

Chapter 9

Non-degenerate Perturbation Theory

Problem :

$$\underline{H}|\varphi_n\rangle = E_n|\varphi_n\rangle$$

can't solve exactly.

But

$$\underline{H} = \underline{H}^0 + \lambda \underline{H}' + \lambda^2 \underline{H}'' + \dots$$

with

$$\lim_{\lambda \rightarrow 0} \underline{H}|\varphi_n^0\rangle = E_n^0|\varphi_n^0\rangle$$

Unperturbed eigenvalue problem.

Can solve exactly.

Therefore, know E_n^0 and $|\varphi_n^0\rangle$.

$$\lambda \underline{H}' + \lambda^2 \underline{H}'' + \dots$$

called perturbations

Solutions of

$$\underline{H}^0 |\varphi_n^0\rangle = E_n^0 |\varphi_n^0\rangle$$

complete, orthonormal set of ket vectors

$$\{|\varphi_n^0\rangle\} \Rightarrow |\varphi_0^0\rangle, |\varphi_1^0\rangle, |\varphi_2^0\rangle \dots$$

with eigenvalues $E_0^0, E_1^0, E_2^0, \dots$

and

$$\langle \varphi_n^0 | \varphi_m^0 \rangle = \delta_{mn}$$

Kronecker delta

$$\delta_{nm} = \begin{cases} 1 & n = m \\ 0 & n \neq m \end{cases}$$

Expand wavefunction

$$|\varphi_n\rangle = |\varphi_n^0\rangle + \lambda |\varphi_n'\rangle + \lambda^2 |\varphi_n''\rangle + \dots$$

and

$$E_n = E_n^0 + \lambda E_n' + \lambda^2 E_n'' + \dots$$

also have

$$\underline{H} = \underline{H}^0 + \lambda \underline{H}' + \lambda^2 \underline{H}'' + \dots$$

Have series for

$$\underline{H} \quad |\varphi_n\rangle \quad E_n$$

Substitute these series into the original eigenvalue equation

$$\underline{H} |\varphi_n\rangle = E_n |\varphi_n\rangle$$

Sum of infinite number of terms for all powers of λ equals 0.

$$\left(\underline{H}^0 |\varphi_n^0\rangle - E_n^0 |\varphi_n^0\rangle \right) \lambda^0 + \left(\underline{H}^0 |\varphi_n'\rangle + \underline{H}' |\varphi_n^0\rangle - E_n^0 |\varphi_n'\rangle - E_n' |\varphi_n^0\rangle \right) \lambda$$

$$+ \left(\underline{H}^0 |\varphi_n''\rangle + \underline{H}' |\varphi_n'\rangle + \underline{H}'' |\varphi_n^0\rangle - E_n^0 |\varphi_n''\rangle - E_n' |\varphi_n'\rangle - E_n'' |\varphi_n^0\rangle \right) \lambda^2$$

$$+ \dots = 0$$

Coefficients of the individual powers of λ must equal 0.

zeroth

order - λ^0

$$\underline{H}^0 |\varphi_n^0\rangle - E_n^0 |\varphi_n^0\rangle = 0$$

first

order - λ^1

$$\underline{H}^0 |\varphi_n'\rangle + \underline{H}' |\varphi_n^0\rangle - E_n^0 |\varphi_n'\rangle - E_n' |\varphi_n^0\rangle = 0$$

second

order - λ^2

$$\underline{H}^0 |\varphi_n''\rangle + \underline{H}' |\varphi_n'\rangle + \underline{H}'' |\varphi_n^0\rangle - E_n^0 |\varphi_n''\rangle - E_n' |\varphi_n'\rangle - E_n'' |\varphi_n^0\rangle = 0$$

First order correction

$$\underline{H}^0 |\varphi'_n\rangle - E'_n |\varphi'_n\rangle = (\underline{E}'_n - \underline{H}') |\varphi_n^0\rangle$$

Want to find E'_n and $|\varphi'_n\rangle$.

