

Math 220B - Summer 2003 Homework 4 Solutions

1. Let Ω be an open, bounded set in \mathbb{R}^n . Let $\phi \in C(\overline{\Omega})$, $g \in C(\partial\Omega)$, and suppose $u \in C^2(\overline{\Omega} \times [0, \infty))$ is a solution of

$$(1) \quad \begin{cases} u_t(x, t) = k\Delta u(x, t) & (x, t) \in \Omega \times (0, \infty) \\ u(x, t) = g(x) & (x, t) \in \partial\Omega \times [0, \infty) \\ u(x, 0) = \phi(x) & x \in \Omega. \end{cases}$$

Also suppose that $v \in C^2(\overline{\Omega})$ is a solution of

$$(2) \quad \begin{cases} \Delta v(x) = 0 & x \in \Omega \\ v(x) = g(x) & x \in \partial\Omega. \end{cases}$$

Prove $u(x, t) \rightarrow v(x)$ in $L^2(\Omega)$ as $t \rightarrow \infty$ as follows.

- (a) For Ω an open, bounded set in \mathbb{R}^n , there exists a constant C (depending only on Ω) such that for every $f \in C^1(\overline{\Omega})$ with $f(x) = 0$ for all $x \in \partial\Omega$,

$$(*) \quad \|f\|_{L^2(\Omega)} \leq C \|\nabla f\|_{L^2(\Omega)}.$$

Prove this inequality for the case when $\Omega = (a, b) \subset \mathbb{R}$.

Answer: Let

$$F(x) = \int_a^x f(y) dy.$$

Then integrating by parts and using the Cauchy-Schwartz inequality, we have

$$\begin{aligned} \|f\|_{L^2([a,b])}^2 &= \int_a^b f^2(x) dx \\ &= \left| - \int_a^b F(x) f'(x) dx + F(x) f(x) \Big|_{x=a}^{x=b} \right| \\ &\leq \int_a^b |F(x) f'(x)| dx \\ &\leq \left(\int_a^b |F(x)|^2 dx \right)^{1/2} \left(\int_a^b |f'(x)|^2 dx \right)^{1/2}. \end{aligned}$$

Now

$$\begin{aligned}
\left(\int_a^b F^2(x) dx\right)^{1/2} &= \left(\int_a^b \left(\int_a^x f(y) dy\right)^2 dx\right)^{1/2} \\
&\leq \left(\int_a^b \left(\int_a^b |f(y)| dy\right)^2 dx\right)^{1/2} \\
&\leq (b-a)^{1/2} \left(\int_a^b |f(y)| dy\right) \\
&\leq (b-a)^{1/2} \left(\int_a^b 1^2 dy\right)^{1/2} \left(\int_a^b |f(y)|^2 dy\right)^{1/2} \\
&\leq (b-a) \left(\int_a^b |f(y)|^2 dy\right)^{1/2}.
\end{aligned}$$

Therefore,

$$\int_a^b |f(x)|^2 dx \leq (b-a) \left(\int_a^b |f(x)|^2 dx\right)^{1/2} \left(\int_a^b |f'(x)|^2 dx\right)^{1/2}$$

Dividing both sides by $\left(\int_a^b |f(x)|^2 dx\right)^{1/2}$, we get the desired result.

- (b) Using (*), prove that for u and v solutions of (1) and (2) respectively, $u(x, t) \rightarrow v(x)$ in $L^2(\Omega)$ as $t \rightarrow +\infty$. (*Hint: Let $w(x, t) \equiv u(x, t) - v(x)$. Consider the PDE that w solves. Show that $\|w(x, t)\|_{L^2(\Omega)}^2 \rightarrow 0$ as $t \rightarrow +\infty$.)*)

Answer: Let $w(x, t) = u(x, t) - v(x)$. Then we see that w is a solution of

$$\begin{cases} w_t - k\Delta w = 0 & (x, t) \in \Omega \times (0, \infty) \\ w(x, 0) = \phi(x) - v(x) & x \in \Omega \\ w(x, t) = 0 & x \in \partial\Omega, t > 0 \end{cases}$$

