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Name: _____

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This exam is closed notes, closed book. The exam is worth a total of 116 points. The point value of each problem is indicated. Please show all work and clearly mark your answer.

Number	Points
1.	
2.	
3.	
4.	
5.	
6.	
7.	
8.	
9.	
10.	
Total	

1. (14 points) Let Ω be the upper half of the unit disk in \mathbb{R}^2 . That is, let

$$\Omega \equiv \{(x, y) \in \mathbb{R}^2 : x^2 + y^2 < 1, y > 0\}.$$

Use separation of variables to solve

$$\begin{cases} u_{xx} + u_{yy} = 0 & (x, y) \in \Omega \\ u(r, 0) = 0 = u(r, \pi) \\ u(1, \theta) = \theta(\theta - \pi). \end{cases}$$

You do **not** need to evaluate any integrals.

Answer: First, we write the equation in polar coordinates,

$$u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} = 0.$$

Now, using separation of variables, we look for a solution of the form $u(r, \theta) = R(r)\Theta(\theta)$. Plugging this into our equation, we have

$$R''\Theta + \frac{1}{r}R'\Theta + \frac{1}{r^2}R\Theta'' = 0.$$

Dividing by $R\Theta$ and multiplying by r^2 , we get

$$\frac{r^2R''}{R} + \frac{rR'}{R} = -\frac{\Theta''}{\Theta} = \lambda.$$

Our boundary condition

$$u(0, \theta) = 0 = u(r, \theta)$$

leads us to the eigenvalue problem

$$\begin{cases} -\Theta'' = \lambda\Theta & 0 < \theta < \pi \\ \Theta(0) = 0 = \Theta(\pi). \end{cases}$$

We know the solutions of this eigenvalue problem are

$$\Theta_n(\theta) = \sin(n\theta) \quad \lambda_n = n^2, n = 1, 2, \dots$$

Now we look at our equation for R_n , for $n = 1, 2, \dots$. We know

$$r^2R_n'' + rR_n' - n^2R_n = 0$$

has solutions

$$R_n = C_n r^n + D_n r^{-n} \quad n = 1, 2, \dots$$

As we do not want our solution to blow up as $r \rightarrow 0$, we discard the solutions r^{-n} . Therefore, our solutions for R_n are

$$R_n(r) = C_n r^n \quad n = 1, 2, \dots$$

Therefore, we let

$$\begin{aligned} u(r, \theta) &= \sum_{n=1}^{\infty} R_n(r) \Theta_n(\theta) \\ &= \sum_{n=1}^{\infty} C_n r^n \sin(n\theta). \end{aligned}$$

Our other boundary condition $u(1, \theta) = \theta(\theta - \pi)$ implies we want to find constants C_n such that

$$u(1, \theta) = \sum_{n=0}^{\infty} C_n \sin(n\theta) = \theta(\theta - \pi).$$

This is the Fourier sine series for our boundary data. We know our coefficients C_n must be given by

$$\begin{aligned} C_n &= \frac{\langle \theta(\theta - \pi), \sin(n\theta) \rangle}{\langle \sin(n\theta), \sin(n\theta) \rangle} \\ &= \frac{\int_0^\pi \theta(\theta - \pi) \sin(n\theta) d\theta}{\int_0^\pi \sin^2(n\theta) d\theta}. \end{aligned}$$

Therefore, our solution is given by

$$u(r, \theta) = \sum_{n=1}^{\infty} C_n r^n \sin(n\theta)$$

where

$$C_n = \frac{\int_0^\pi \theta(\theta - \pi) \sin(n\theta) d\theta}{\int_0^\pi \sin^2(n\theta) d\theta}.$$

2. (10 points) Let Ω be the upper half-plane in \mathbb{R}^2 . That is, let

$$\Omega = \{(x, y) \in \mathbb{R}^2, y > 0\}.$$

Consider the boundary-value problem,

$$\begin{cases} \Delta u = 0 & (x, y) \in \Omega \\ u = 0 & (x, y) \in \partial\Omega. \end{cases} \quad (1)$$

(a) Prove uniqueness of *bounded* solutions of (1). *Hint:* Suppose u is a solution of (1). Consider the odd reflection of u across the x -axis; that is, consider the function v defined as

$$v(x, y) = \begin{cases} u(x, y) & y > 0 \\ -u(x, -y) & y < 0. \end{cases}$$

Answer: Suppose u is a solution of (1). Then by defining v as the odd reflection of u , we note that v is a harmonic function in all of \mathbb{R}^2 . Also, by assumption, u is bounded. Therefore, v is bounded. By Liouville's Theorem, the only bounded, harmonic functions in all of \mathbb{R}^n are constant functions. Therefore, v must be constant. But, $v(x, 0) = 0$. Therefore, $v(x, y) \equiv 0$, which implies $u(x, y) \equiv 0$.