Expand $|\varphi'_n\rangle$

$$|\varphi'_n\rangle = \sum_i c_i |\varphi_i^0\rangle \quad \text{also substituting}$$

Then

$$\underline{H}^0 |\varphi'_n\rangle = \sum_i c_i \underline{H}^0 |\varphi_i^0\rangle = \sum_i c_i E_i^0 |\varphi_i^0\rangle \quad \text{Substituting this result.}$$

After substitution

$$\sum_i c_i (E_i^0 - E'_n) |\varphi_i^0\rangle = (\underline{E}'_n - \underline{H}') |\varphi_n^0\rangle$$

After substitution

$$\sum_i c_i (E_i^0 - E_n^0) |\varphi_i^0\rangle = (E_n' - \underline{H}') |\varphi_n^0\rangle$$

Left multiply by

$$\langle \varphi_n^0 |$$

$$\langle \varphi_n^0 | \sum_i c_i (E_i^0 - E_n^0) |\varphi_i^0\rangle = \langle \varphi_n^0 | (E_n' - \underline{H}') |\varphi_n^0\rangle$$

$$\sum_i c_i (E_i^0 - E_n^0) \langle \varphi_n^0 | \varphi_i^0 \rangle = \langle \varphi_n^0 | (E_n' - \underline{H}') |\varphi_n^0\rangle$$


$\langle \varphi_n^0 | \varphi_i^0 \rangle = 0$ unless $n = i$,
but then

$$E_n^0 - E_n^0 = 0$$

Therefore, the left side is 0.

We have

$$\langle \varphi_n^0 | (E'_n - \underline{H}') | \varphi_n^0 \rangle = 0$$


$$\langle \varphi_n^0 | E'_n | \varphi_n^0 \rangle - \langle \varphi_n^0 | \underline{H}' | \varphi_n^0 \rangle = 0$$



E'_n number, kets normalized, and transposing,

$$E'_n = \langle \varphi_n^0 | \underline{H}' | \varphi_n^0 \rangle$$

The first order correction to the energy.
(Expectation value of \underline{H}' in zeroth order state φ_n^0)

Then

$$E_n = E_n^0 + \lambda E'_n$$

Absorbing λ into H'_{nn} and E'_n

$$E_n = E_n^0 + E'_n$$

$$E'_n = \langle \varphi_n^0 | \underline{H}' | \varphi_n^0 \rangle = H'_{nn}$$

The first order correction to the energy is the expectation value of \underline{H}' .

First order correction to the wavefunction

Again using the equation obtained after substituting series expansions

$$\sum_i c_i (E_i^0 - E_n^0) |\varphi_i^0\rangle = (E_n' - \underline{H}') |\varphi_n^0\rangle$$

Left multiply by $\langle \varphi_j^0 |$

$$\langle \varphi_j^0 | \sum_i c_i (E_i^0 - E_n^0) |\varphi_i^0\rangle = \langle \varphi_j^0 | (E_n' - \underline{H}') |\varphi_n^0\rangle$$

← Equals zero unless $i = j$.

$$c_j (E_j^0 - E_n^0) = \langle \varphi_j^0 | (E_n' - \underline{H}') |\varphi_n^0\rangle$$

$$c_j (E_j^0 - E_n^0) = -\langle \varphi_j^0 | \underline{H}' | \varphi_n^0 \rangle$$

$$c_j = \frac{\langle \varphi_j^0 | \underline{H}' | \varphi_n^0 \rangle}{(E_n^0 - E_j^0)} \quad j \neq n$$

Coefficients in expansion of ket in terms of the zeroth order kets.

$$c_j = \frac{\langle \varphi_j^0 | \underline{H}' | \varphi_n^0 \rangle}{(E_n^0 - E_j^0)} \quad j \neq n$$

$$c_j = \frac{H'_{jn}}{(E_n^0 - E_j^0)} \quad H'_{jn} \text{ is the bracket of } \underline{H}' \text{ with } \langle \varphi_j^0 | \text{ and } | \varphi_n^0 \rangle .$$

Therefore

$$|\varphi_n\rangle = |\varphi_n^0\rangle + \sum_j' \frac{H'_{jn}}{(E_n^0 - E_j^0)} |\varphi_j^0\rangle$$

correction to zeroth order ket

The prime on the sum mean
 $j \neq n$.

