

# Solutions For Homework #3

## Problem 1:[10 pts]

- (a) The diagram in Figure 1 shows the two-dimensional function  $f(x, y) = \text{rect}(x - 2, y - 3) + 3\text{rect}(x - 2, y - 1)$  and the two projections  $P_0 f(x, y)$  and  $P_{90} f(x, y)$ . We can evaluate the projection  $P_0 f(x, y)$  mathematically

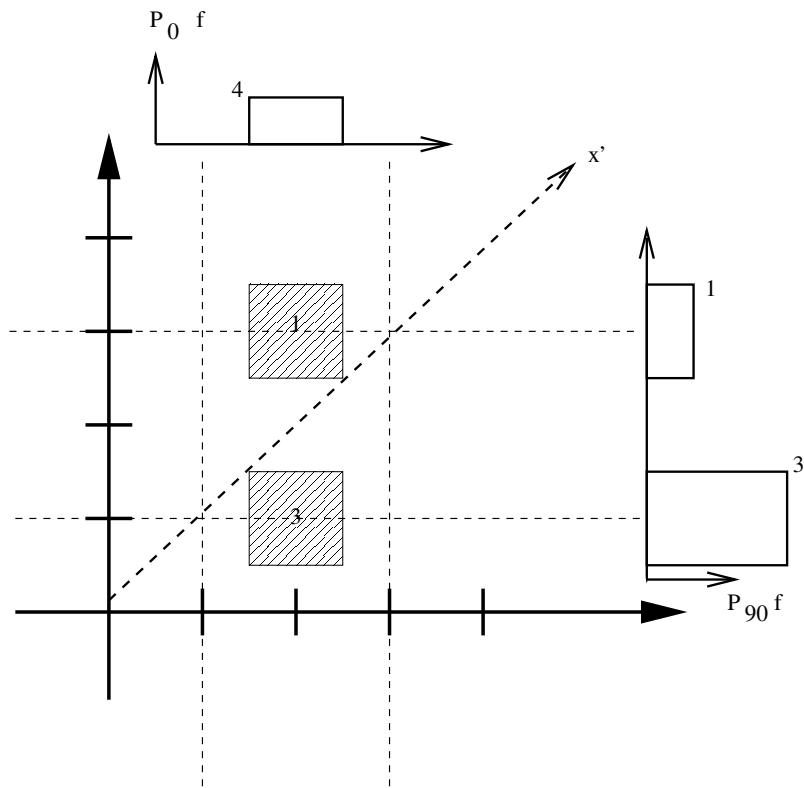


Figure 1:

as follows

$$\begin{aligned}
 P_0(R) &= \int_{-\infty}^{\infty} f(R, y) dy & (1) \\
 &= \int_{-\infty}^{\infty} \text{rect}(R - 2, y - 3) + 3\text{rect}(R - 2, y - 1) \\
 &= 4\text{rect}(R - 2)
 \end{aligned}$$

where  $P_0(R)$  is the projection of the two dimensional function  $f(x, y)$  along  $x = R$ .

- (b) Figure 1 shows the two dimensional function  $f(x, y) = \text{rect}(x - 2, y - 3) + 3\text{rect}(x - 2, y - 1)$  and the projection along a line parallel to the y-axis,  $P_{90} f(x, y)$ . Analogous to part (a), we compute the projection mathematically as follows

$$\begin{aligned}
 P_{90}(R) &= \int_{-\infty}^{\infty} f(x, R) dx & (2) \\
 &= \int_{-\infty}^{\infty} \text{rect}(x - 2, R - 3) + 3\text{rect}(x - 2, R - 1) \\
 &= \text{rect}(R - 3) + 3\text{rect}(R - 1)
 \end{aligned}$$

where  $P_{90}(R)$  is the projection of the two dimensional function  $f(x, y)$  along  $y = R$ .