(b) Give an unbounded counterexample.

Answer: We note that $u(x, y) = Cy$ is a solution of (1) for any $C \in \mathbb{R}$. Therefore, for any $C_1 \neq C_2$ (neither of which is zero), $u_1(x, y) = C_1 y$ and $u_2(x, y) = C_2 y$ are two unbounded solutions of (1).

3. (18 points) Let Ω be the triangle with vertices at $(1, 0)$, $(-1, 0)$ and $(0, 2)$. Consider the eigenvalue problem

$$\begin{cases} -\Delta u = \lambda u & x \in \Omega \\ u = 0 & x \in \partial\Omega. \end{cases} \quad (2)$$

(a) Use the Comparison Principle to get an upper bound on the first eigenvalue for this eigenvalue problem. In particular, find the best upper bound on λ_1 among all rectangles contained within Ω with sides parallel to the coordinate axes.

Answer: For any rectangle R contained within Ω with sides parallel to the axes, and vertices $(x, 0)$, $(-x, 0)$, (x, y) , $(-x, y)$, the eigenvalues of R are given by

$$\lambda_{nm}(R) = \left(\frac{n\pi}{2x}\right)^2 + \left(\frac{m\pi}{y}\right)^2.$$

For any rectangle contained within Ω , we can get a better estimate if we extend the rectangle so that its vertices intersect the boundary. Therefore, we want to minimize the first eigenvalue for all rectangles contained within Ω such that $y = 2 - 2x$. That is, we want to minimize the function

$$f(x) = \frac{1}{(2x)^2} + \frac{1}{y(x)^2} = \frac{1}{4x^2} + \frac{1}{(2-2x)^2} = \frac{1}{4} \left[\frac{1}{x^2} + \frac{1}{(1-x)^2} \right].$$

For simplicity, we neglect the coefficient $1/4$. We just need to minimize

$$g(x) = \frac{1}{x^2} + \frac{1}{(1-x)^2}.$$

We look for critical points.

$$g'(x) = \frac{-2}{x^3} + \frac{2}{(1-x)^3}.$$

Now $g'(x) = 0$ implies

$$\frac{1}{(1-x)^3} = \frac{1}{x^3}$$

or

$$x^3 = (1-x)^3 \implies x = 1-x \implies x = \frac{1}{2}.$$

We see that $x = 1/2$ minimizes g . Therefore, the best upper bound on the first eigenvalue of Ω given by rectangles with sides parallel to the coordinate axes is the rectangle with vertices $(1/2, 0)$, $(-1/2, 0)$, $(1/2, 1)$, $(-1/2, 1)$. Denote this rectangle by R^* . Therefore,

$$\boxed{\lambda_1(R^*) = 2\pi^2 \geq \lambda_1(\Omega).}$$

(b) Let w_1, w_2 be two C^2 functions which are identically zero for $(x, y) \in \partial\Omega$. Explain how to use the Rayleigh-Ritz method with w_1 and w_2 to approximate the second eigenvalue of (2).

Answer: Define the 2×2 matrices

$$A = (\langle \nabla w_j, \nabla w_k \rangle)$$

$$B = (\langle w_j, w_k \rangle).$$

Consider the equation

$$\det(A - \lambda B) = 0.$$

The larger of the roots of this second-order equation gives an approximation for the second eigenvalue of (2).

(c) Find two linearly independent functions w_1 and w_2 which can be used for this approximation. (You do **not** need to apply the Rayleigh-Ritz approximation to these functions.)

Answer: We need to find two functions w_1 and w_2 , both of which vanish for $(x, y) \in \partial\Omega$. In particular, we want both functions to vanish along the lines $y = 0$, $y = -2x + 2$ and $y = 2x + 2$. Therefore, for example, we let

$$\begin{aligned} w_1(x, y) &= y(y - (-2x + 2))(y - (2x + 2)) \\ w_2(x, y) &= y^2(y - (-2x + 2))^2(y - (2x + 2))^2. \end{aligned}$$

4. (8 points) Let Ω be an open, bounded subset of \mathbb{R}^n . Consider

$$\begin{cases} \Delta u = 0 & x \in \Omega \\ \frac{\partial u}{\partial \nu} + u = g & x \in \partial\Omega. \end{cases} \quad (3)$$

(a) State the definition of a single-layer potential with moment h .