zeroth order ket

energy denominator

First order corrections

$$E_n = E_n^0 + H'_{nn} + \dots$$

$$H'_{nn} = \langle \varphi_n^0 | H' | \varphi_n^0 \rangle$$

$$|\varphi_n\rangle = |\varphi_n^0\rangle + \sum_j' \frac{H'_{jn}}{(E_n^0 - E_j^0)} |\varphi_j^0\rangle + \dots$$

$$H'_{jn} = \langle \varphi_j^0 | H' | \varphi_n^0 \rangle$$

Second Order Corrections

Using λ^2 coefficient

Expanding $|\varphi'_n\rangle$ $|\varphi''_n\rangle$

Substituting and following same type of procedures yields

$$E''_n = \sum'_i \frac{H'_{ni}H'_{in}}{(E_n^0 - E_i^0)} + H''_{nn}$$

λ^2 coefficients have been absorbed.

$$H'_{ni}H'_{in} = \langle \varphi_n^0 | \underline{H}' | \varphi_i^0 \rangle \langle \varphi_i^0 | \underline{H}' | \varphi_n^0 \rangle$$

Second order correction due to first order piece of \underline{H} .

Second order correction due to an additional second order piece of \underline{H} .

$$|\varphi''_n\rangle = \sum'_k \left[\sum'_m \frac{H'_{km}H'_{mn}}{(E_n^0 - E_k^0)(E_n^0 - E_m^0)} - \frac{H'_{nn}H'_{kn}}{(E_n^0 - E_k^0)^2} \right] |\varphi_k^0\rangle + \sum'_k \frac{H''_{kn}}{(E_n^0 - E_k^0)} |\varphi_k^0\rangle$$

Second order correction due to first order piece of \underline{H} .

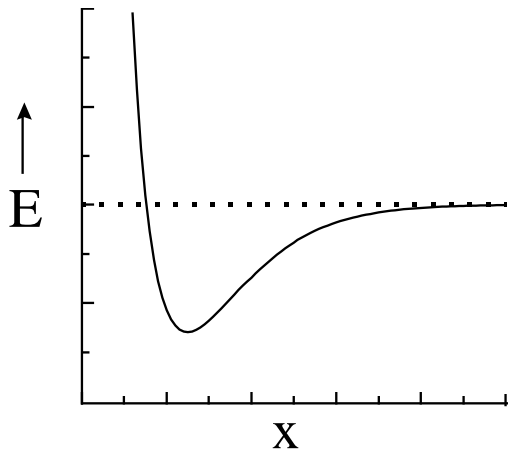
Second order correction due to an additional second order piece of \underline{H} .

Energy and Ket Corrected to First and Second Order

$$E = E^0 + H'_{nn} + \sum_i' \frac{H'_{ni}H'_{in}}{(E_n^0 - E_i^0)} + H''_{nn} + \dots$$

$$|\varphi_n\rangle = |\varphi_n^0\rangle + \sum_j' \frac{H'_{jn}}{(E_n^0 - E_j^0)} |\varphi_j^0\rangle + \sum_k' \left[\sum_m' \frac{H'_{km}H'_{mn}}{(E_n^0 - E_k^0)(E_n^0 - E_m^0)} - \frac{H'_{nn}H'_{kn}}{(E_n^0 - E_k^0)^2} \right] |\varphi_k^0\rangle \\ + \sum_k' \frac{H''_{kn}}{(E_n^0 - E_k^0)} |\varphi_k^0\rangle + \dots$$

Example: x^3 and x^4 perturbation of the Harmonic Oscillator



Vibrational potential of molecules not harmonic.

Approximately harmonic near potential minimum.

Expand potential in power series.

First additional terms in potential after x^2 term are x^3 and x^4 .

$$\underline{H} = \frac{\underline{p}^2}{2m} + \frac{1}{2}k\underline{x}^2 + c\underline{x}^3 + q\underline{x}^4$$

quartic “force constant”

cubic “force constant”

$$\underline{H}^0 = \frac{\underline{p}^2}{2m} + \frac{1}{2}k\underline{x}^2$$

harmonic oscillator – know solutions

$$\underline{H}^0 = \frac{1}{2}\hbar\omega(\underline{a}\underline{a}^+ + \underline{a}^+\underline{a})$$

$$E^0 = \left(n + \frac{1}{2}\right)\hbar\omega_0$$

zeroth order eigenvalues

$$|n\rangle$$

zeroth order eigenkets

$$\underline{H}' = c\underline{x}^3 + q\underline{x}^4$$

perturbation

c and q are expansion coefficients like λ .