- (c) In Figure 1, we see a line oriented 45 degrees from the x-axis, counterclockwise, labeled  $x'$ . This is the projection direction from which the one dimensional function  $P_{45} f(x, y)$  will be computed. The most direct way to compute this diagonal projection is to rotate the given  $x - y$  coordinate system, counterclockwise, by 45 degrees. We find the corresponding location  $(x', y')$  of a point  $(x, y)$  in the original coordinate system as follows

$$\begin{pmatrix} x' \\ y' \end{pmatrix} = \begin{pmatrix} \cos 45^\circ & \sin 45^\circ \\ -\sin 45^\circ & \cos 45^\circ \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} \quad (3)$$

Using the formula above, we can find the locations of the centers of the two squares,  $(2, 3)$  and  $(2, 1)$  respectively, shown in Figure 2 in the new  $x' - y'$

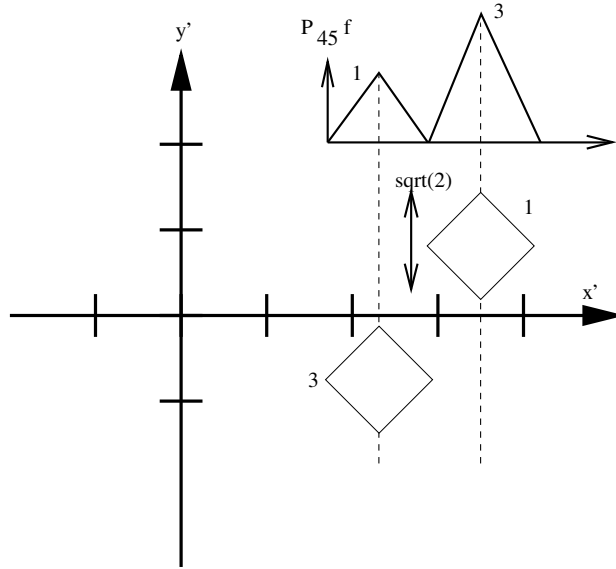


Figure 2:

coordinate system. We find that

$$\begin{aligned} (x, y) = (2, 3) &\rightarrow (x', y') = \left(\frac{5}{\sqrt{2}}, \frac{1}{\sqrt{2}}\right) \\ (x, y) = (2, 1) &\rightarrow (x', y') = \left(\frac{3}{\sqrt{2}}, -\frac{1}{\sqrt{2}}\right) \end{aligned} \quad (4)$$

By repeatedly applying the formula Equation 3 above to the  $(x, y)$  coordinates of the vertices of the two squares, we can deduce the shape of the objects in the  $x' - y'$  coordinate system. This is shown in Figure 2. We know that the two “diamonds” are centered at  $\left(\frac{5}{\sqrt{2}}, \frac{1}{\sqrt{2}}\right)$  and  $\left(\frac{3}{\sqrt{2}}, -\frac{1}{\sqrt{2}}\right)$  respectively. Hence, the two projected functions will resemble triangles centered on  $R = \frac{3}{\sqrt{2}}$  and  $R = \frac{5}{\sqrt{2}}$  respectively, where  $R = x'$ , see Figure 2.

Defining a triangle function as follows

$$\Lambda(x) = \begin{cases} 1 - |x| & |x| < 1 \\ 0 & \text{else} \end{cases} \quad (5)$$

we see that

$$P_{45}f(x, y) = 3\sqrt{2}\Lambda(\sqrt{2}(R - \frac{3}{\sqrt{2}})) + \sqrt{2}\Lambda(\sqrt{2}(R - \frac{5}{\sqrt{2}})) \quad (6)$$

**Problem 2:**[10 pts]

Figure 3 shows the result of backprojecting the 1-dimensional functions  $P_0 f(x, y)$  and  $P_{90} f(x, y)$ . The result of these two backprojection operations is the set of 4 objects that might comprise the desired function  $f(x, y)$ . In Figure 3, these 4 objects are denoted A,B,C and D respectively. We still need to apply backprojection to the remaining profile  $P_{45} f(x, y)$  which will determine which of these four points are consistent with all three projection profiles.

