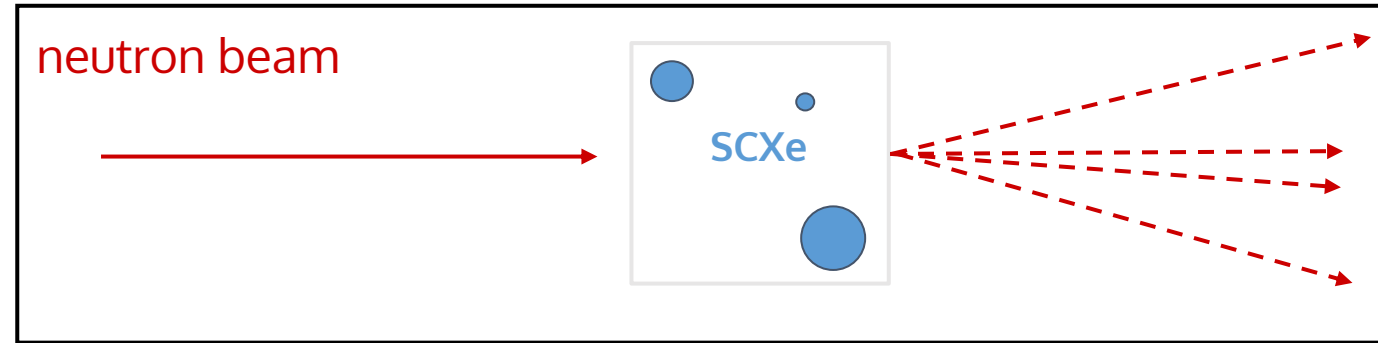
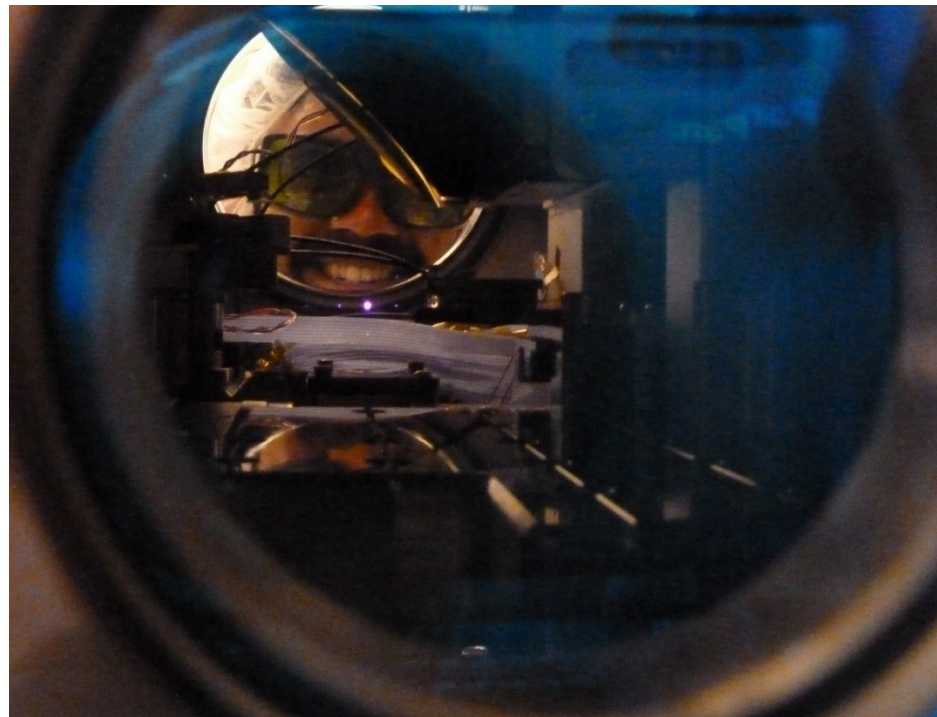
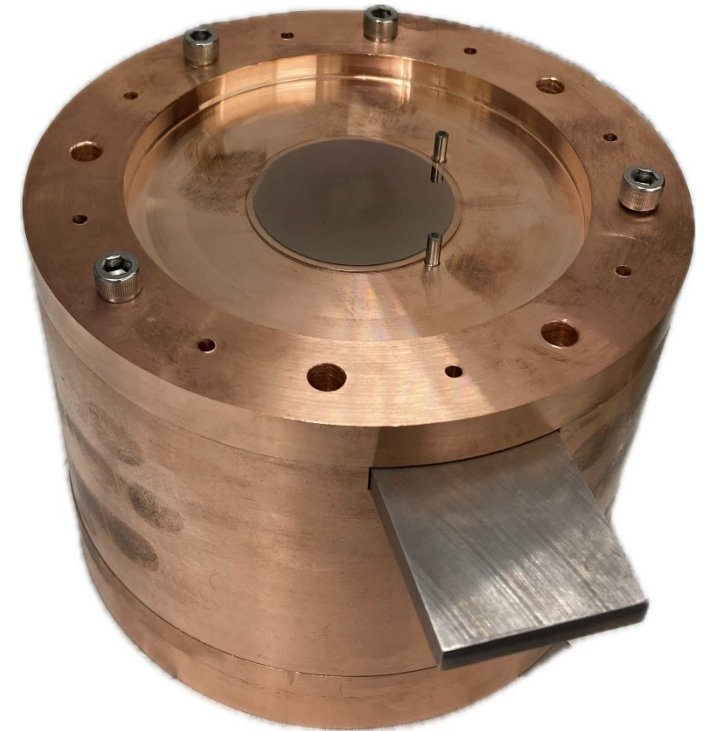


# Testing Gravity at Ever Shorter Scale

*a trip into exotic experimental physics*



*Giorgio Gratta*  
*Physics Dept, Stanford*



***“My advice is to try crazy ideas and innovative experiments.”***

*Steven Weinberg, APS News Feb 2019*

The inverse square law is generally assumed to work all the way down to the Planck

$$\text{length } R_P = \sqrt{\frac{G\hbar}{c^3}} = 1.6 \times 10^{-35} \text{ m}$$

Of course, this is a bold assumption that requires experimental verification.

In addition, there are important theoretical reasons to suspect that deviation from  $1/R^2$  may actually arise naturally and be more than just plausible.

Gravity is a notoriously rebellious interaction. Its quantum field theory is not as well established as that of other fundamental interactions.

And gravity at ordinary energies/distances is so much weaker than any of the other fundamental interactions.

Why? Are those issues related to each other?

Will the solutions of these puzzles simultaneously solve other modern puzzles in physics, such as those of Dark Matter or Dark Energy.

*Anyway, one can take the point of view that exploring the law of gravity at any distance is such an important endeavor that should be carried out irrespective of theoretical prejudice.*

**So, how well do we know that the inverse square law applies?**

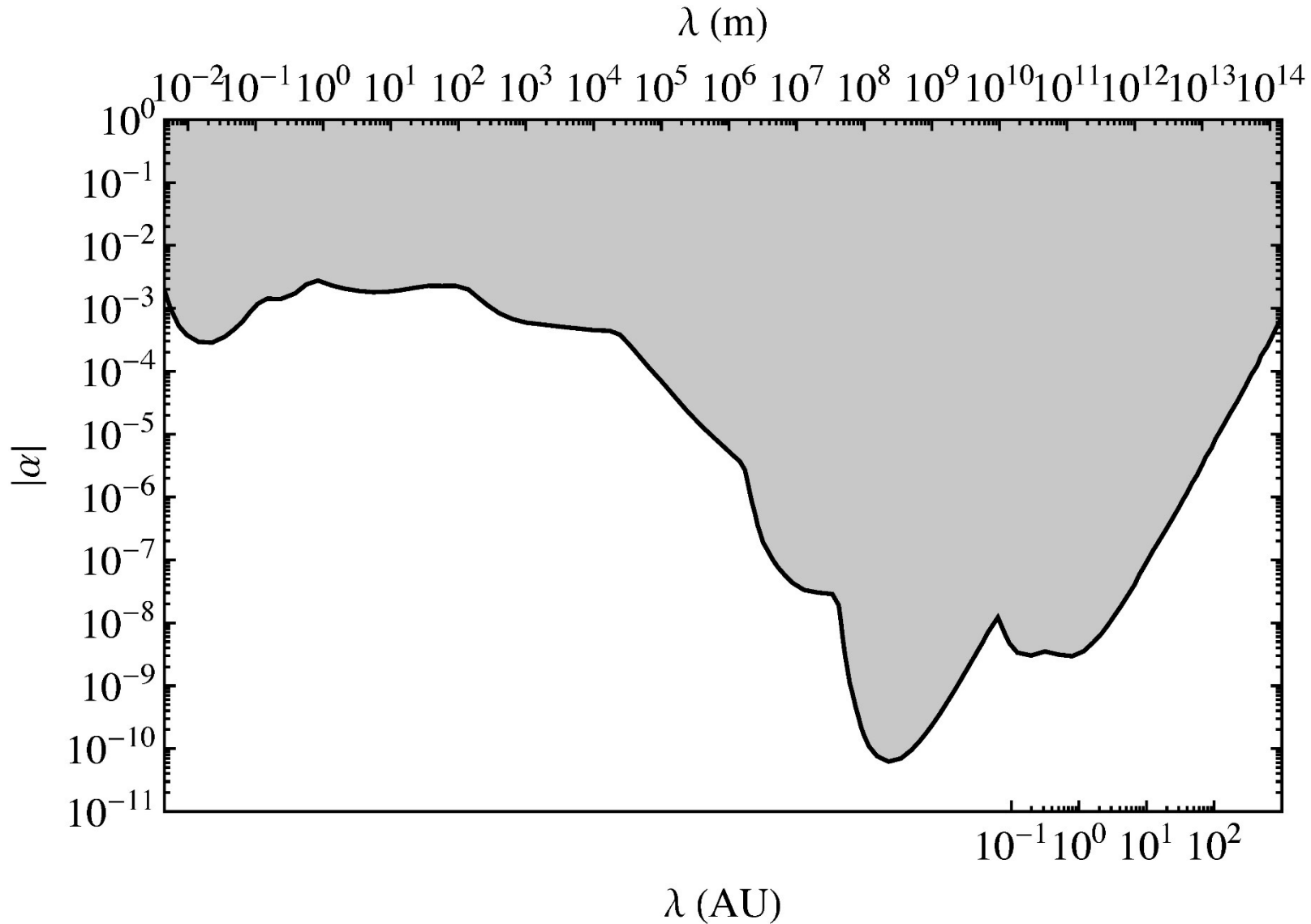
It is customary to express potential deviations from the  $1/R^2$  law by modifying the potential with a Yukawa term, obtaining:

$$V(R) = G \frac{M_1 M_2}{R} (1 + \alpha e^{-R/\lambda})$$

$\alpha$ : magnitude of the effect

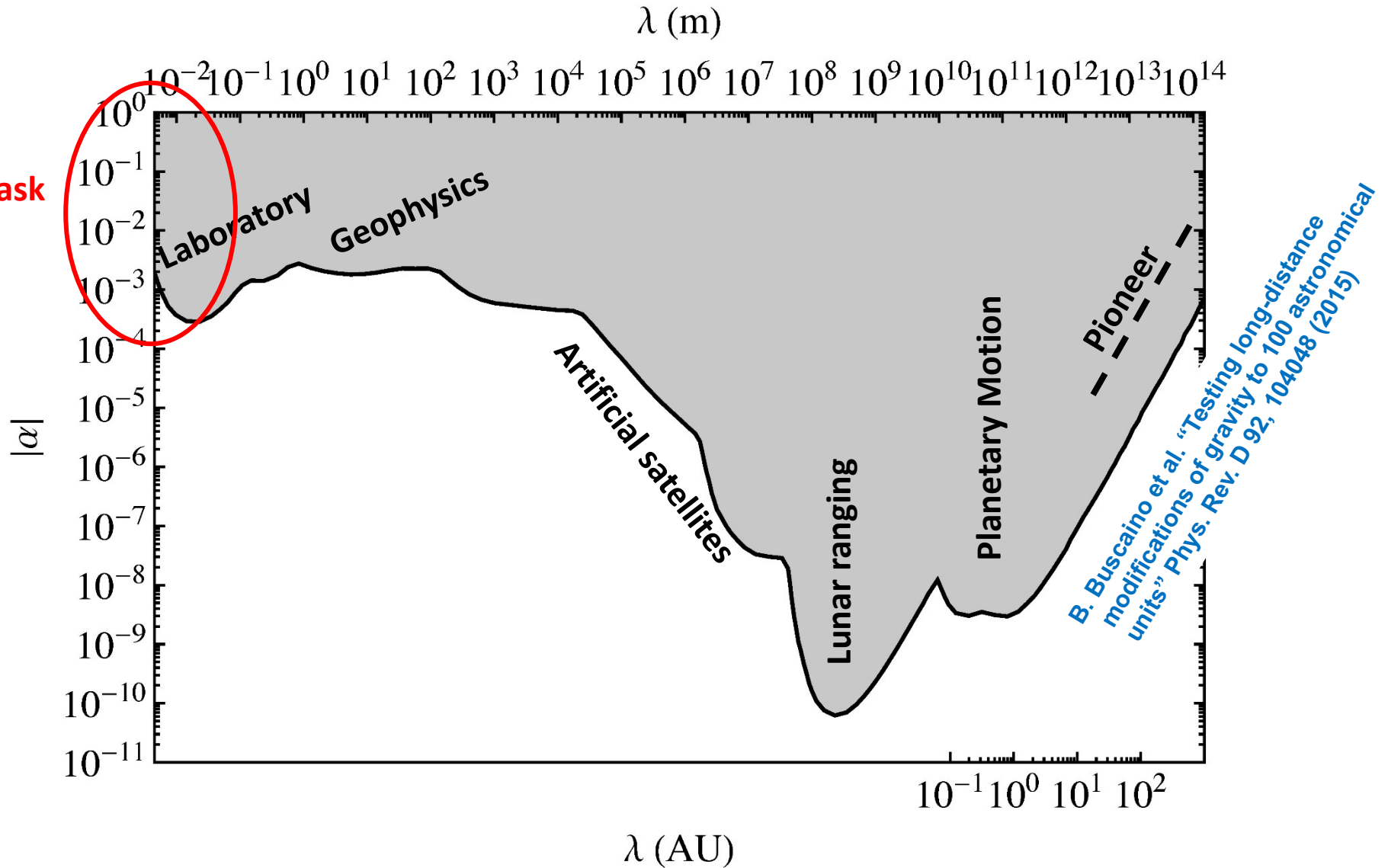
$\lambda$ : scale of the effect

# What do we empirically know



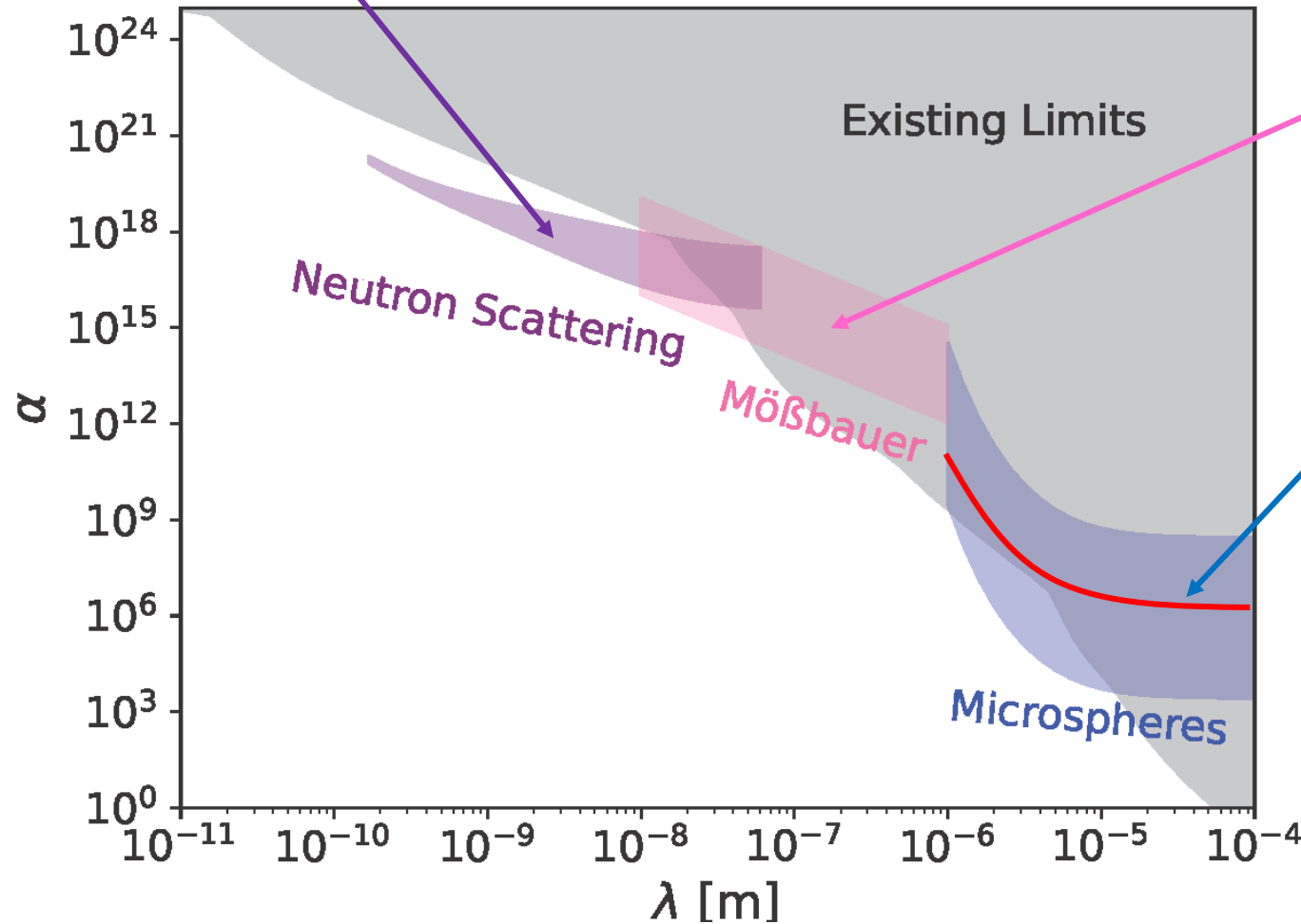
# What do we empirically know

Today's primary task



# Aspirations of our program (the bands are projections)

Z.Bogorad, P.W.Graham, GG,  
*Phys Rev D 108 (2023) 055005*



Also a motivation to “nuclear quantum optics”, with various potential applications.

GG, D.E. Kaplan, S. Rajendran,  
*Phys Rev D 102 (2020) 115031.*

Recent result:

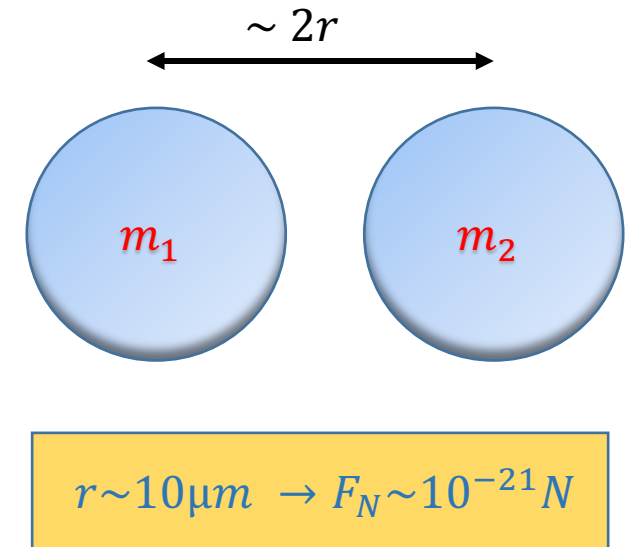
G.Venugopalan et al., *Nature Sci Rep 16 (2026) 5180*

Also, application to neutrality of matter, inertial sensing, quantum S&T

It is also desirable to have some overlap between techniques, convenient in case of a discovery.