When c and $q \rightarrow 0$, $\underline{H} \rightarrow \underline{H}_0$

$$\begin{aligned}
H'_{nn} &= \langle n | \underline{H}' | n \rangle \\
&= \langle n | c \underline{x}^3 + q \underline{x}^4 | n \rangle \\
&= c \langle n | \underline{x}^3 | n \rangle + q \langle n | \underline{x}^4 | n \rangle
\end{aligned}$$

In Dirac representation

$$\underline{x} = \left(\frac{\hbar \omega_0}{2k} \right)^{\frac{1}{2}} (\underline{a} + \underline{a}^+)$$

First consider cubic term.

$$\underline{x}^3 \propto (\underline{a} + \underline{a}^+)^3$$

Multiply out. Many terms.

$$\underline{a}^3, \underline{a}^2 \underline{a}^+, \underline{a} \underline{a}^+ \underline{a}, \dots, \underline{a}^{\dagger 3}.$$

None of the terms have the same number of raising and lowering operators.

$$\langle n | \underline{x}^3 | n \rangle = 0 \quad (\text{At second order will not be zero.})$$

$$\langle n | \underline{x}^4 | n \rangle = \frac{\hbar^2 \omega_0^2}{4k^2} \langle n | (\underline{a} + \underline{a}^+)^4 | n \rangle$$

$(\underline{a} + \underline{a}^+)^4$ has terms with same number of raising and lowering operators.

Therefore, $\langle n | \underline{x}^4 | n \rangle \neq 0$

Using $a | n \rangle = n^{1/2} | n-1 \rangle$ and $a^+ | n \rangle = (n+1)^{1/2} | n+1 \rangle$

$$\langle n | \underline{a} \underline{a} \underline{a}^+ \underline{a}^+ | n \rangle = (n+1)(n+2)$$

$$\langle n | \underline{a}^+ \underline{a}^+ \underline{a} \underline{a} | n \rangle = n(n-1)$$

$$\langle n | \underline{a} \underline{a}^+ \underline{a} \underline{a}^+ | n \rangle = (n+1)^2$$

$$\langle n | \underline{a}^+ \underline{a} \underline{a}^+ \underline{a} | n \rangle = n^2$$

$$\langle n | \underline{a} \underline{a}^+ \underline{a}^+ \underline{a} | n \rangle = n(n+1)$$

$$\langle n | \underline{a}^+ \underline{a} \underline{a} \underline{a}^+ | n \rangle = (n+1)n$$

Only terms with the same number of raising and lowering operators are non-zero.

There are six terms.

Sum of the six terms

$$\langle n | (\underline{a} + \underline{a}^+)^4 | n \rangle = 6(n^2 + n + 1/2)$$

Therefore

$$H'_{nn} = \frac{q \hbar^2 \omega_0^2}{k^2} \frac{3}{2} \left(n^2 + n + \frac{1}{2} \right)$$

With $\omega_0 = \sqrt{k/m}$ $k^2 = \omega_0^4 m^2$

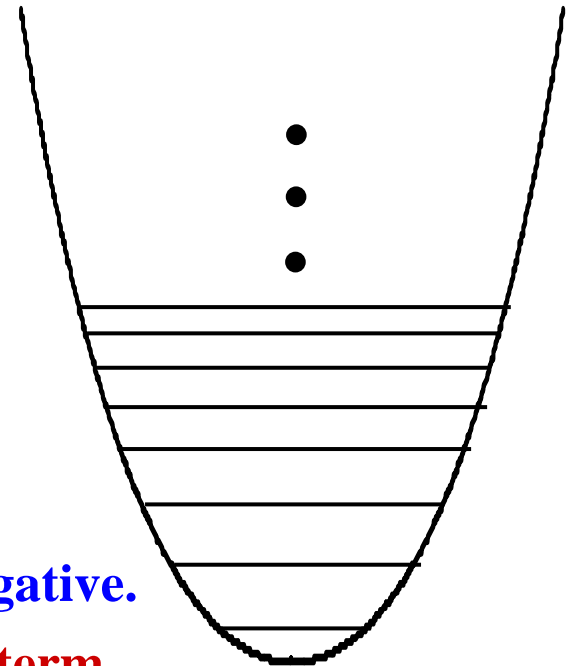
$$E_n = \left(n + \frac{1}{2} \right) \hbar \omega_0 + q \frac{3}{2} \left(n^2 + n + \frac{1}{2} \right) \frac{\hbar^2}{m^2 \omega_0^2}$$