As in Problem 1, we will map the  $(x, y)$  coordinates of these 4 candidate objects to a rotated coordinate system  $(x', y')$ . Subsequently, we will see that the projection profile  $P_{45} f(x, y)$  can be obtained by integrating along the  $x'$ -direction. The four objects, A,B,C and D in the rotated  $x' - y'$  coordinate system are shown in Figure 4 The profile  $P_{45} f(x, y)$  is shown along with the 4 points. By inspection, we see that only points B,C and D can yield the 3 projection profiles  $P_0 f(x, y)$ ,  $P_{90} f(x, y)$  and  $P_{45} f(x, y)$ .

Now that we have determined the number and locations of the three objects, we need to determine their intensities. Let  $S_i$  be the intensities of the three points B,C and D. Then,

$$\begin{aligned} S_C &= 0 && \text{from } P_{45} f(x, y) \\ S_A + S_D &= 1.5 && \text{from } P_{45} f(x, y) \\ S_B &= 0.5 && \text{from } P_{45} f(x, y) \\ S_C + S_D &= 0.5 && \text{from } P_{90} f(x, y) \end{aligned} \quad (7)$$

Solving the above equations, we find that  $S_A = 1$ ,  $S_B = 0.5$ , and  $S_D = 0.5$ .

**Problem 3:**[10 pts]

We apply the three-step rule:

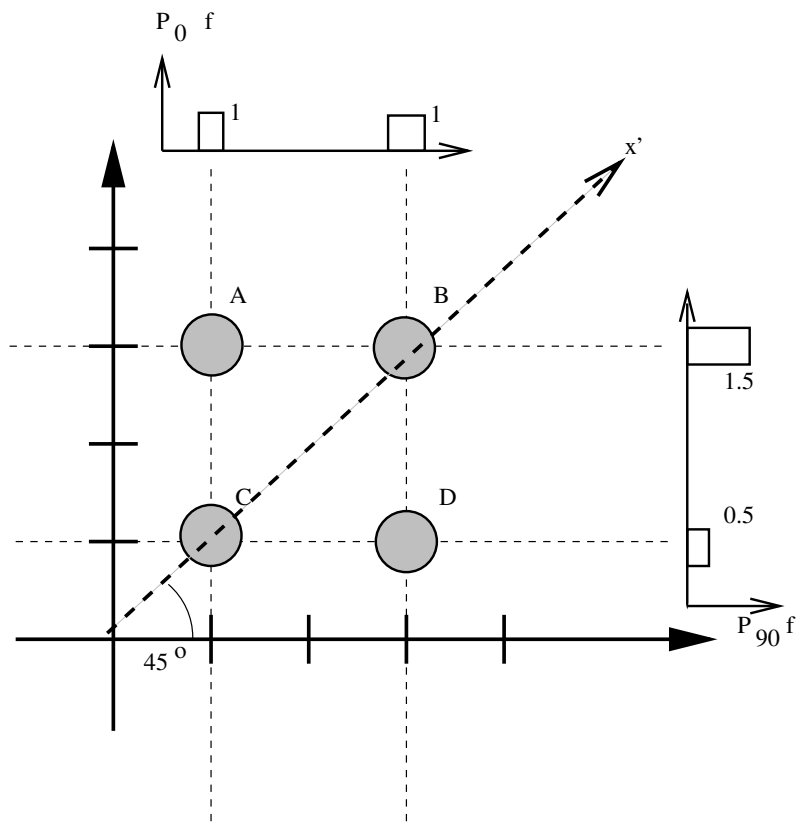


Figure 3:

- (i)  $\delta(f(x)) \rightarrow \frac{1}{\tau} \text{rect} \left( \frac{f(x)}{\tau} \right)$
- (ii) perform operation
- (iii) evaluate as  $\tau \rightarrow 0$

$$\delta(f(x)) = \lim_{\tau \rightarrow 0} \frac{1}{\tau} \text{rect} \left( \frac{f(x)}{\tau} \right) \quad (8)$$

where

$$\text{rect} \left( \frac{f(x)}{\tau} \right) = \begin{cases} 1 & |x| < \frac{\tau}{2} \\ 0 & \text{else} \end{cases} \quad (9)$$