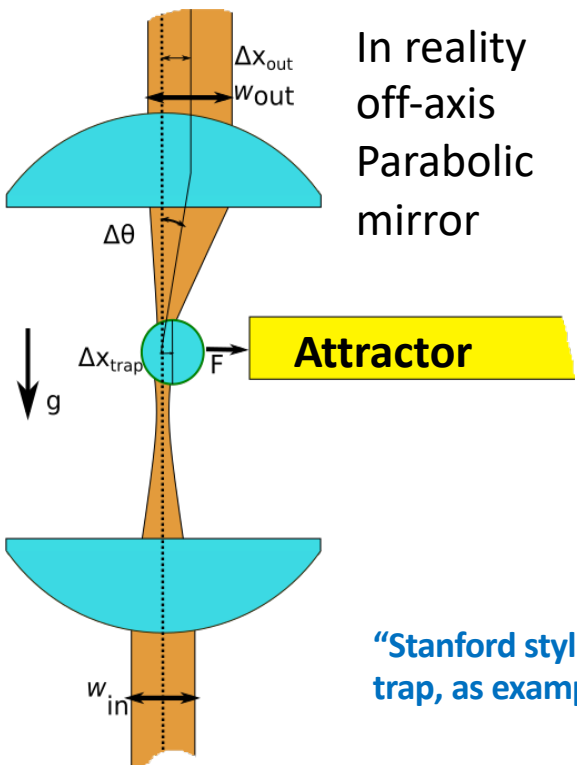
# Experimental challenges

- Since  $F = G \frac{M_1 M_2}{R^2} = G \frac{\rho_1 V_1 \rho_2 V_2}{R^2}$   
for atomic materials (we can't use Neutron Stars!)  
 $\rho_1 \sim \rho_2 < 20 \text{ g/cm}^3$ , there is no silver bullet.  
In addition, the volume  $V \sim R^3$ , so  $F \sim G \frac{\rho^2 R^6}{R^2}$   
and it is clear that measurements at short  
distance become exceedingly difficult.
- At distances  $< 100 \mu\text{m}$  even neutral matter results  
in residual E&M interaction that are a dangerous  
background for these measurements



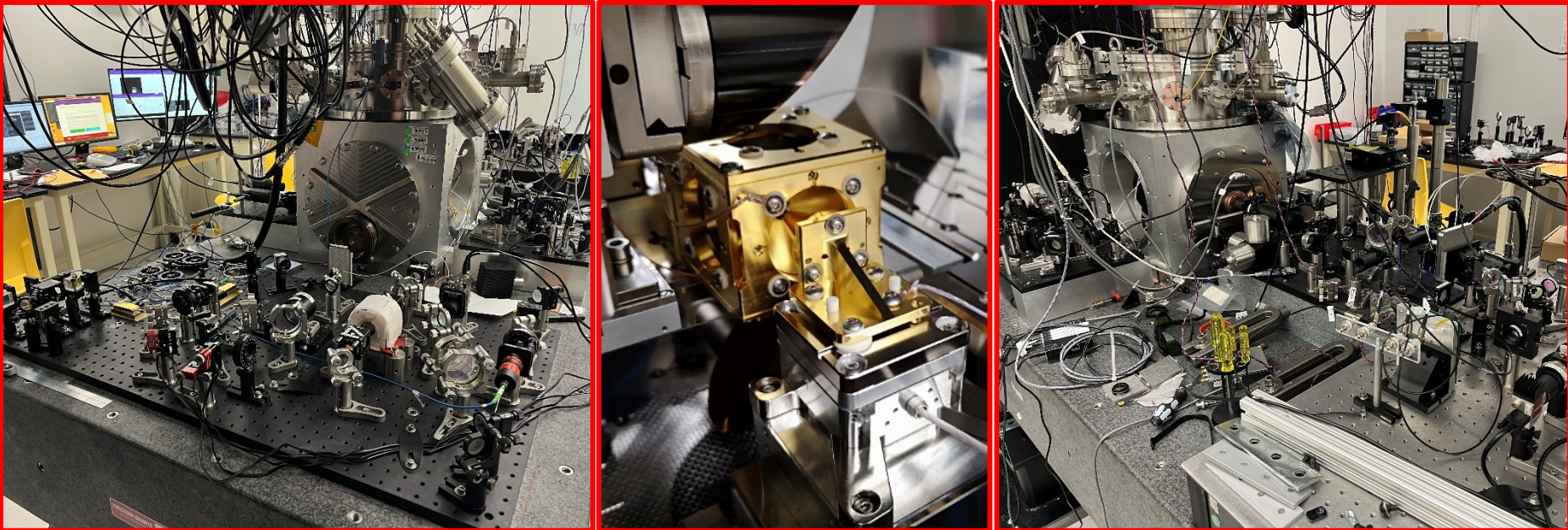
In the last 50 years, optical tweezers have developed into broad field, primarily applied to biology, in water.

Microspheres optically trapped in vacuum make superb force sensors. No mechanical springs!



In reality off-axis Parabolic mirror

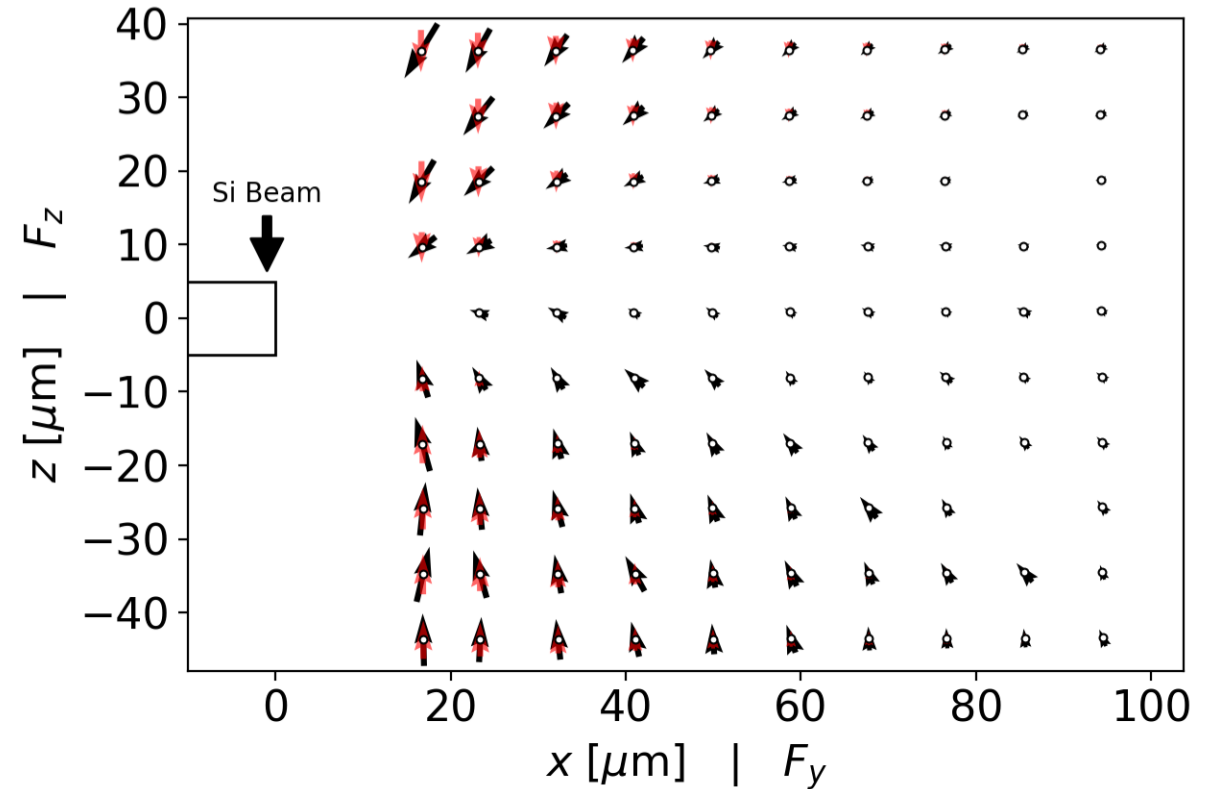
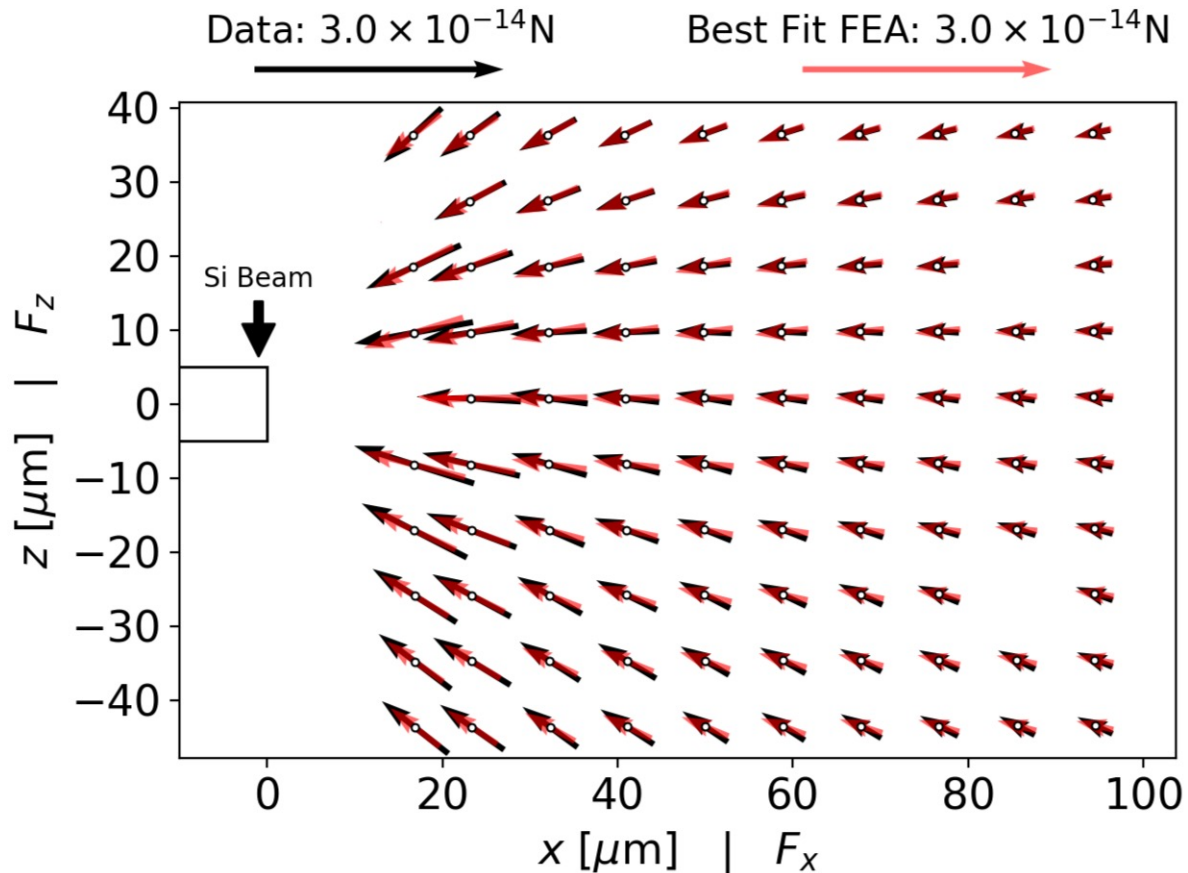
“Stanford style” trap, as example