Energy levels not equally spaced.

Real molecules, levels get closer together – q is negative.

Correction grows with n faster than zeroth order term

→ decrease in level spacing.



Perturbation Theory for Degenerate States

$$\underline{H}|\varphi_1\rangle = E|\varphi_1\rangle \quad |\varphi_1\rangle \quad \text{and} \quad |\varphi_2\rangle$$

$$\underline{H}|\varphi_2\rangle = E|\varphi_2\rangle$$

normalize and orthogonal

$|\varphi_1\rangle$ and $|\varphi_2\rangle$ **Degenerate, same eigenvalue, E .**

If $|\psi\rangle = c_1|\varphi_1\rangle + c_2|\varphi_2\rangle$

with $\bar{c}_1 c_1 + \bar{c}_2 c_2 = 1$

$$\underline{H}|\psi\rangle = E|\psi\rangle$$

Any superposition of degenerate eigenstates is also an eigenstate with the same eigenvalue.

n linearly independent states with same eigenvalue



system n -fold degenerate

Can form n orthonormal $|\psi_i\rangle$ from the n degenerate $|\varphi_n\rangle$.

Can form an infinite number of sets of $|\psi_i\rangle$.

Nothing unique about any one set of n degenerate eigenkets.

Want approximate solution to

$$\left(\underline{H}^0 + \lambda \underline{H}'\right) |\varphi_j\rangle = E_j |\varphi_j\rangle$$

zeroth order
Hamiltonian

perturbation

$$\lambda \rightarrow 0 \quad \underline{H}^0 |\varphi_j^0\rangle = E_j^0 |\varphi_j^0\rangle$$

zeroth order
eigenket

zeroth order
energy

But E_i^0 is m -fold degenerate.

Call these m eigenkets belonging to the m -fold degenerate E_1^0

$|\varphi_1^0\rangle, |\varphi_2^0\rangle \cdots |\varphi_m^0\rangle$ orthonormal

With $E_1^0 = E_2^0 = \cdots = E_m^0 \equiv E_1^0$

Here is the difficulty

$$\lambda \rightarrow 0$$

$$|\varphi_i\rangle \rightarrow |\psi_i^0\rangle$$

perturbed ket **zeroth order ket having eigenvalue, E_1^0**

But, $|\psi_i^0\rangle$ is a linear combination of the $|\varphi_i^0\rangle$.

$$|\psi_i^0\rangle = c_1 |\varphi_1^0\rangle + c_2 |\varphi_2^0\rangle + \cdots + c_m |\varphi_m^0\rangle$$

We don't know which particular linear combination it is.

$|\psi_i^0\rangle$ is the correct zeroth order ket, but we don't know the c_i .

The correct zero order ket depends on the nature of the perturbation.

p states of the H atom in external

magnetic field – p_1, p_0, p_{-1}

electric field – p_x, p_z, p_y

To solve problem

Expand E and $|\varphi_i\rangle$

$$E = E_1^0 + \lambda E' + \dots$$

$$|\varphi_i\rangle = \sum_{j=1}^m c_j |\varphi_j^0\rangle + \lambda |\varphi_i'\rangle + \dots$$

Some superposition, but we don't know the c_j .

→ Don't know correct zeroth order function.