We want to show that $\|w(x, t)\|_{L^2(\Omega)}^2 \rightarrow 0$ as $t \rightarrow \infty$. Now

$$\frac{1}{2} \|w(x, t)\|_{L^2(\Omega)}^2 = \frac{1}{2} \int_{\Omega} w(x, t)^2 dx$$

is the Energy function associated with the heat equation. Let

$$E(t) = \frac{1}{2} \int_{\Omega} w(x, t)^2 dx.$$

As we have seen,

$$\begin{aligned}
 E'(t) &= \int_{\Omega} w w_t \, dx \\
 &= \int_{\Omega} w k \Delta w \, dx \\
 &= -k \int_{\Omega} |\nabla w|^2 \, dx + k \int_{\partial\Omega} w \frac{\partial w}{\partial \nu} \, dS(x) \\
 &= -k \int_{\Omega} |\nabla w|^2 \, dx.
 \end{aligned}$$

But, using part (a), we see that

$$\begin{aligned}
 E'(t) &= -k \int_{\Omega} |\nabla w|^2 \, dx \\
 &\leq -C \int_{\Omega} |w|^2 \, dx \\
 &= -CE(t).
 \end{aligned}$$

Therefore, we see that

$$E(t) \leq K e^{-Ct}.$$

Therefore, $E(t) \rightarrow 0$ as $t \rightarrow +\infty$. In particular, this means $\|w(x, t)\|_{L^2(\Omega)}^2 \rightarrow 0$ as $t \rightarrow +\infty$, as claimed.

2. Let $B_n(0, a)$ be the unit ball in \mathbb{R}^n centered at 0 with radius $a > 0$.

(a) Let $\alpha > 0$. Show that

$$\int_{B_n(0, a)} \frac{1}{|x|^\alpha} \, dx < \infty$$

if and only if $\alpha < n$. In particular, evaluate the integral for $n > \alpha > 0$.

Answer: First, consider the case when $\alpha \neq n$. We have

$$\begin{aligned}
 \int_{B_n(0, a)} \frac{1}{|x|^\alpha} \, dx &= \int_0^a \int_{\partial B_n(0, r)} \frac{1}{|x|^\alpha} \, dS(x) \, dr = \int_0^a \int_{\partial B_n(0, r)} \frac{1}{r^\alpha} \, dS(x) \, dr \\
 &= \int_0^a \frac{1}{r^\alpha} \int_{\partial B_n(0, r)} dS(x) \, dr \\
 &= \int_0^a \frac{1}{r^\alpha} n\alpha(n) r^{n-1} \, dr \\
 &= \int_0^a n\alpha(n) r^{n-\alpha-1} \, dr \\
 &= \frac{n\alpha(n)}{n-\alpha} r^{n-\alpha} \Big|_{r \rightarrow 0}^{r=a}
 \end{aligned}$$

We see that this limit is finite if and only if $\alpha < n$. In particular, for $0 < \alpha < n$, we have

$$\boxed{\int_{B_n(0,a)} \frac{1}{|x|^\alpha} dx = \frac{n\alpha(n)}{n-\alpha} a^{n-\alpha}.}$$

Lastly, we consider the case when $\alpha = n$. In this case, we see that the integral becomes

$$\int_0^a n\alpha(n)r^{n-\alpha-1} dr = \int_0^a n\alpha(n)r^{-1} dr = n\alpha(n) \ln(r) \Big|_{r \rightarrow 0}^{r=a} \rightarrow +\infty.$$

Therefore, the integral is not finite for $\alpha = n$.

(b) Give conditions on α for which

$$\int_{\mathbb{R}^n \setminus B_n(0,a)} \frac{1}{|x|^\alpha} dx < \infty.$$

Answer: Again, consider $\alpha \neq n$. As in part (a), we write

$$\begin{aligned} \int_{\mathbb{R}^n \setminus B_n(0,a)} \frac{1}{|x|^\alpha} dx &= \int_a^\infty \int_{\partial B_n(0,r)} \frac{1}{|x|^\alpha} dS(x) dr \\ &= \int_a^\infty \frac{1}{r^\alpha} n\alpha(n)r^{n-1} dr \\ &= \int_a^\infty n\alpha(n)r^{n-\alpha-1} dr \\ &= \frac{n\alpha(n)}{n-\alpha} r^{n-\alpha} \Big|_{r=a}^{r \rightarrow +\infty}. \end{aligned}$$

We notice that this integral will only converge if $n - \alpha < 0$.