Rev. Sci. Instrum. 91, 083201 (2020)

# The trapped microsphere (5-10 $\mu\text{m}$ diameter) is an excellent force sensor, with full 3D, vector field mapping capability

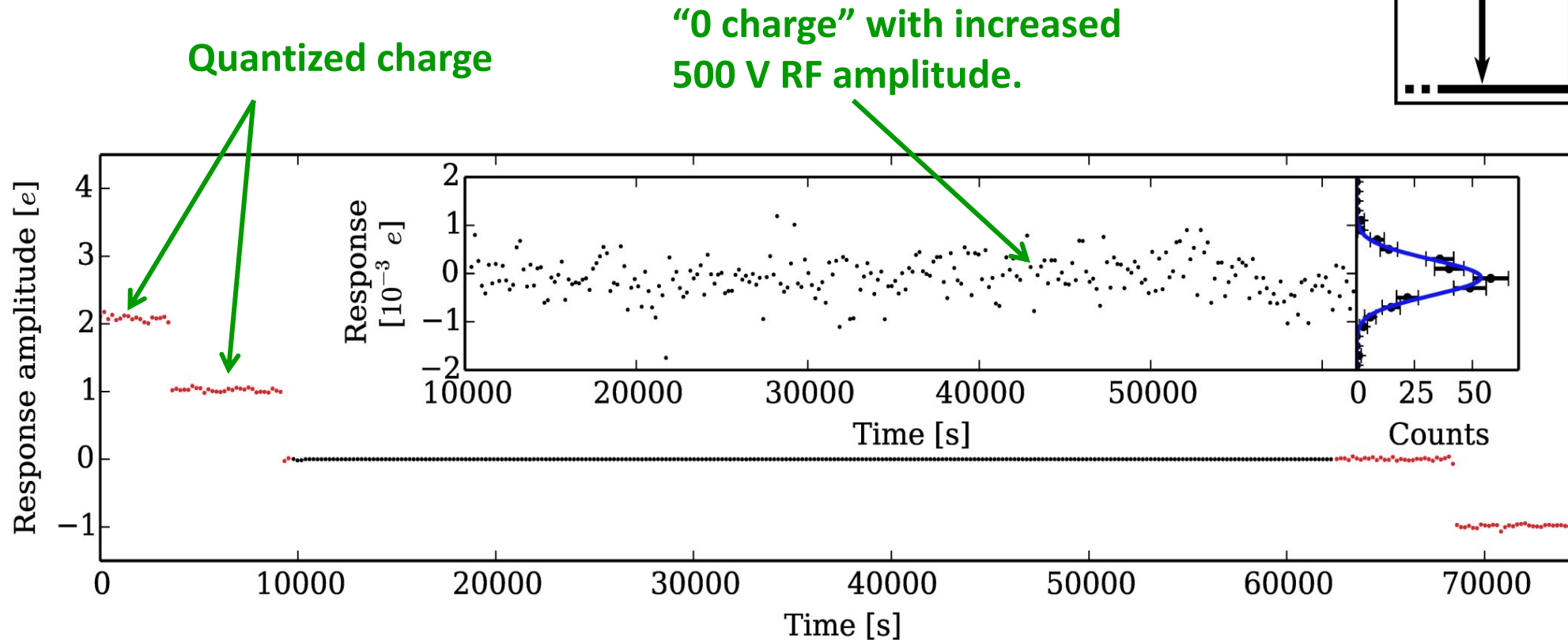
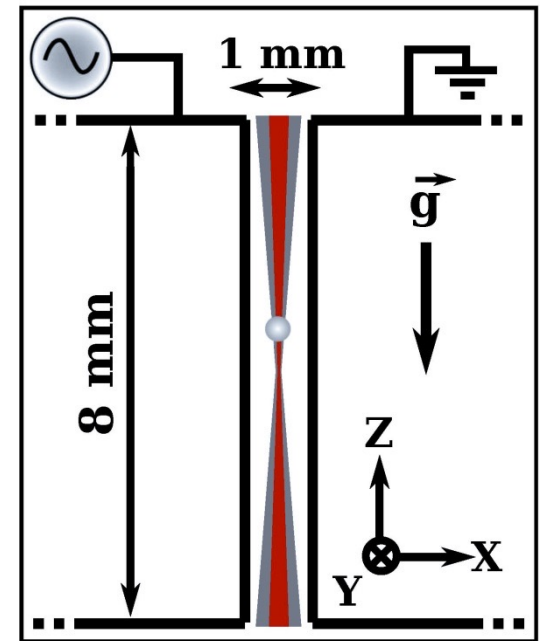
*as can be shown/calibrated by measuring charged microspheres*

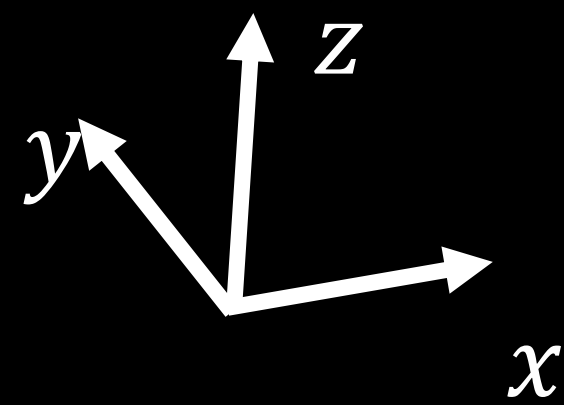
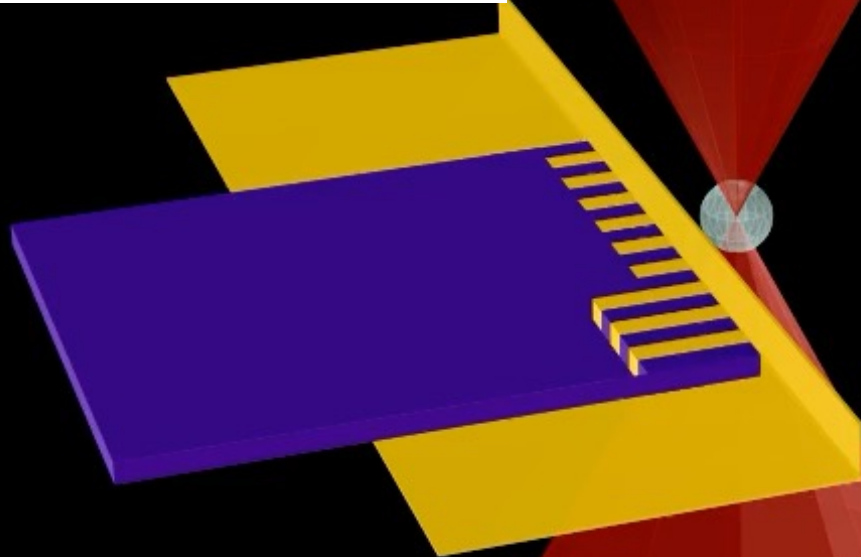
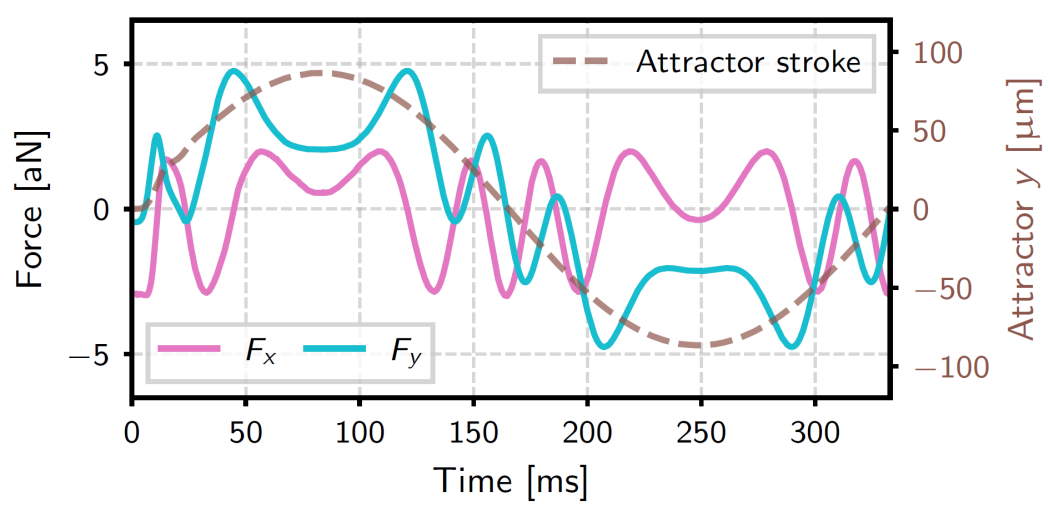


C.Blakemore et al., Phys. Rev. A 99 (2019) 023816

Similar results in G.Winstone et al., Phys. Rev. A 98 (2018) 053831

The possibility of setting the charge of the microsphere at leisure, opens the opportunity to search for very small charges ( $\ll 1$  milli- $e^-$ ) or, more interestingly, search for deviations from neutrality of matter.