Substituting the expansions for E and $|\varphi_i\rangle$ into

$$(\underline{H}^0 + \lambda \underline{H}') |\varphi_i\rangle = E_i |\varphi_i\rangle$$

and obtaining the coefficients of powers of λ , gives

zeroth
order

$$\underline{H}^0 \sum_{j=1}^m c_j |\varphi_j^0\rangle = E_1^0 \sum_{j=1}^m c_j |\varphi_j^0\rangle$$

first
order

$$(\underline{H}^0 - E_1^0) |\varphi_i'\rangle = \sum_{j=1}^m c_j (\underline{E}' - \underline{H}') |\varphi_j^0\rangle$$

want these

$$\left(\underline{H}^0 - E_1^0\right)|\varphi_i'\rangle = \sum_{j=1}^m c_j (E' - \underline{H}')|\varphi_j^0\rangle \quad \text{To solve}$$

substitute $|\varphi_i'\rangle = \sum_k A_k |\varphi_k^0\rangle$

Need $\underline{H}'|\varphi_j^0\rangle$

Use projection operator $|\varphi_k^0\rangle\langle\varphi_k^0|$

$$\underline{H}'|\varphi_j^0\rangle = \sum_k |\varphi_k^0\rangle\langle\varphi_k^0|\underline{H}'|\varphi_j^0\rangle$$

The projection operator gives the piece of $\underline{H}'|\varphi_j^0\rangle$ that is $|\varphi_k^0\rangle$.

Then the sum over all k gives the expansion of $\underline{H}'|\varphi_j^0\rangle$ in terms of the $|\varphi_i^0\rangle$.

Defining $\underline{H}'_{kj} = \langle\varphi_k^0|\underline{H}'|\varphi_j^0\rangle$

Known – know perturbation piece of the Hamiltonian and the zeroth order kets.

$$\underline{H}'|\varphi_j^0\rangle = \sum_k \underline{H}'_{kj} |\varphi_k^0\rangle$$

$$\left(\underline{H}^0 - E_1^0\right)|\varphi_i'\rangle = \sum_{j=1}^m \underline{c}_j \left(\underline{E}' - \underline{H}'\right)|\varphi_j^0\rangle$$

$$\underline{H}'|\varphi_j^0\rangle = \sum_k H'_{kj} |\varphi_k^0\rangle$$

this piece becomes

$$\sum_{j=1}^m \underline{c}_j \underline{H}'|\varphi_j^0\rangle = \sum_{j=1}^m \sum_k \underline{c}_j H'_{kj} |\varphi_k^0\rangle$$

Substituting this and $|\varphi_i'\rangle = \sum_k A_k |\varphi_k^0\rangle$ **gives**

$$\sum_k \left(E_k^0 - E_1^0\right) A_k |\varphi_k^0\rangle = \sum_{j=1}^m E' c_j |\varphi_j^0\rangle - \sum_k \left(\sum_{j=1}^m \underline{c}_j H'_{kj}\right) |\varphi_k^0\rangle$$

Result of operating \underline{H}^0 on the zeroth order kets.

Left multiplying by $\langle\varphi_i^0|$

$$\sum_k \left(E_k^0 - E_1^0\right) A_k \langle\varphi_i^0|\varphi_k^0\rangle = \sum_{j=1}^m E' c_j \langle\varphi_i^0|\varphi_j^0\rangle - \sum_k \left(\sum_{j=1}^m \underline{c}_j H'_{kj}\right) \langle\varphi_i^0|\varphi_k^0\rangle$$

$$\sum_k (E_k^0 - E_1^0) A_k \langle \varphi_i^0 | \varphi_k^0 \rangle = \sum_{j=1}^m E' c_j \langle \varphi_i^0 | \varphi_j^0 \rangle - \sum_k \left(\sum_{j=1}^m c_j H'_{kj} \right) \langle \varphi_i^0 | \varphi_k^0 \rangle$$

Correction to the Energies

Two cases: $i \leq m$ (the degenerate states) and $i > m$.

$i \leq m$

Left hand side – sum over k equals zero unless $k = i$.

But with $i \leq m$,

$$E_i^0 = E_1^0 \quad \text{Therefore, } E_i^0 - E_1^0 = 0$$

The left hand side of the equation = 0.

Right hand side, first term non-zero when $j = i$. Bracket = 1, normalization.