Now if $\alpha = n$, the integral becomes

$$\int_a^\infty n\alpha(n)r^{-1} dr = n\alpha(n) \ln r \Big|_{r=a}^{r \rightarrow +\infty} \rightarrow +\infty.$$

Therefore, we conclude that this integral is finite if and only if

$$\boxed{\alpha > n.}$$

3. Find all radial solutions of

$$-\Delta u + u = 0 \quad x \in \mathbb{R}^3 \setminus \{0\}.$$

Answer: Suppose there is a radial solution $u(x) = v(|x|)$. From class, we calculated that if $u(x) = v(|x|)$, then letting $r = |x|$,

$$\Delta u = \frac{n-1}{r} v'(r) + v''(r).$$

Therefore, in $n = 3$ dimensions, we have

$$\begin{aligned} -\Delta v + v &= -v'' - \frac{2}{r}v' + v = 0 \\ \implies -rv'' - 2v' + rv &= 0 \\ \implies -(rv)'' + rv &= 0. \end{aligned}$$

Let $w = rv$. Therefore, we see that w satisfies

$$w'' - w = 0$$

which implies

$$w(r) = Ae^r + Be^{-r}.$$

Solving for v in terms of w , we have

$$v(r) = \frac{w}{r} = \frac{Ae^r + Be^{-r}}{r}.$$

Therefore, we see that all radial solutions of our equation must be given by

$$u(x) = v(|x|) = \frac{Ae^{|x|} + Be^{-|x|}}{|x|}.$$

4. Prove that

$$u(x) \equiv \frac{e^{-|x|}}{4\pi|x|}$$

satisfies

$$-\Delta u + u = \delta_0 \quad x \in \mathbb{R}^3$$

in the sense of distributions.

Answer: Let $F_u : \mathcal{D}(\mathbb{R}^3) \rightarrow \mathbb{R}$ be defined such that

$$(F_u, \phi) = \int_{\mathbb{R}^3} u(x)\phi(x) dx \quad \forall \phi \in \mathcal{D}(\mathbb{R}^3).$$

To say that

$$-\Delta u + u = \delta_0$$

in the sense of distributions means that

$$(-\Delta F_u + F_u, \phi) = (\delta_0, \phi) = \phi(0).$$

Or, using the definition of derivative,

$$(F_u, -\Delta \phi) + (F_u, \phi) = \phi(0).$$

Therefore, we just need to show that

$$(F_u, -\Delta \phi) + (F_u, \phi) = \phi(0) \quad \forall \phi \in \mathcal{D}(\mathbb{R}^3).$$

We write

$$\begin{aligned}
(F_u, -\Delta\phi) &= - \int_{\mathbb{R}^3} u \Delta\phi \, dx \\
&= - \int_{B(0,\delta)} u \Delta\phi \, dx - \int_{K-B(0,\delta)} u \Delta\phi \, dx \\
&\equiv A + B
\end{aligned}$$

where $K = \text{supp}(\phi)$. In addition we write

$$\begin{aligned}
(F_u, \phi) &= \int_{B(0,\delta)} u \phi \, dx + \int_{K-B(0,\delta)} u \phi \, dx \\
&\equiv C + D.
\end{aligned}$$

We will show that $A + B + C + D \rightarrow \phi(0)$ as $\delta \rightarrow 0^+$, and, thus, the desired result follows.

First, we look at terms A and C . We see that

$$\begin{aligned}
|A + C| &= \left| - \int_{B(0,\delta)} u \Delta\phi \, dx + \int_{B(0,\delta)} u \phi \, dx \right| \\
&\leq (|\Delta\phi|_{L^\infty(B(0,\delta))} + |\phi|_{L^\infty(B(0,\delta))}) \int_{B(0,\delta)} |u| \, dx \\
&\leq C \int_{B(0,\delta)} \frac{e^{-|x|}}{|x|} \, dx \\
&\leq C \int_{B(0,\delta)} \frac{1}{|x|} \, dx.
\end{aligned}$$