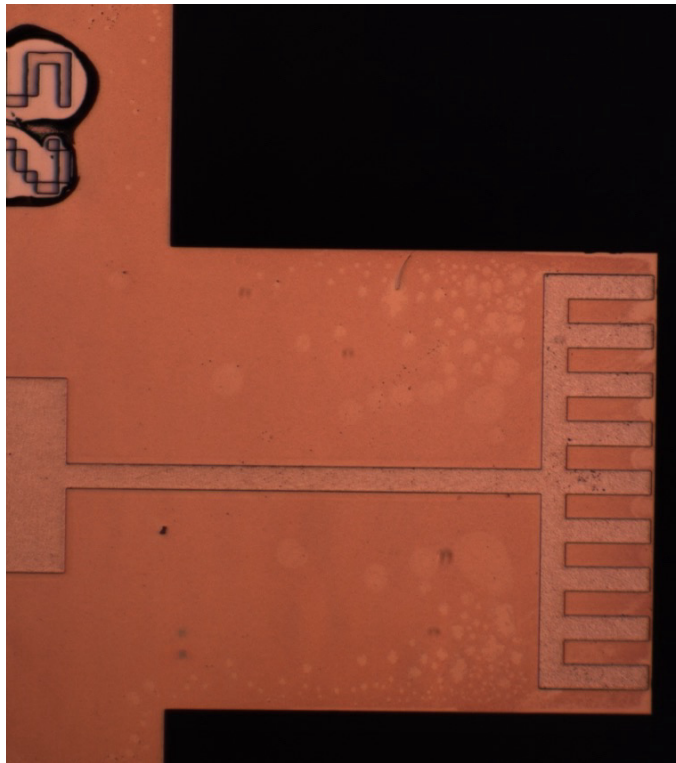


\*Microsphere motion greatly exaggerated

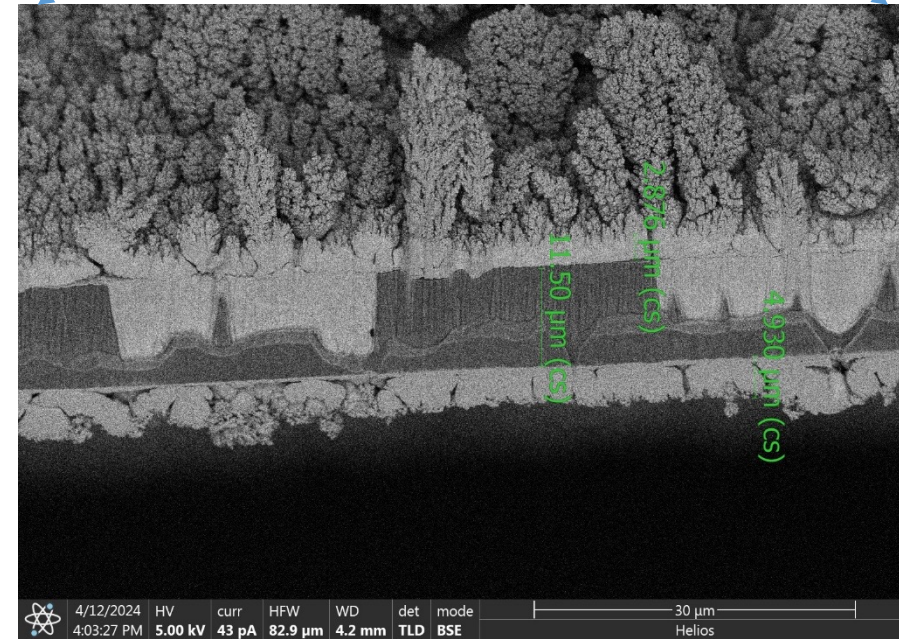
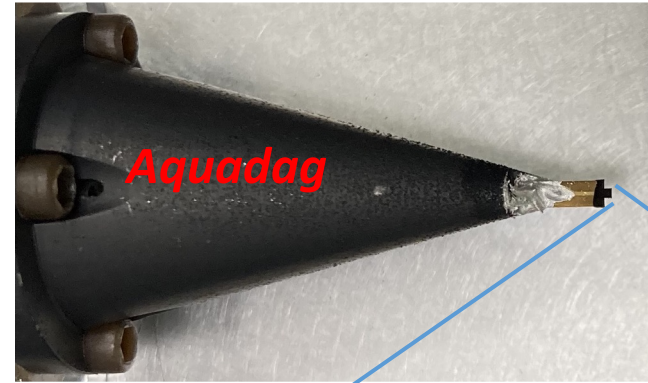
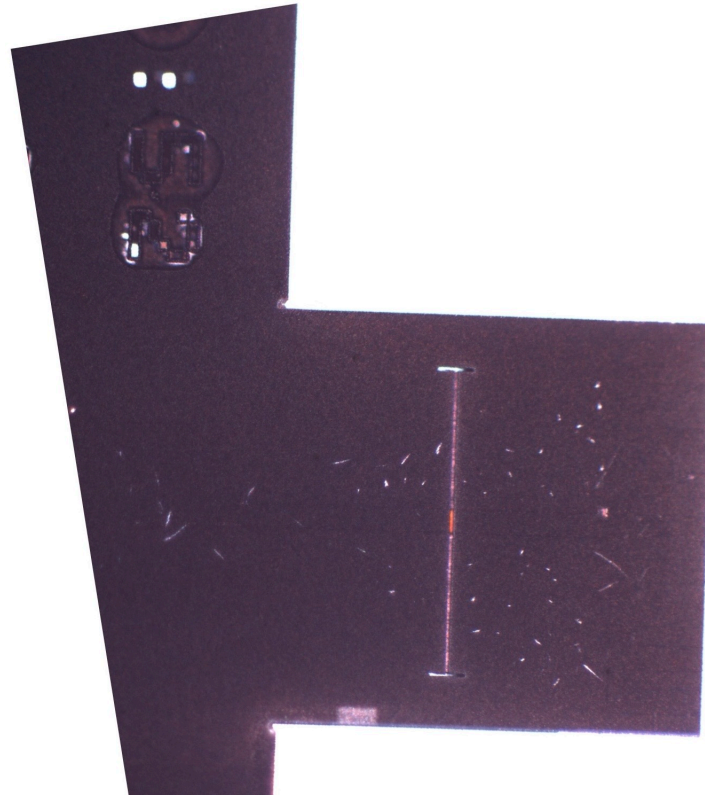
Animation by G. Venugopalan

