underline{\alpha}$ , the operator is “self-adjoint.”  
It corresponds to a real dynamical variable.

$\underline{\bar{\alpha}} = \underline{\alpha}$  means  $\underline{\bar{\alpha}}|Q\rangle = \underline{\alpha}|Q\rangle$  all  $|Q\rangle$



$$\overline{\underline{\beta} \underline{\alpha}} = \underline{\underline{\alpha}} \underline{\underline{\beta}}$$

$$\overline{\underline{\alpha} \underline{\beta} \underline{\gamma}} = \underline{\underline{\gamma}} \underline{\underline{\beta}} \underline{\underline{\alpha}}$$

$$\overline{|A\rangle\langle B|} = |B\rangle\langle A|$$

**The complex conjugate of any product of bras, kets, and linear operators is the complex conjugate of each factor with factors in reverse order.**

# Eigenvalues and Eigenvectors

$$\underline{\alpha} |P\rangle = p |P\rangle$$

linear operator   ket   number (real)   same ket

**Eigenvalue problem**      (Mathematical problem in linear algebra.)

Know  $\underline{\alpha}$ .

want to find  $|P\rangle$  and  $p$ .

eigenvector      eigenvalue

**Apply linear operator to ket, get same ket back multiplied by a number.**

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Can also have  $\langle Q | \underline{\alpha} = q \langle Q |$

$\langle Q |$  is an eigenbra of  $\underline{\alpha}$ .

# Observables and Linear Operators

**Kets (or bras)**  $\longrightarrow$  **State of dynamical system.**

**Linear operators**  $\longrightarrow$  **Dynamical variables.**  
**Questions you can ask about the state of a system.**

**Observables**  $\longrightarrow$  **Real Dynamical Variables.**

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**Observables are the eigenvalues of Hermitian Linear Operators.**

**Hermitian operators**  $\longrightarrow$  **Real Dynamical Variables.**

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**Hermitian operator**

$$\langle a | \underline{\gamma} | b \rangle = \langle a | \overline{\gamma} | b \rangle$$

**If**  $|b\rangle = |a\rangle$

$$\langle a | \overline{\gamma} | a \rangle = \langle a | \underline{\gamma} | a \rangle$$

**Observables  $\Rightarrow$**

**Eigenvalues of Hermitian Operators (real eigenvalues)**

$$\underline{\alpha} |P\rangle = p |P\rangle$$

**linear operator**  $\rightarrow$   $\underline{\alpha}$

**eigenvector  
eigenket**  $\rightarrow$   $|P\rangle$

**observable value**  $\rightarrow$   $p$

**For every observable there is a linear operator**

**(more than one form, different representations of QM).**

**The eigenvalues are the numbers you will measure  
in an experiment.**

## Some important theorems and proofs

If an eigenvector of  $\underline{\alpha}$  is multiplied by any number  $C$  other than zero, then the product is still an eigenvector with the same eigenvalue. (Length doesn't matter, only direction.)

$$\underline{\alpha} |A\rangle = a |A\rangle$$

$$\begin{aligned} \underline{\alpha} [C |A\rangle] &= C [\underline{\alpha} |A\rangle] \longleftarrow \text{A constant commutes with an operator. A constant can be thought of as a number times the identity operator. The identity operator commutes with any other operator.} \\ &= C [a |A\rangle] \\ &= a [C |A\rangle] \end{aligned}$$

This is the reason that the direction, not the length, of a vector describes the state of a system. Eigenvalues (observable values) not changed by length of vector as long as  $C \neq 0$  .