Answer:

$$\boxed{\bar{u}(x) = - \int_{\Omega} h(y) \Phi(y - x) dS(y).}$$

(b) In order to write the solution of (3) as a single-layer potential with moment h , what integral equation must h satisfy?

Answer: Fix $x_0 \in \partial\Omega$. Let $\nu(x_0)$ be the outer unit normal to $\partial\Omega$ at x_0 . For all $t < 0$ such that $x_0 + t\nu(x_0) \in \Omega$, let

$$i^{x_0}(t) = \nabla \bar{u}(x_0 + t\nu(x_0)) \cdot \nu(x_0).$$

We will say the boundary condition is satisfied if

$$\lim_{t \rightarrow 0^-} i^{x_0}(t) + \bar{u}(x) = g(x_0).$$

for all $x \in \Omega$, $x_0 \in \partial\Omega$. Now

$$\lim_{t \rightarrow 0^-} i^{x_0}(t) = -\frac{1}{2}h(x_0) - \int_{\partial\Omega} h(y) \frac{\partial \Phi}{\partial \nu_x}(x_0 - y) dS(y).$$

In addition, we recall that a single-layer potential is continuous for all $x \in \mathbb{R}^n$. Therefore, in order to find a solution as a single-layer potential, we need h to satisfy

$$\boxed{-\frac{1}{2}h(x_0) - \int_{\partial\Omega} h(y) \frac{\partial\Phi}{\partial\nu_x}(x_0 - y) dS(y) - \int_{\partial\Omega} h(y)\Phi(y - x_0) dS(y) = g(x_0)}$$

for all $x_0 \in \partial\Omega$.

5. (10 points) Recall that a Neumann function satisfies

$$\begin{cases} -\Delta_y N(x, y) = \delta_x & y \in \Omega \subset \mathbb{R}^n \\ \frac{\partial N}{\partial \nu}(x, y) = C & y \in \partial\Omega \end{cases}$$

for each $x \in \Omega$, where $C = \frac{1}{\int_{\partial\Omega} dS}$. Find the Neumann function for the first quadrant

$$\Omega = \{(x_1, x_2) \in \mathbb{R}^2 : x_1, x_2 > 0\}.$$

Answer: We note that $\int_{\partial\Omega} dS = 0$ for Ω the first quadrant. Therefore, our Neumann function must satisfy

$$\begin{cases} -\Delta_y N(x, y) = \delta_x & y \in \Omega \subset \mathbb{R}^n \\ \frac{\partial N}{\partial \nu}(x, y) = 0 & y \in \partial\Omega. \end{cases}$$

In particular, letting $\Phi(y)$ denote the fundamental solution of Laplace's equation in \mathbb{R}^2 , we can write $N(x, y) = \Phi(y - x) - \tilde{h}^x(y)$ if $h^x(y)$ is a solution of

$$\begin{cases} \Delta_y h^x(y) = 0 & y \in \Omega \\ \frac{\partial h^x(y)}{\partial \nu} = \frac{\partial \Phi(y-x)}{\partial \nu} y \in \partial\Omega. \end{cases}$$

We see that if $x = (x_1, x_2) \in \Omega$, then considering the reflected points $x^* = (x_1, -x_2)$, $\tilde{x} = (-x_1, x_2)$ and $\tilde{x}^* = (-x_1, -x_2)$, then

$$\boxed{N(x, y) = \Phi(y - x) + \Phi(y - x^*) + \Phi(y - \tilde{x}) + \Phi(y - \tilde{x}^*)}$$

satisfies the necessary conditions, where

$$\boxed{\Phi(y) = -\frac{1}{2\pi} \ln |y|}.$$

6. (8 points) Let $F_n : \mathcal{D} \rightarrow \mathbb{R}$ be the distribution defined such that

$$(F_n, \phi) = \int_{-\infty}^{\infty} \sin(nx) \phi(x) dx \quad \forall \phi \in \mathcal{D}.$$

Show that F_n converges to 0 weakly as $n \rightarrow +\infty$.