# Electroplating "Platinum Black"

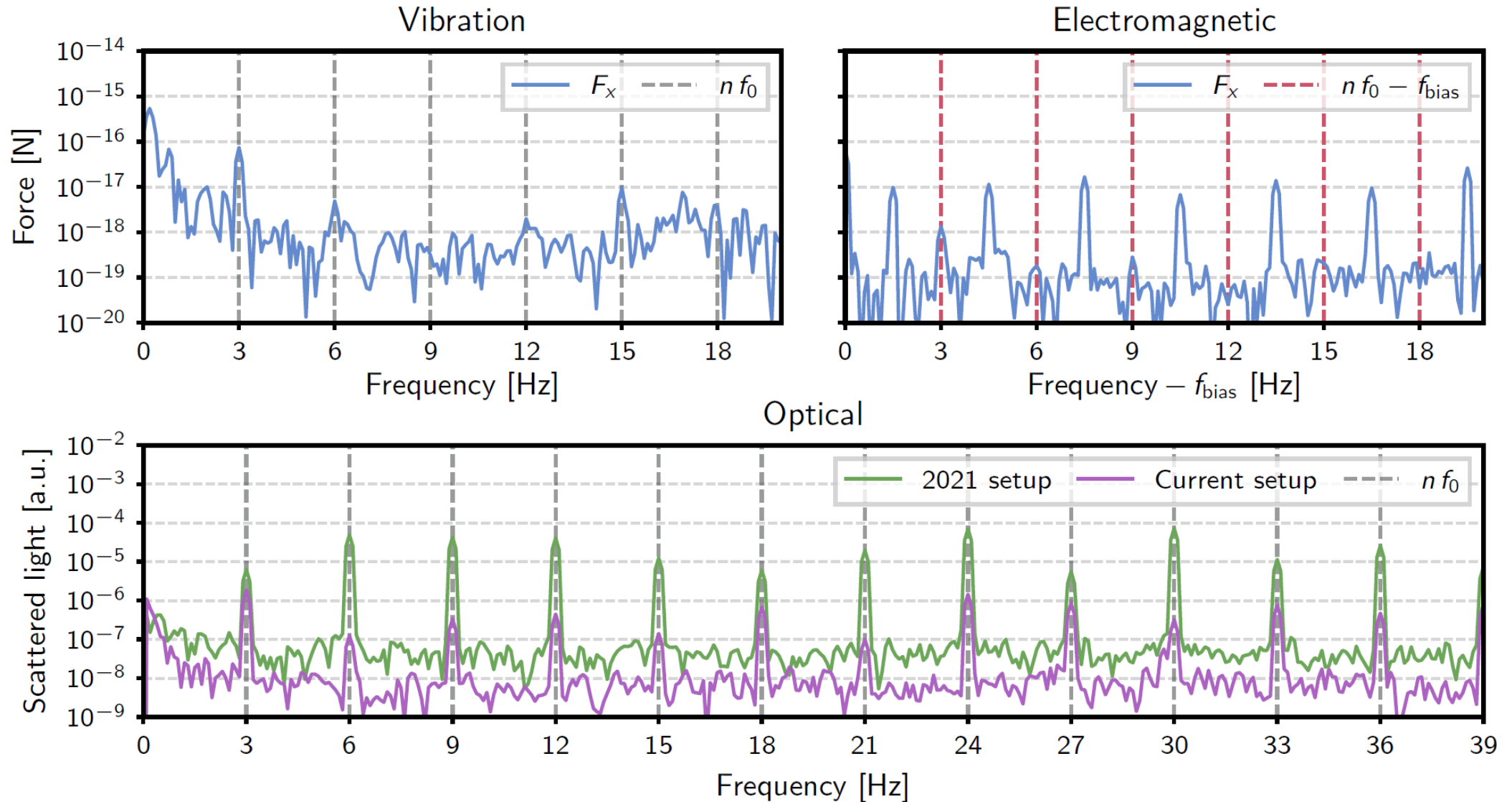
*Before plating*



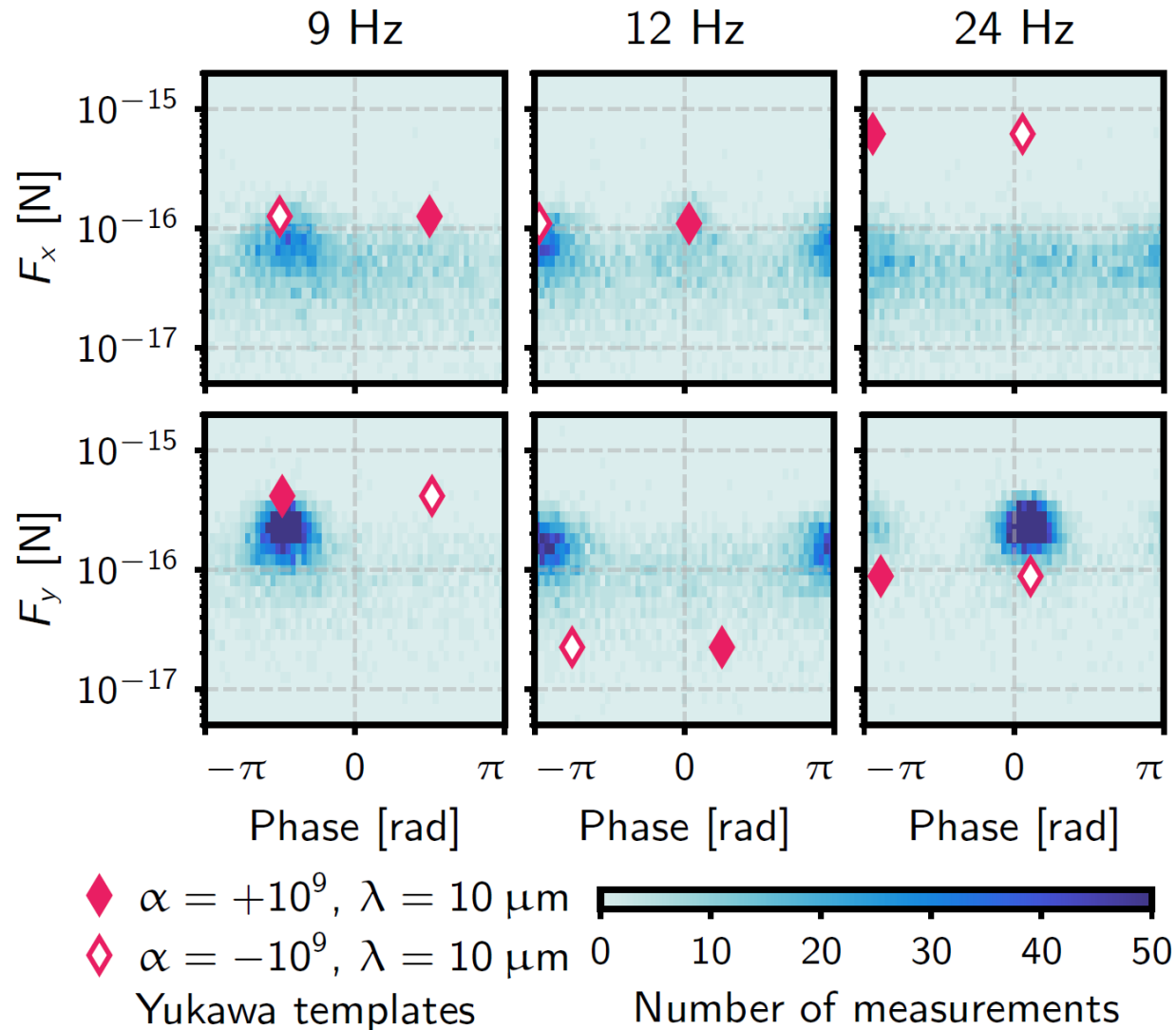
*After plating,  
longer exposure*



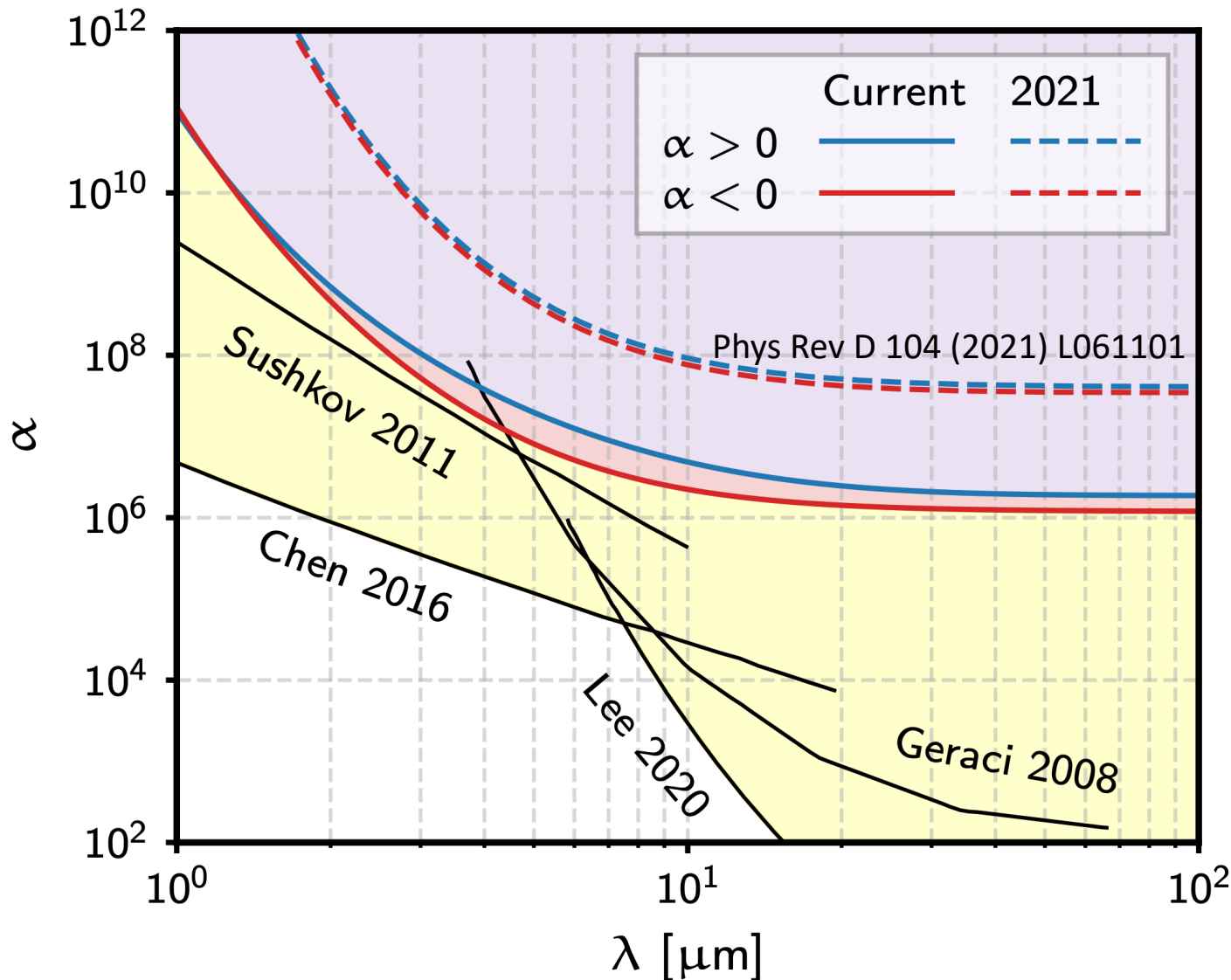
**Backgrounds are still the most serious limitation to the sensitivity,**  
*but the dominant optical one was reduced by  $\sim 2$  orders of magnitude*



The force measurement in 2 coordinates (eventually 3), its phase, and the use of the six harmonics with the most power better constrains the measurement.



# Result



- Limit is set using profile-likelihood approach
- Better understanding and reduction of background:  
➔ 100x improvement
- Setting limit on positive and negative coupling constants
- First measurement of the force vector
- Sensitivity still limited by diffuse light backgrounds (and, more is likely lurking under that)
- Working on more background abating techniques  
➔ Most notably a 100pixel, fast, hi res photodetector

## Now I want to describe a different technique, using nuclei (as opposed to atoms) as force sensors

*GG, D.E. Kaplan, S. Rajendran, Phys Rev D 102 (2020) 115031  
See also Poster by C. Brandenstein at this meeting.*

***“Nuclei are well-protected affairs”***

They have electric charge, but that is screened by the electron cloud and so has little coupling to external E&M disturbances.