Assuming  $\tau$  is small, then  $\text{rect} \left( \frac{f(x)}{\tau} \right)$  is non-zero only when the function  $f(x)$  lies in a band of width  $\tau$  about the  $x$ -axis. **We refer the reader to pages 3 and 4**

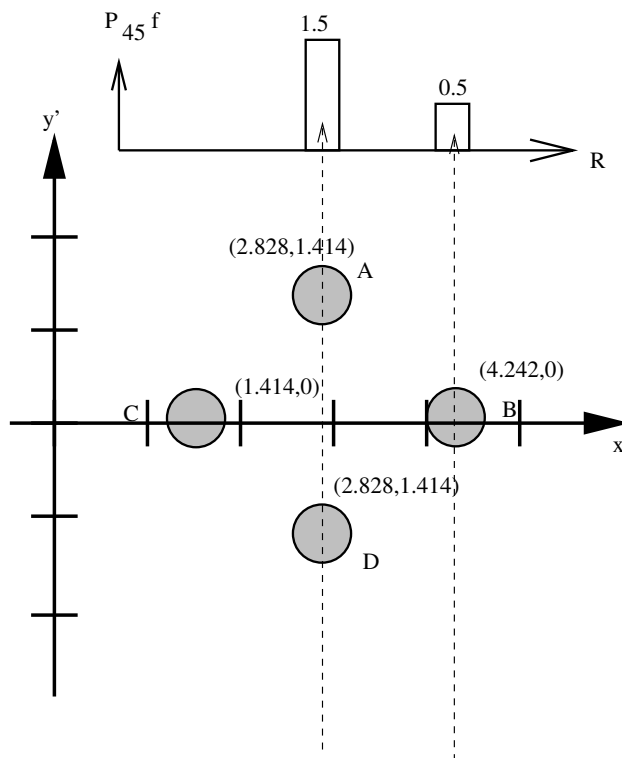


Figure 4:

**of Handout 12 for a graphical explanation.** Suppose, in particular, that in this function crosses the  $x$ -axis at  $x_0$ . In other words, let  $x_0$  be a root of the function  $f(x)$ .

$$f(x_0) = 0 \quad (10)$$

We can expand this function as a Taylor series about this root. Furthermore, we keep only the first order term in the expansion.

$$\begin{aligned} f(x) &= f(x_0) + f'(x_0)(x - x_0) + \frac{f''(x_0)}{2}(x - x_0)^2 \\ &\approx f(x_0) + f'(x_0)(x - x_0) \end{aligned} \quad (11)$$

Substituting the above expansion into the rect

$$\frac{1}{\tau} \text{rect} \left( \frac{f(x)}{\tau} \right) \approx \frac{1}{\tau} \text{rect} \left( \frac{f(x_0) + f'(x_0)(x - x_0)}{\tau} \right) \quad (12)$$

and making a change of variables  $\lambda = \frac{\tau}{f'(x_0)}$ , we have

$$\frac{1}{\lambda f'(x_0)} \text{rect} \left( \frac{(x - x_0)}{\lambda} \right) \quad (13)$$

Now, note that as  $\tau \rightarrow 0$ ,  $\lambda \rightarrow 0$ . So,

$$\frac{1}{\lambda f'(x_0)} \text{rect} \left( \frac{(x - x_0)}{\lambda} \right) \rightarrow \frac{1}{f'(x_0)} \delta(x - x_0) \text{ as } \lambda \rightarrow 0 \quad (14)$$

**Problem 4:**[10 pts]

(a) The line integral in question can be expressed as follows

$$I = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x, y) \delta(g(r)) dx dy \quad (15)$$

where,  $f(x, y) = \cos(\pi x)$  and  $g(r) = r - R$ . To evaluate this, and subsequent integrals, we make use of our result from Problem 3 which, in 2 dimensions, can be stated as

$$\delta(f(x, y)) = \frac{1}{|\nabla f|} \delta(x - x_0, y - y_0) \quad (16)$$

where  $x_0$  and  $y_0$  denote a locus of points in the  $x - y$  plane where the function  $f(x, y)$  is zero. As the argument of the delta function is in polar coordinates, we should compute the polar coordinate gradient magnitude in order

to find the strength of the impulse,  $|\nabla(r - R)| = \sqrt{\left(\frac{\partial(r-R)}{\partial r}\right)^2 + \left(\frac{1}{r} \frac{\partial(r-R)}{\partial \theta}\right)^2} =$