**Several independent eigenkets of an operator  $\underline{\alpha}$  can have the same eigenvalues.  $\longrightarrow$  Degenerate**

**Any superposition of these eigenkets is an eigenket with the same eigenvalue.**

$$\underline{\alpha}|P_1\rangle = p|P_1\rangle$$

$$\underline{\alpha}|P_2\rangle = p|P_2\rangle$$

$$\underline{\alpha}|P_3\rangle = p|P_3\rangle$$

$$\begin{aligned}\underline{\alpha}\{C_1|P_1\rangle + C_2|P_2\rangle + C_3|P_3\rangle\} &= C_1\underline{\alpha}|P_1\rangle + C_2\underline{\alpha}|P_2\rangle + C_3\underline{\alpha}|P_3\rangle \\ &= p\{C_1|P_1\rangle + C_2|P_2\rangle + C_3|P_3\rangle\}\end{aligned}$$

**Example:  $p_x$ ,  $p_y$ , and  $p_z$  orbitals of the H atom.**

The eigenvalues associated with the eigenkets are the same as those for the eigenbras for the same linear operator.

Assume not the same.

$$\underline{\gamma} |P\rangle = a |P\rangle$$

$$\langle P | \underline{\gamma} = b \langle P | \quad \bar{\underline{\gamma}} = \underline{\gamma} \quad \text{Hermitian property}$$

Left multiply top by  $\langle P |$

Right multiply bottom by  $|P\rangle$

$$\langle P | \underline{\gamma} |P\rangle = a \langle P | P\rangle$$

$$\langle P | \underline{\gamma} |P\rangle = b \langle P | P\rangle$$

Subtract

$$0 = (a - b) \langle P | P\rangle$$

But,  $\langle P | P\rangle > 0$

Therefore,  $a = b$

Theory is symmetrical in kets and bras.

**The eigenvalues of Hermitian operators are real.**

$$\underline{\gamma} |P\rangle = a |P\rangle \quad \text{take complex conjugate}$$

$$\langle P | \underline{\gamma} = \langle P | \bar{a}$$

$$\langle P | \underline{\gamma} = \langle P | \bar{a} \quad \text{Hermitian property}$$

**Left multiply the first equation by  $\langle P |$**

**Right multiply third equation by  $|P\rangle$**

$$\langle P | \underline{\gamma} |P\rangle = a \langle P | P\rangle$$

$$\langle P | \underline{\gamma} |P\rangle = \bar{a} \langle P | P\rangle$$

**Subtract**

$$\mathbf{0} = (a - \bar{a}) \langle P | P\rangle$$

$$\langle P | P\rangle > \mathbf{0} \quad \text{Length of vector greater than zero.}$$

**Therefore,  $\bar{a} = a$ ; and a number equal to its complex conjugate is real.  
Observables are real numbers.**



## The Orthogonality Theorem

**Two eigenvectors of a real dynamical variable (Hermitian operator - observable) belonging to different eigenvalues are orthogonal.**

$$\underline{\gamma}|\gamma'\rangle = \gamma'|\gamma'\rangle \quad (1)$$

$$\underline{\gamma}|\gamma''\rangle = \gamma''|\gamma''\rangle \quad (2)$$

**Take complex conjugate of (1).**

$$\langle\gamma'|\underline{\gamma} = \langle\gamma'|\gamma' \longleftarrow \text{Operator Hermitian; eigenvalue real}$$

$$\langle\gamma'|\underline{\gamma}|\gamma''\rangle = \gamma'\langle\gamma'|\gamma''\rangle \quad (3) \quad \text{Right multiply by } |\gamma''\rangle$$

$$\langle\gamma'|\underline{\gamma}|\gamma''\rangle = \gamma''\langle\gamma'|\gamma''\rangle \quad (4) \quad \text{Left multiply (2) by } \langle\gamma'|$$

$$\mathbf{0} = (\gamma' - \gamma'')\langle\gamma'|\gamma''\rangle \quad \text{Subtract (4) from (3)}$$

**But**  $\gamma' \neq \gamma''$  **different eigenvalues**

$$\text{Therefore, } \langle\gamma'|\gamma''\rangle = \mathbf{0}$$

$|\gamma'\rangle$  **and**  $|\gamma''\rangle$  **are orthogonal.**