Answer: To show that F_n converges to 0 weakly as $n \rightarrow +\infty$, we need to show that $(F_n, \phi) \rightarrow 0$ for all $\phi \in \mathcal{D}$. Consider $\phi \in \mathcal{D}$. Therefore, ϕ has compact support. We note that

$$\begin{aligned} (F_n, \phi) &= \int_{-\infty}^{\infty} \sin(nx)\phi(x) dx \\ &= - \int_{-\infty}^{\infty} \frac{1}{n} \cos(nx)\phi'(x) dx + \frac{1}{n} \cos(nx)\phi'(x) \Big|_{-\infty}^{\infty}. \end{aligned}$$

Using the fact that ϕ has compact support, we know that the boundary terms go to zero. Therefore, we are left with

$$\begin{aligned} |(F_n, \phi)| &\leq \left| \int_{-\infty}^{\infty} \frac{1}{n} \cos(nx)\phi'(x) dx \right| \\ &\leq \frac{C}{n} \rightarrow 0 \text{ as } n \rightarrow +\infty. \end{aligned}$$

7. (8 points) Let Ω be an open, bounded subset of \mathbb{R}^n . Assume $a(x) > 0$ for all $x \in \partial\Omega$. Consider the eigenvalue problem

$$\begin{cases} -\Delta X = \lambda X & x \in \Omega \\ \frac{\partial X}{\partial \nu} + a(x)X = 0 & x \in \partial\Omega. \end{cases}$$

Prove that all eigenvalues are positive.

Answer: Suppose λ is an eigenvalue with eigenfunction X . Then

$$\begin{aligned} \lambda \int_{\Omega} X^2 dx &= - \int \Delta X X dx \\ &= + \int |\nabla X|^2 dx - \int_{\partial\Omega} \frac{\partial X}{\partial \nu} X dS(x) \\ &= \int_{\Omega} |\nabla X|^2 dx + \int_{\partial\Omega} a(x)X^2 dx \geq 0, \end{aligned}$$

using the assumption that $a(x) > 0$ for all $x \in \partial\Omega$. Therefore, we see that $\lambda \geq 0$.

It just remains to show that $\lambda \neq 0$. Suppose $\lambda = 0$. Then

$$\int_{\Omega} |\nabla X|^2 dx = 0 = \int_{\partial\Omega} a(x)X^2 dx.$$

The first equality implies that $X \equiv C$. The second equality implies $X \equiv 0$ for all $x \in \partial\Omega$. Therefore, $X \equiv 0$ for all $x \in \Omega$. But, the zero function is not an eigenfunction. Therefore, all eigenvalues are positive.

8. (10 points) Let Ω be an open, bounded subset of \mathbb{R}^n . Suppose $u \in C^2(\bar{\Omega})$ is a solution of

$$\Delta u = f \geq 0 \quad x \in \Omega.$$

Show that

$$u(x) \leq \int_{\partial B(x,r)} u(y) dS(y)$$

for all $B(x,r) \subset \Omega$.

Answer: Define the function

$$\phi(r) = \int_{\partial B(x,r)} u(y) dS(y)$$

for $r > 0$ and $\phi(0) = u(x)$. Using the assumption that u is continuous, we conclude that ϕ is continuous. We now look at $\phi'(r)$.

$$\begin{aligned} \phi'(r) &= \frac{d}{dr} \int_{\partial B(x,r)} u(y) dS(y) \\ &= \frac{d}{dr} \int_{\partial B(0,1)} u(x + rz) dS(z) \\ &= \int_{\partial B(0,1)} \nabla u(x + rz) \cdot z dS(z) \\ &= \int_{\partial B(x,r)} \nabla u(y) \cdot \frac{y - x}{r} dS(y) \\ &= \int_{\partial B(x,r)} \frac{\partial u}{\partial \nu} dS(y) \\ &= \frac{1}{n\alpha(n)r^{n-1}} \int_{\partial B(x,r)} \frac{\partial u}{\partial \nu} dS(y) \\ &= \frac{1}{n\alpha(n)r^{n-1}} \int_{B(x,r)} \Delta u dy \\ &\geq 0. \end{aligned}$$

Therefore, ϕ' is an increasing function of r . Therefore, we conclude that $\phi(0) \leq \phi(r)$. Therefore, we have

$$u(x) \leq \int_{\partial B(x,r)} u(y) dS(y),$$

as claimed.