1. Thus, the impulse  $\delta(r - R)$  has unit strength.

$$\begin{aligned} I &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \cos(\pi x) \delta(r - R) dx dy & (17) \\ &= \int_0^{2\pi} \int_0^{\infty} \cos(\pi r \cos(\theta)) \delta(r - R) r dr d\theta \\ &= \int_0^{\infty} \left[ \int_0^{2\pi} \cos(\pi r \cos(\theta)) d\theta \right] \delta(r - R) r dr \end{aligned}$$

where the second line in the above results from a change of Cartesian to Polar coordinates. Now, the  $\theta$  integral, in square brackets above, is reminiscent of the definition of a Bessel function of first kind, order zero, defined as follows

$$\begin{aligned} J_0(x) &= \frac{1}{\pi} \int_0^\pi \cos(x \sin(\theta)) d\theta \\ &= \frac{1}{\pi} \int_{-\pi/2}^{\pi/2} \cos(x \cos(\theta')) d\theta', \quad \theta' = \theta - \frac{\pi}{2} \\ &= \frac{1}{2\pi} \int_{-\pi}^\pi \cos(x \cos(\theta'')) d\theta'', \quad \theta'' = 2\theta' \end{aligned} \quad (18)$$

Thus, substituting the Bessel function above in the line integral, we have

$$\begin{aligned} I &= \int_0^\infty 2\pi J_0(\pi r) \delta(r - R) r dr \\ &= 2\pi R J_0(\pi R) \end{aligned} \quad (19)$$

by using the sifting property of delta functions.

(b) In this problem, the line integral is as follows

$$I = \int_0^{2\pi} \int_0^\infty \cos(\pi r \cos(\theta)) \delta(r^2 - R^2) r dr d\theta \quad (20)$$

We first need to compute the strength of the impulse, using our result from problem 3. The gradient magnitude (in polar coordinates) of the argument of the delta-function in this part is

$$|\nabla(r^2 - R^2)| = 2r \quad (21)$$

which gives us the strength of the impulse as  $\frac{1}{2r}$ . Furthermore, the function  $r^2 - R^2 = 0$  only when  $r = R$ . Consequently, the delta function can be expressed as

$$\delta(r^2 - R^2) = \frac{1}{2R} \delta(r - R) \quad (22)$$

Now, performing the  $\theta$  integral first, plugging in the expression for the delta function above and recalling the definition of the Bessel function we used in part (a),

$$\begin{aligned} I &= \int_0^\infty 2\pi J_0(\pi r) \frac{1}{2R} \delta(r - R) r dr \\ &= \pi J_0(\pi R) \end{aligned} \quad (23)$$

by the sifting property of delta functions.

(c) In this problem, the line integral is as follows

$$I = \int_0^{2\pi} \int_0^\infty \cos(\pi r \cos(\theta)) \delta(\tan^{-1}(r - R)) r dr d\theta \quad (24)$$

The strength of the impulse function is proportional to the gradient magnitude of its argument

$$|\nabla(\tan^{-1}(r - R))| = \frac{1}{1 + (r - R)^2} \quad (25)$$

This gradient magnitude, a function of  $r$ , needs to be evaluated at values of  $r$  such that  $\tan^{-1}(r - R) = 0$ . Of course, the arctangent is zero only when  $r = R$ , which leads us to find that the strength of the delta function is unity. Thus, the line integral expression now becomes

$$\begin{aligned} I &= \int_0^{2\pi} \int_0^\infty \cos(\pi r \cos(\theta)) \delta(r - R) r dr d\theta \\ &= 2\pi R J_0(\pi R) \end{aligned} \quad (26)$$