Second term non-zero when $k = i$. Bracket = 1, normalization.

The result is

$$\sum_{j=1}^m H'_{ij} c_j - E' c_i = 0$$

We don't know the c 's and the E 's .

$\sum_{j=1}^m H'_{ij} c_j - E' c_i = 0$ is a system of m of equations for the c_j 's.

$$(H'_{11} - E')c_1 + H'_{12}c_2 + \cdots + H'_{1m}c_m = 0$$

$$H'_{21}c_1 + (H'_{22} - E')c_2 + \cdots + H'_{2m}c_m = 0$$

•
•
•

$$H'_{m1}c_1 + H'_{m2}c_2 + \cdots + (H'_{mm} - E')c_m = 0$$

One equation for each index i of c_i .

Besides trivial solution of $c_1 = c_2 = \cdots = c_m = 0$

only get solution if the determinant of the coefficients vanish.

$$\begin{vmatrix} (H'_{11} - E') & H'_{12} & \cdots & H'_{1m} \\ \vdots & (H'_{22} - E') & \cdots & H'_{2m} \\ & \vdots & & \vdots \\ H'_{m1} & H'_{m2} & \cdots & (H'_{mm} - E') \end{vmatrix} = 0$$

We know the

$$H'_{jk} = \langle \varphi_j^0 | H' | \varphi_k^0 \rangle$$

Have m^{th} degree equation for the E 's .

Solve m^{th} degree equation – get the E'_i 's . Now have the corrections to energies.

To find the correct zeroth order eigenvectors, one for each E'_i , substitute E'_i (one at a time) into system of equations.

Get system of equations for the coefficients, c_j 's.

$$(H'_{11} - E'_i)c_1 + H'_{12}c_2 + \cdots + H'_{1m}c_m = 0$$

Know the H'_{ij} .

$$H'_{21}c_1 + (H'_{22} - E'_i)c_2 + \cdots + H'_{2m}c_m = 0$$

•
•
•

$$H'_{m1}c_1 + H'_{m2}c_2 + \cdots + (H'_{mm} - E'_i)c_m = 0$$

There are only $m - 1$ conditions because can multiply everything by constant. Use normalization for m^{th} condition.

$$c_1^*c_1 + c_2^*c_2 + \cdots + c_m^*c_m = 1$$

Now we have the correct zeroth order functions.

The solutions to the m^{th} degree equation (expanding determinant) are

$$E'_1, E'_2, \dots, E'_m$$

Therefore, to first order, the energies of the perturbed initially degenerate states are

$$E_i = E_1^0 + E'_i \quad 1 \leq i \leq m$$

Have m different E'_i 's (unless some still degenerate).

With $E_i \rightarrow E_1^0$

as $\lambda \rightarrow 0$

Correction to wavefunctions

Again using equation found substituting the expansions into the first order equation

$$\sum_k (E_k^0 - E_1^0) A_k |\varphi_k^0\rangle = \sum_{j=1}^m E' c_j |\varphi_j^0\rangle - \sum_k \left(\sum_{j=1}^m c_j H'_{kj} \right) |\varphi_k^0\rangle$$

Left multiply by

$$\langle \varphi_i^0 | \quad i = k > m$$

gives 1

gives 0

Orthogonality makes other terms zero.

Normalization gives 1 for non-zero brackets.

$$(E_k^0 - E_1^0) A_k = - \sum_{j=1}^m c_j H'_{kj}$$

Therefore

$$A_k = \frac{\sum_{j=1}^m c_j H'_{kj}}{(E_1^0 - E_k^0)} \quad k > m$$

Normalization gives $A_j = 0$ for $j \leq m$.

Already have part of wavefunction for $j \leq m$

First order degenerate perturbation theory results

$$E_i = E_1^0 + \lambda E_i' + \dots$$

$$|\varphi_i\rangle = |\psi_i^0\rangle + \lambda \sum_{k>m} \frac{\sum_{j=1}^m c_j H'_{kj}}{(E_1^0 - E_k^0)} |\varphi_k^0\rangle + \dots$$

**Correct zeroth order function.
Coefficients c_k determined from
system of equations.**

**Correction to
zeroth order function.**