From our answer to 2(a), we know that this integral is equal to

$$\frac{3\alpha(3)}{2} \delta^2 = O(\delta^2).$$

Therefore,

$$|A + C| \rightarrow 0 \quad \text{as } \delta \rightarrow 0^+.$$

We now need to look at terms B and D . Integrating B by parts, we have

$$- \int_{K-B(0,\delta)} u \Delta\phi \, dx = \int_{K-B(0,\delta)} \nabla u \cdot \nabla\phi \, dx - \int_{\partial B(0,\delta)} u \frac{\partial\phi}{\partial\tilde{\nu}} \, dS$$

where $\partial\phi/\partial\tilde{\nu} \equiv \nabla\phi \cdot \tilde{\nu}$ for $\tilde{\nu}$ the inner unit normal to $B(0, \delta)$. Above, we used the fact that $\phi \equiv 0$ on ∂K . We now need to look at the boundary term above. We have

$$\begin{aligned}
\left| \int_{\partial B(0,\delta)} u \frac{\partial\phi}{\partial\tilde{\nu}} \, dS \right| &\leq C \left| \frac{\partial\phi}{\partial\tilde{\nu}} \right|_{L^\infty(\partial B(0,\delta))} \int_{\partial B(0,\delta)} \frac{e^{-|x|}}{|x|} \, dS \\
&\leq \frac{C}{|x|} \int_{\partial B(0,\delta)} dS \\
&= \frac{C}{\delta} \delta^{n-1} \\
&= C\delta \rightarrow 0 \text{ as } \delta \rightarrow 0^+.
\end{aligned}$$

Therefore, we see that

$$-\int_{K-B(0,\delta)} u \Delta \phi \, dx = \int_{K-B(0,\delta)} \nabla u \cdot \nabla \phi \, dx + O(\delta).$$

Integrating by parts again, we have

$$\begin{aligned} -\int_{K-B(0,\delta)} u \Delta \phi \, dx &= -\int_{K-B(0,\delta)} \Delta u \phi \, dx + \int_{\partial B(0,\delta)} \frac{\partial u}{\partial \tilde{\nu}} \phi \, dS + O(\delta) \\ &\equiv B(1) + B(2). \end{aligned}$$

Now $B(1) + D = 0$ by direct calculation, or using the result of problem 3. In particular, we know that $-\Delta(e^{-|x|}/|x|) + e^{-|x|}/|x| = 0$ for $x \neq 0$. Finally, it remains to look at $B(2)$. By direct calculation we see that

$$\nabla u = -\frac{1}{4\pi} \left[\frac{e^{-|x|x}}{|x|^2} + \frac{e^{-|x|x}}{|x|^3} \right].$$

Further, we know that the inner unit normal on $\partial B(0, \delta)$ is given by

$$\tilde{\nu} = \frac{-x}{|x|}.$$

Therefore, we see that

$$\frac{\partial u}{\partial \tilde{\nu}} = \frac{1}{4\pi} \left[\frac{e^{-|x|}}{|x|} + \frac{e^{-|x|}}{|x|^2} \right].$$

Therefore, we have two terms to look at. First, we see that

$$\begin{aligned} \left| \frac{1}{4\pi} \int_{\partial B(0,\delta)} \frac{e^{-|x|}}{|x|} \phi \, dS \right| &= \left| \frac{e^{-\delta}}{4\pi\delta} \int_{\partial B(0,\delta)} \phi \, dS \right| \\ &\leq \frac{C}{\delta} \int_{\partial B(0,\delta)} dS \\ &= \frac{C}{\delta} \delta^2 = C\delta \rightarrow 0 \text{ as } \delta \rightarrow 0^+. \end{aligned}$$

Now for the second term we have

$$\begin{aligned} \frac{1}{4\pi} \int_{\partial B(0,\delta)} \frac{e^{-|x|}}{|x|^2} \phi \, dS &= \frac{e^{-\delta}}{4\pi\delta^2} \int_{\partial B(0,\delta)} \phi \, dS \\ &= e^{-\delta} \int_{\partial B(0,\delta)} \phi \, dS \rightarrow e^{-0} \phi(0) = \phi(0) \text{ as } \delta \rightarrow 0^+, \end{aligned}$$

using the fact that ϕ is continuous. Combining all of the above estimates, we conclude that

$$A + B + C + D \rightarrow \phi(0) \quad \text{as } \delta \rightarrow 0^+.$$