(d) The strength per unit length of the impulse function  $\delta((r - R)^2)$  is computed from the magnitude of the gradient of the argument of the delta function,

$$|\nabla(r - R)^2| = 2(r - R) \quad (27)$$

This gradient magnitude is evaluated at the locus of points where  $(r - R)^2 = 0$ . Clearly, this locus is a circle of radius  $R$ . However, we find that, on this circle, the gradient magnitude is zero yielding an infinite strength per unit length. To understand this, we approximate the impulse function by a rect

$$\delta((r - R)^2) \approx \frac{1}{\tau} \text{rect}\left(\frac{(r - R)^2}{\tau}\right) \quad (28)$$

Note that for a given value of  $\tau$ ,

$$\text{rect}\left(\frac{(r - R)^2}{\tau}\right) = \begin{cases} 1 & (r - R)^2 < \tau/2 \\ 0 & \text{else} \end{cases} \quad (29)$$

This means that the width of the rect function is  $2\sqrt{\tau/2}$ . Consequently, the area of the rect function is  $\frac{1}{\tau} \times 2\sqrt{\tau/2} = \sqrt{\frac{2}{\tau}}$ . The area, therefore, is not

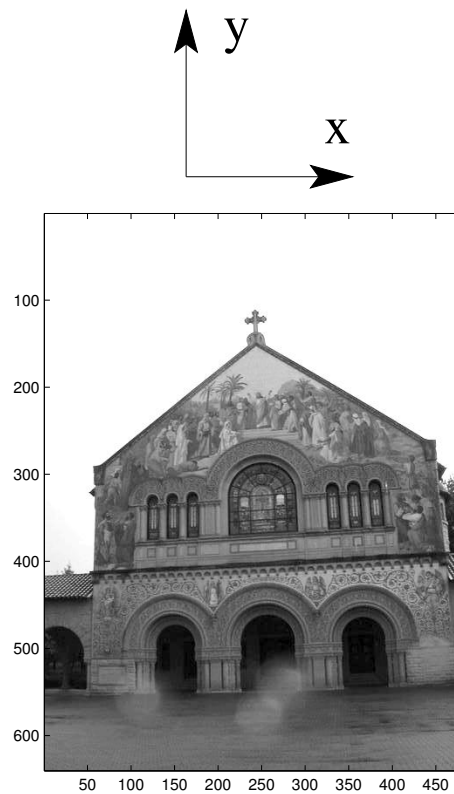


Figure 5:

independent of  $\tau$  and increases to infinity as  $\tau$  becomes infinitesimal. As the area the rect function is proportional to the strength of the delta, we see then why  $\delta((r - R)^2)$  has infinite strength.

**Problem 5:**[10 pts]

- (a) The image *hw3p5image* is shown in Figure 5. Figure 5 shows a grayscale image of Memorial Church. In addition, a  $y-x$  coordinate system is shown. We need to establish a mapping between row/column -  $(i, j)$  - pixel indices and  $(x, y)$  coordinates, under the assumption that the pixel spacings in both

dimensions are unity.

$$\begin{aligned}x &= j - 241 \\y &= 321 - i\end{aligned}\tag{30}$$

Thus, evaluating the line integral of the product of the image with the impulse function  $\delta(x)$  is simply the sum of all pixel values in column 241.

$$I = \sum_{i=1}^{640} \text{image}(i, 241) = 87164\tag{31}$$

- (b) Evaluating the line integral of the product of the image with the impulse function  $\delta(2y) = \frac{1}{2}\delta(y)$  is simply the sum of all pixel values in row 321, scaled by one half

$$I = \frac{1}{2} \sum_{j=1}^{480} \text{image}(321, j) = 28433\tag{32}$$

- (c) In this part, we sum all pixel values located on the line  $y - x = 0$  or  $(321 - i) - (j - 241) = 562 - i - j = 0$ .