9. (18 points) Determine whether the following statements are true or false. **Provide a reason for your answer.**

(a) Let h be a continuous function. The function

$$\bar{u}(x) = - \int_{\partial \Omega} h(y) \Phi(x - y) dS(y)$$

is harmonic for all $x \in \mathbb{R}^n$.

Answer: False. It is not harmonic for $x \in \partial \Omega$. (It is not even continuous for $x \in \partial \Omega$.)

(b) Let h be a continuous function. The function

$$\bar{u}(x) = - \int_{\partial\Omega} h(y) \frac{\partial\Phi}{\partial\nu_y}(x-y) dS(y)$$

is continuous for all $x \in \mathbb{R}^n$.

Answer: False. For example, let $h = 1$. By Gauss' Lemma, we know that

$$\bar{u}(x) = \begin{cases} 1 & x \in \Omega \\ \frac{1}{2} & x \in \partial\Omega \\ 0 & x \in \Omega^c \end{cases}$$

(c) Let h be a continuous function. Let $n \geq 2$. The function

$$\bar{u}(x) = - \int_{\partial\Omega} h(y) \frac{\partial\Phi}{\partial\nu_y}(x-y) dS(y)$$

is $O(|x|^{2-n})$.

Answer: True. We note that $\Phi(x-y) = \frac{C}{|x-y|^{n-2}}$. Therefore, $\frac{\partial\Phi}{\partial\nu_y}(x-y) = C|x-y|^{n-1} \cdot \nu(y) = O(|x|^{n-1}) = O(|x|^{n-2})$.

(d) Let Ω be an open, bounded subset of \mathbb{R}^n . Assume $a(x) \not\equiv 0$. There exists at most one solution of

$$\begin{cases} \Delta u = 0 & x \in \Omega \\ \frac{\partial u}{\partial\nu} - a(x)u = 0 & x \in \partial\Omega. \end{cases}$$

False. Suppose $\Omega = (0, 1)$. If u is harmonic in Ω , then $u(x) = A + Bx$. Now

$$\begin{aligned} \frac{\partial u}{\partial\nu}(0) - a(0)u(0) &= -u_x(0) - a(0)u(0) \\ \frac{\partial u}{\partial\nu}(1) - a(1)u(1) &= u_x(1) - a(1)u(1). \end{aligned}$$

Suppose $a(0) = -1$, $a(1) = 1/2$. Then our boundary conditions read

$$\begin{aligned} -u_x(0) + u(0) &= 0 \\ u_x(1) - \frac{1}{2}u(1) &= 0. \end{aligned}$$

The first boundary condition implies

$$-B + A = 0.$$

The second boundary condition implies

$$B - \frac{1}{2}(A + B) = 0.$$

We see that there are an infinite number of solutions to this system of equations, and, therefore, an infinite number of solutions to this boundary-value problem.

(e) Let $\phi : \mathbb{R}^n \rightarrow \mathbb{R}$ be a bounded function. There exists at most one bounded solution u of

$$\begin{cases} u_t - \Delta u = 0 & x \in \mathbb{R}^n, t > 0 \\ u(x, 0) = \phi(x). \end{cases}$$

Answer: True This follows from the uniqueness theorem for solutions of the heat equation on \mathbb{R}^n . In particular, we know there exists at most one solution $u(x, t)$ which satisfies the growth estimate

$$|u(x, t)| \leq Ae^{a|x|^2}$$

for some constants A, a . Therefore, in particular, there exists at most one bounded solution.

(f) If u is a harmonic function on the rectangle

$$\Omega \equiv \{(x, y) \in \mathbb{R}^2 : 0 < x < a, 0 < y < b\},$$

then

$$\int_0^a u_y(x, 0) dx + \int_0^b u_x(a, y) dy + \int_0^a u_y(x, b) dx + \int_0^b u_x(0, y) dy = 0.$$

Answer: False Let $u(x, y) = x$. This function is harmonic on Ω . Clearly, $u_y = 0$. But, $u_x(0, y) = 1 = u_x(a, y)$. Therefore,

$$\int_0^a u_y(x, 0) dx + \int_0^b u_x(a, y) dy + \int_0^a u_y(x, b) dx + \int_0^b u_x(0, y) dy = 2b \neq 0.$$

10. (12 points) Consider the initial/boundary value problem

$$\begin{cases} u_t - u_{xx} + u_x = 0 & 0 < x < l, t > 0 \\ u(x, 0) = \phi(x) & 0 < x < l \\ u_x(0, t) = 0 = u_x(l, t). \end{cases} \quad (4)$$

(a) Suppose u is a solution of this problem. Find a function g such that the function $v = gu$ satisfies

$$v_t - v_{xx} + v = 0.$$

Write the initial/boundary value problem that v satisfies.