In addition, nuclear level shifts due to E&M coupling, occurs through coupling to multipole (mainly dipole) moments and these are suppressed by the size of the nucleus.

And, this is further suppressed, in the case of unpolarized nuclei, by  $\sqrt{N}$ , when looking for the shift of a spectroscopy line that is measured by  $N$  events.

**→ Can nuclear (gamma) spectroscopy used to investigate new forces?**



Glen Rebka at the basement station

Zeitschrift für Physik, Bd. 151, S. 124–143 (1958)

Aus dem Institut für Physik im Max-Planck-Institut für medizinische Forschung,  
Heidelberg

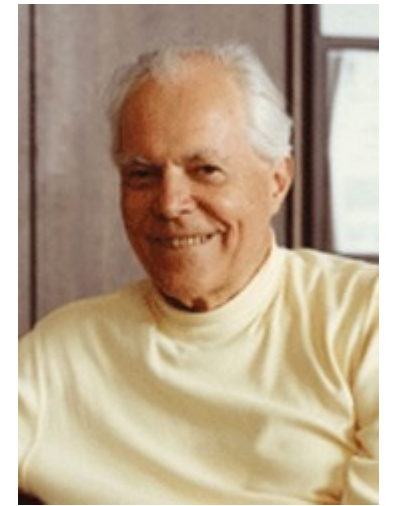
### **Kernresonanzfluoreszenz von Gammastrahlung in $\text{Ir}^{191}$**

Von

RUDOLF L. MÖSSBAUER\*

Mit 8 Figuren im Text

(Eingegangen am 9. Januar 1958)



Source: Pontifical Academy of Sciences

**Mössbauer spectroscopy was used in an elegant experiment to detect, for the first time, gravitational red/blue shift of photons.** Pound & Rebka Physical Review Letters. 4 (1960) 337

**Somewhat disconcertingly, Mössbauer spectroscopy was then “appropriated” by Chemists!**

**→ It's time to reclaim it back for fundamental physics!**

*Incidentally  $^{229}\text{Th}$  clock transition is a special case of this.*

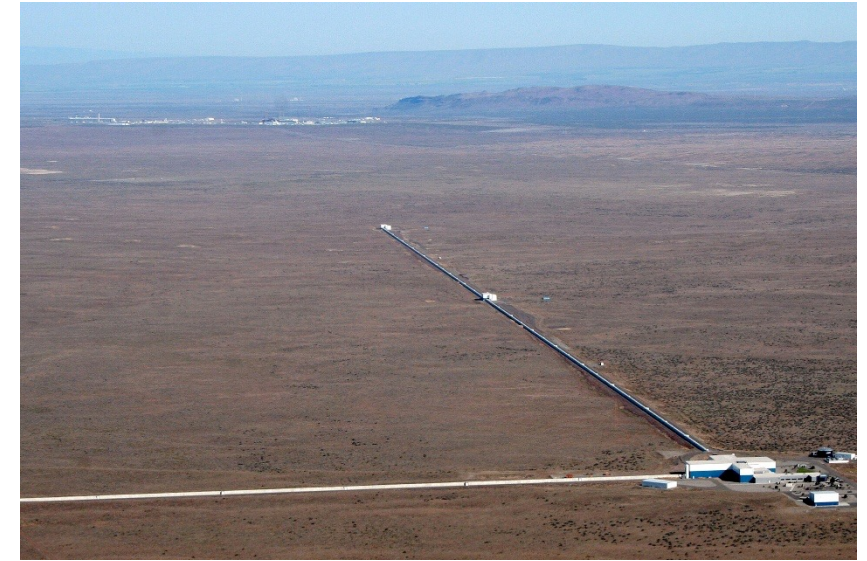
# In some cases, the natural line width is directly accessible

Used mainly in chemistry

Nuclide	$E$ (eV)	$T_{1/2}$	$\Gamma$ (eV)	$\Gamma/E$
$^{57}_{26}\text{Fe}$	14,413	98.3 ns	$4.7 \times 10^{-9}$	$6.4 \times 10^{-13}$
$^{73}_{32}\text{Ge}$	13,328	2.92 $\mu\text{s}$	$1.6 \times 10^{-10}$	$1.2 \times 10^{-14}$
$^{181}_{73}\text{Ta}$	6,237	6.05 $\mu\text{s}$	$7.5 \times 10^{-11}$	$1.2 \times 10^{-14}$
$^{67}_{30}\text{Zn}$	93,300	9.07 $\mu\text{s}$	$5.0 \times 10^{-11}$	$5.4 \times 10^{-16}$

Essentially unexplored

$^{45}_{21}\text{Sc}$	12,400	318 ms	$1.4 \times 10^{-15}$	$1.13 \times 10^{-19}$
$^{107}_{47}\text{Ag}$	93,125	44.3 s	$1.03 \times 10^{-17}$	$1.1 \times 10^{-22}$
$^{103}_{45}\text{Rh}$	39,753	56.1 min	$1.36 \times 10^{-19}$	$3.4 \times 10^{-24}$
$^{189}_{76}\text{Os}$	30,814	5.8 hr	$2.2 \times 10^{-20}$	$7.0 \times 10^{-25}$



For reference, aLIGO strain sensitivity:  
 $\delta l/l \sim 10^{-23}/\sqrt{\text{Hz}}$

# Principle of the experiment

The new, hypothetical interaction, perturbs the isomeric state, slightly changing the position of the line.

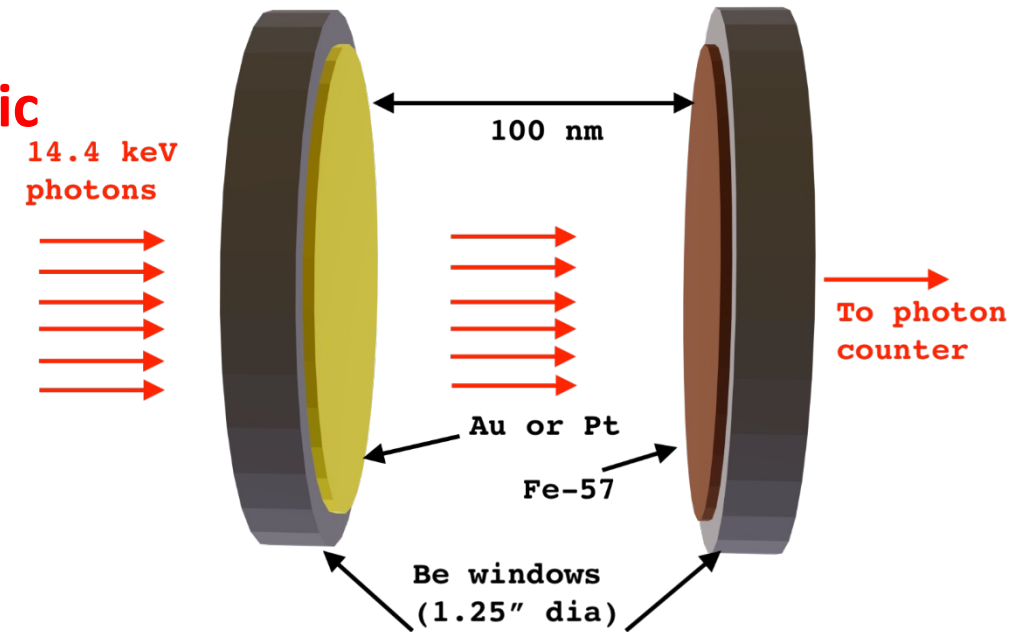
By design, the experiment is insensitive to vector interactions, highly suppressing EM backgrounds.

Note that we are not sensing a force but the effect of the field on a bound (nuclear) state:  
*a rather new concept of measurement.*