$$I = \sum_k \text{image}(k, 562 - k) = 69628\tag{33}$$

We do not need to scale this sum,  $I$ , by the strength per unit length of the line impulse  $\delta(y - x)$  because, in this case, the strength per unit length is unity. Recall from the class notes, Handout 11, page 12, that the strength of a unit length impulse situated on a line  $y = mx + c$  is  $\cos \alpha$ , where  $\alpha$  is the angle that the line makes with the horizontal  $x$ -axis. So, in this problem, the strength per unit length of that line impulse would be  $\cos 45^\circ = \frac{1}{\sqrt{2}}$ . Figure 6 shows a unit length rect function of width  $\tau$  approximating the line impulse drawn on a grid with unit spacing. The grid represents the pixel points of the image which we wish to sample. We note that the minimum diagonal distance between two pixels is  $\sqrt{2}$ . Because of this, the strength per unit length of  $\frac{1}{\sqrt{2}}$  of our line impulse  $\delta(y - x)$  is multiplied by a distance of  $\sqrt{2}$  to give a strength per pixel of 1.

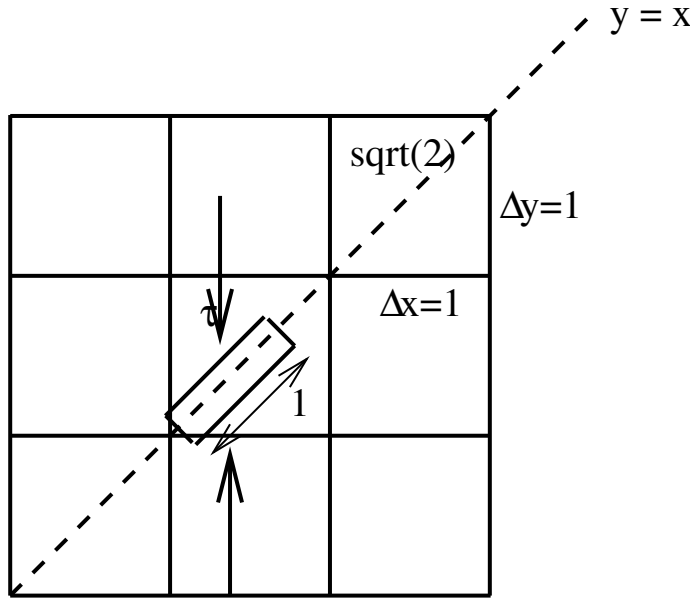


Figure 6:

(d) The line integrate we wish to evaluate is, in polar coordinates,

$$I = \int_0^{2\pi} \int_0^{\infty} f(r \cos \theta, r \sin \theta) \delta(r - 100) r dr d\theta \quad (34)$$

We approximate the above integral by a sum,

$$I \approx \sum_i \sum_{j=1}^{\infty} f(r_j \cos \theta_i, r_j \sin \theta_i) \delta(r_j - 100) r_j \Delta r \Delta \theta \quad (35)$$

where the  $i$  summation in the above is over the number of distinct angles that can be formed using the unit spacing  $x - y$  grid of the digital image. Of course,  $f(r_j \cos \theta_i, r_j \sin \theta_i)$  in the above refers to the  $ij$ -th pixel value of the image. We assume in the above that there exists a certain bin  $j$  for which  $r_j = 100$ . By virtue of the delta function above, we can then eliminate the inner,  $j$ , summation

$$I \approx \sum_{i=1}^N f(100 \cos \theta_i, 100 \sin \theta_i) 100 \Delta \theta \quad (36)$$

At this point, we assume that the length of the element of arc, a distance 100 pixels from the origin, is unity, i.e.  $100\Delta\theta = 1$ . This allows us to find  $N$ , the number of distinct angles, in the above. We mention here that different assumptions regarding arc length will yield different values of  $N$ , thereby affecting the final sum. In our case, we find that  $I = 75887$

Note: Using a more direct method of finding all pixels in the image that have distance 100 from the origin will yield a sum of 77723. This method is acceptable as well.