Answer: If $v = gu$, then we see that

$$\begin{aligned} v_t - v_{xx} + v &= g_t u + g u_t - g_{xx} u - 2g_x u_x - g u_{xx} + g u \\ &= g(u_t - u_{xx} + u_x) - g u_x + g_t u - g_{xx} u - 2g_x u_x + g u \\ &= u(g_t - g_{xx} + g) - u_x(g + 2g_x). \end{aligned}$$

In order for this to equal zero, we need

$$g + 2g_x = 0$$

$$g_t - g_{xx} + g = 0.$$

The first equation implies $g(x, t) = C(t)e^{-\frac{1}{2}x}$. Plugging this into the second equation, we have

$$C''(t)e^{-\frac{1}{2}x} - \frac{1}{4}C(t)e^{-\frac{1}{2}x} + C(t)e^{-\frac{1}{2}x},$$

which implies

$$C'(t) + \frac{3}{4}C(t) = 0.$$

Therefore, $C(t) = e^{-\frac{3}{4}t}$. Therefore, we let

$$g(x, t) = e^{-\frac{3}{4}t}e^{-\frac{1}{2}x}.$$

Then for $v = gu$, using the fact that

$$\begin{aligned} v_x(0, t) &= g_x(0, t)u(0, t) + g(0, t)u_x(0, t) = -\frac{1}{2}e^{-\frac{3}{4}t}u(0, t) = -\frac{1}{2}v(0, t) \\ v_x(l, t) &= g_x(l, t)u(l, t) + g(l, t)u_x(l, t) = -\frac{1}{2}e^{-\frac{3}{4}t}e^{-\frac{1}{2}l}u(l, t) = -\frac{1}{2}v(l, t) \end{aligned}$$

we see that v is a solution of

$$\begin{cases} v_t - v_{xx} + v = 0 & 0 < x < l \\ v(x, 0) = g(x, 0)u(x, 0) = e^{-\frac{1}{2}x}\phi(x) \\ v_x(0, t) + \frac{1}{2}v(0, t) = 0 \\ v_x(l, t) + \frac{1}{2}v(l, t) = 0. \end{cases}$$

(b) In part (a) you show that $v = gu$ will satisfy an initial/boundary value problem of the form

$$\begin{cases} v_t - v_{xx} + v = 0 & 0 < x < l, t > 0 \\ v(x, 0) = f(x) & 0 < x < l \\ v \text{ satisfies symmetric B.C.s} & \end{cases} \quad (5)$$

Suppose λ_n, X_n are the eigenvalues and corresponding eigenfunctions of the eigenvalue problem

$$\begin{cases} -X'' = \lambda X & 0 < x < l \\ X \text{ satisfies (*)} & x = 0, l, \end{cases}$$

where (*) denotes the symmetric boundary conditions in (5). Write the solution of (5) in terms of X_n, λ_n and f .

Answer: We look for a solution of (5) by using separation of variables. Look for a solution of the form $v(x, t) = X(x)T(t)$. Plugging this into (5), we have

$$T'X - TX'' + TX = 0,$$

which implies

$$\frac{T'}{T} - \frac{X''}{X} + 1 = 0 \implies \frac{T'}{T} + 1 = \frac{X''}{X} = -\lambda.$$

Now the solutions of the eigenfunction problem

$$\begin{cases} -X'' = \lambda X & 0 < x < l \\ X \text{ satisfies (*)} & x = 0, l, \end{cases}$$

are assumed to be given by X_n, λ_n . Then we see the solutions of

$$\frac{T'_n}{T_n} + 1 = -\lambda_n$$

are given by $T_n(t) = C_n e^{-(\lambda_n+1)t}$. Therefore, our solution is given by

$$u(x, t) = \sum_{n=1}^{\infty} C_n X_n(x) e^{-(\lambda_n+1)t}$$

where

$$C_n = \frac{\langle f, X_n \rangle}{\langle X_n, X_n \rangle}.$$