**Problem 6:**[10 pts]

- (a) To calculate the strength per unit area of the impulse  $\delta(r - \sqrt{x^2 + y^2 + z^2})$ , we apply the three-step rule

(i)  $\delta(f(r)) \rightarrow \frac{1}{\tau} \text{rect} \left( \frac{f(r)}{\tau} \right)$

(ii) perform operation

(iii) evaluate as  $\tau \rightarrow 0$

Now, in three dimensions,  $\frac{1}{\tau} \text{rect} \left( \frac{r - \sqrt{x^2 + y^2 + z^2}}{\tau} \right)$  can be visualized as a shell of radius  $R = \sqrt{x^2 + y^2 + z^2}$  and thickness  $\tau$ . An intensity of  $\frac{1}{\tau}$  is uniformly distributed at all points on this shell. The volume of this shell,  $V$ , can be computed as follows

$$V = \frac{4}{3}\pi(R + \frac{\tau}{2})^3 - \frac{4}{3}\pi(R - \frac{\tau}{2})^3 = \frac{4}{3}\pi(\frac{\tau^3}{4} + 3R^2\tau) \quad (37)$$

Finally, the total mass of this shell is  $V \times \frac{1}{\tau} = \frac{4}{3}\pi(\frac{\tau^2}{4} + 3R^2)$ . Applying the third rule, we take the limit as  $\tau$  goes to zero and find that the mass of a shell of thickness  $\tau$  is exactly  $4\pi R^2$ . But, this mass is distributed across a surface area of  $4\pi R^2$  as well. This means that the strength per unit area of the impulse  $\delta(r - R)$ , in three dimensions, is unity as it was in the 2D case.

The total mass (strength) integrated over all space is

$$\begin{aligned} M &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \delta(r - R) dx dy dz & (38) \\ &= \int_0^{\infty} \int_0^{2\pi} \int_0^{\pi} \delta(r - R) r^2 \sin \theta d\theta d\phi dr \\ &= 4\pi R^2 \end{aligned}$$

(b) The isotropic intensity pattern we are given is

$$f(r) = \frac{100 \text{ Watts}}{4\pi r^2}, \quad r = \sqrt{x^2 + y^2 + z^2} \quad (39)$$

We find the strength per unit area of the delta function  $\delta(r - R)$  by computing the magnitude of the gradient at  $r = R$

$$|\nabla (r - R)|_{r=R} = 1 \quad (40)$$

Thus, the strength per unit area of the delta function,  $S(r)$  is equal to unity.

Recording the total amount of power emitted by this light source at a certain distance  $R$  is equivalent to summing up the power densities on the surface of a shell of radius  $R$ , weighted by the strength per unit area  $S(r)$  of the delta function. Mathematically,

$$\begin{aligned} I &= \int_0^\infty \int_0^{2\pi} \int_0^\pi f(r)S(r)\delta(r - R)r^2 \sin \theta d\theta d\phi dr \quad (41) \\ &= \frac{100}{4\pi R^2} \int_0^{2\pi} d\theta \int_0^\pi R^2 \sin \phi d\phi \\ &= 100 \end{aligned}$$

Thus, we see that  $I = 100$  Watts, independent of the distance  $R$  where the observation is made. This makes physical sense because, by conservation of energy, the 100 Watts radiated by the light source should be maintained in a spherical shell of any radius  $R$ . Of course, the power *density* decreases as  $\frac{1}{R^2}$  with increasing distance from the location of the light source, as the surface area of that shell increases as  $4\pi R^2$ .

## MATLAB code for Problem 5

```
% load in image file, 480 rows by 640 cols

numrows = 480; numcols = 640;

fid = fopen('hw3p5image', 'rb');
data = fread(fid, inf, 'uint8');
```

```

fclose(fid);

im = reshape(data,[numrows numcols]); im = im';

% Note: x is in column direction, y is in row direction
Oy = 321; Ox = 241;

% create grid
[x,y] = meshgrid((1:numrows)-Ox, (numcols:-1:1)-Oy);

% part(a): delta(x)
%-----%

Ia = sum(im(:,Ox));

% part(b): delta(2y)
%-----%

Ib = sum(im(Oy,:))/2;

% part(c): delta(x-y)
%-----%
mask = (x==y); ind = find(mask == 1);
Ic = sum(im(ind));

% part(d): delta(r - 100)
%-----%

theta = 0:1/100:2*pi;
xx = round(100*cos(theta)); yy = round(100*sin(theta));
Id = sum(im(sub2ind(size(im),Oy-yy,xx+Ox)))

```