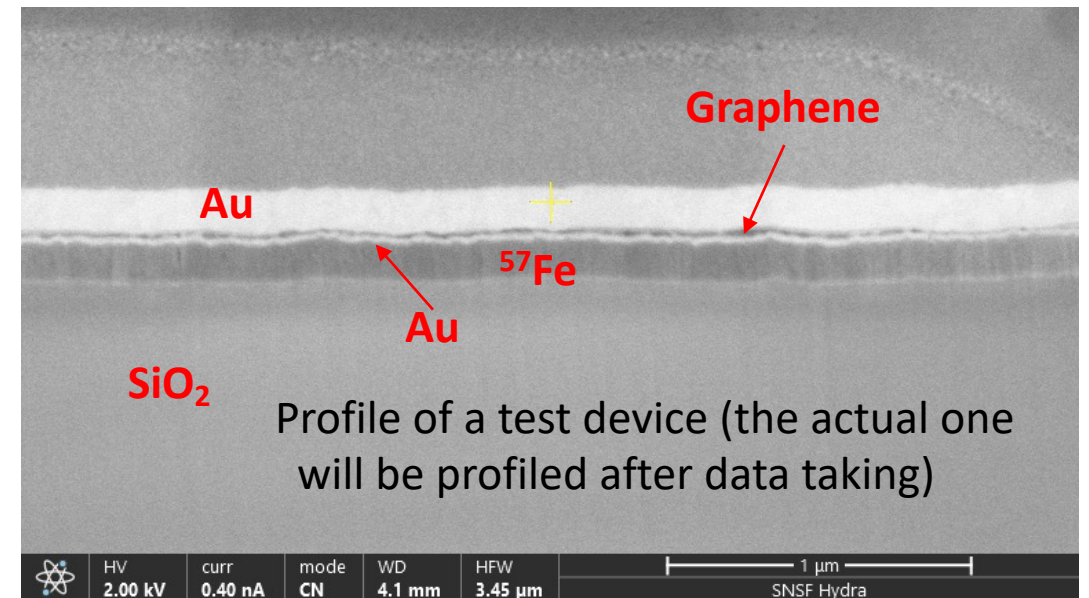
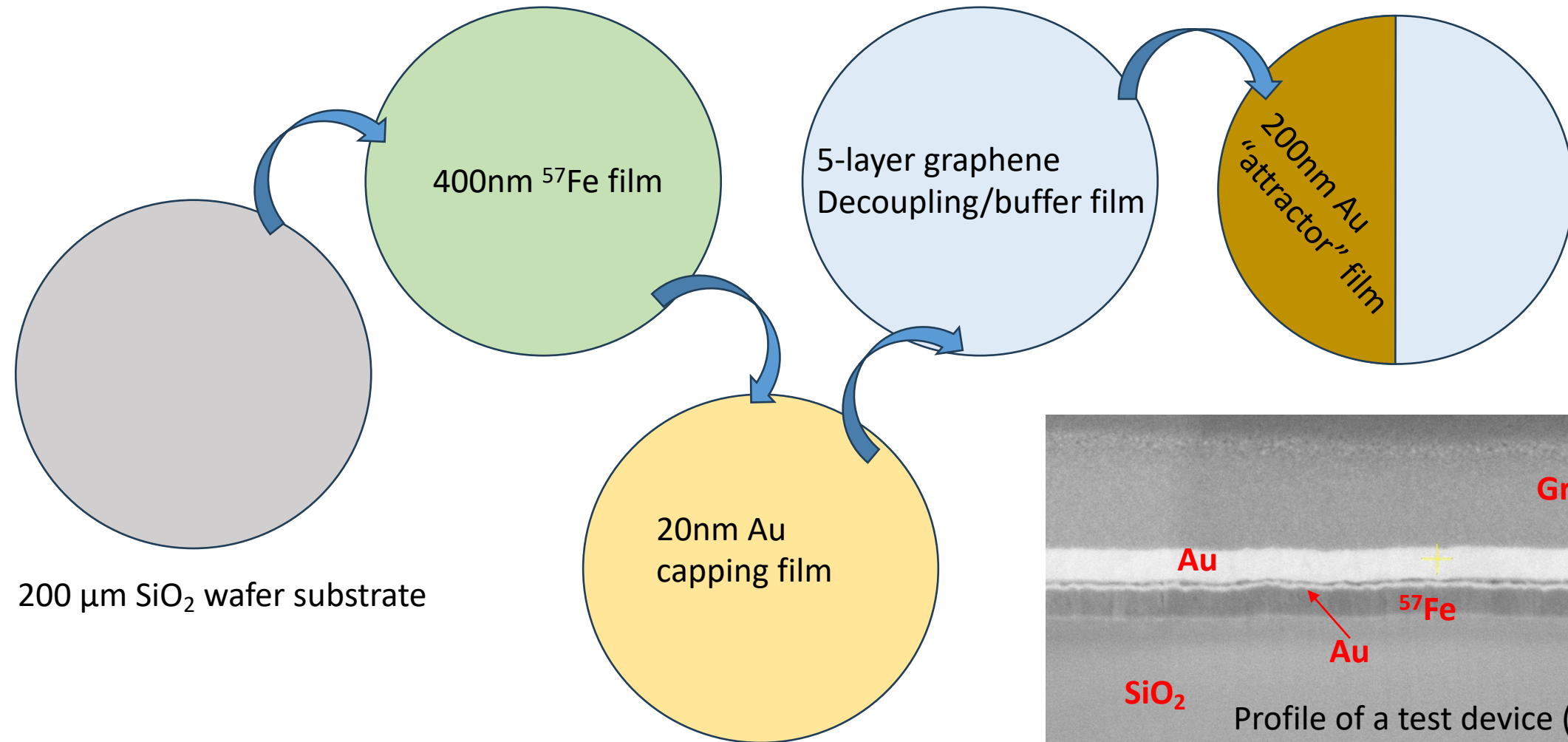
*Unfortunately, there is nuclear physics between a measured effect and the line shift potentially observed by an experiment.*

*Yet, a line shift, as long as real, is a discovery, irrespective of the nuclear physics.*

*...reminds me of neutrinoless double beta decay!*



# A first experiment with $^{57}\text{Fe}$ : *Integrated resonant absorber and attractor sourcing the new interaction*



The sensitivity to an energy shift depends on the linewidth ( $\Gamma$ ), the contrast ( $C$ ), and the counting statistics ( $n$ )

$$\delta E_{min} = \frac{T(v_0)}{2\sqrt{n}} \left| \frac{dT}{dv} \right|_{v_0}^{-1} \sim \frac{T(v_0) \Gamma}{2\sqrt{n} C}$$

We are initially shooting for

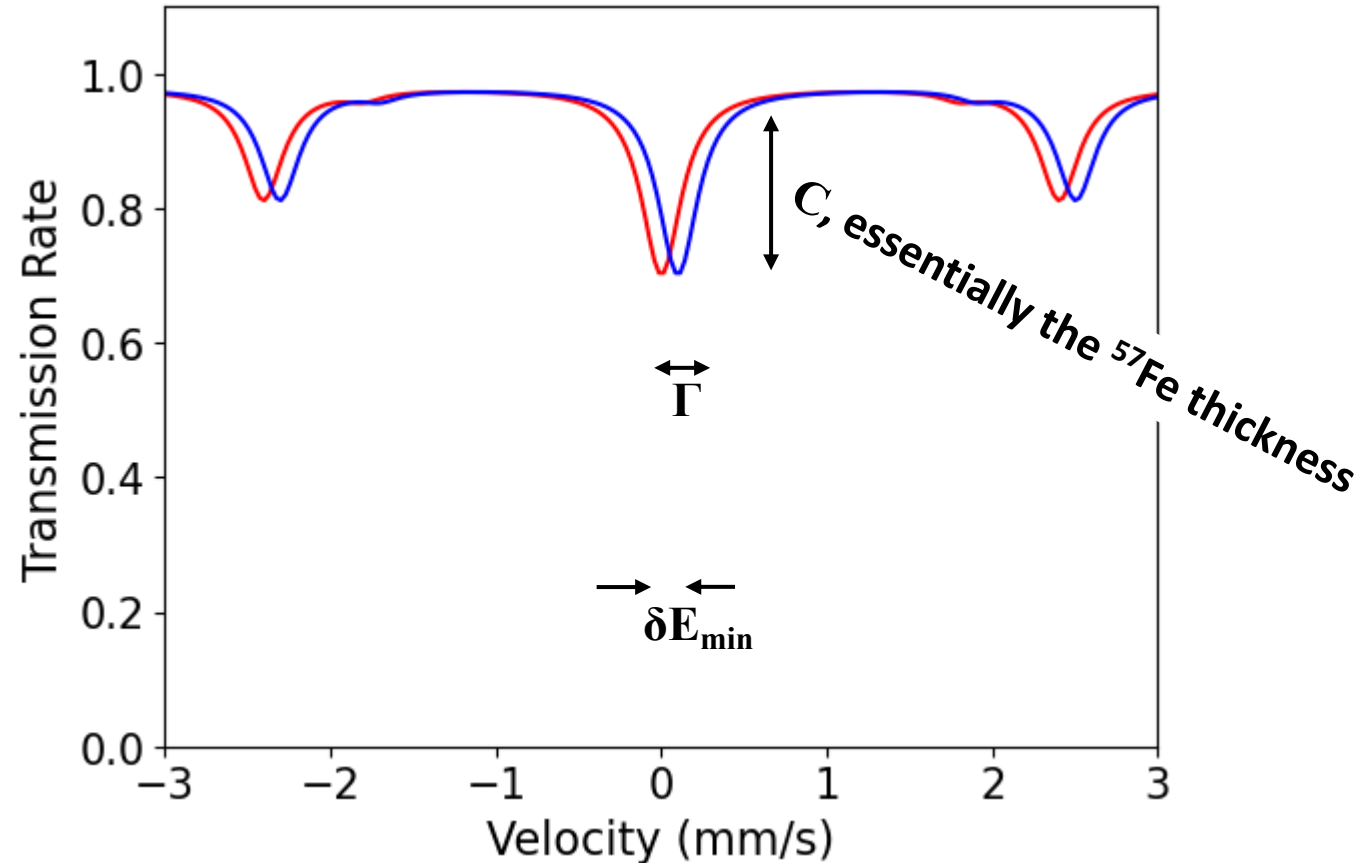
$$\delta E_{min} = 3 \times 10^{-14} \text{ eV}^*$$

$$\rightarrow \Gamma = 0.28 \text{ mm/s} = 1.3 \times 10^{-8} \text{ eV}$$

$$C = 0.18$$

$$n = 2 \times 10^{12} \text{ on each side}$$

$$(100 \text{ days @ } 2.5 \times 10^5 \text{ evts/s})$$



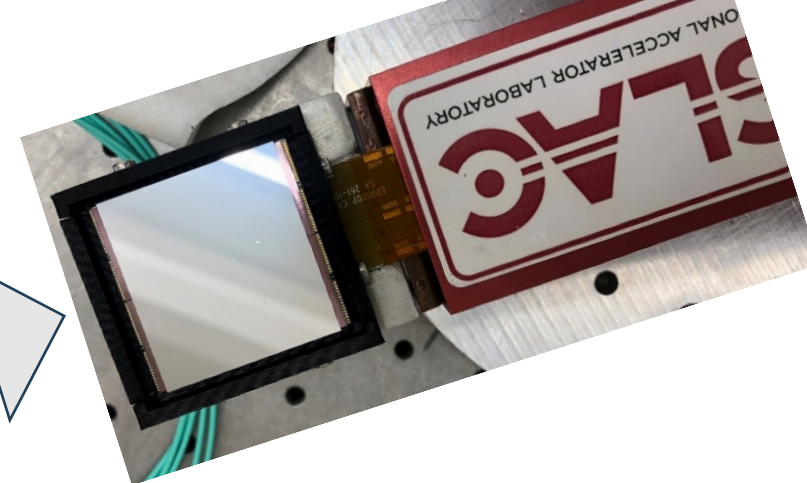
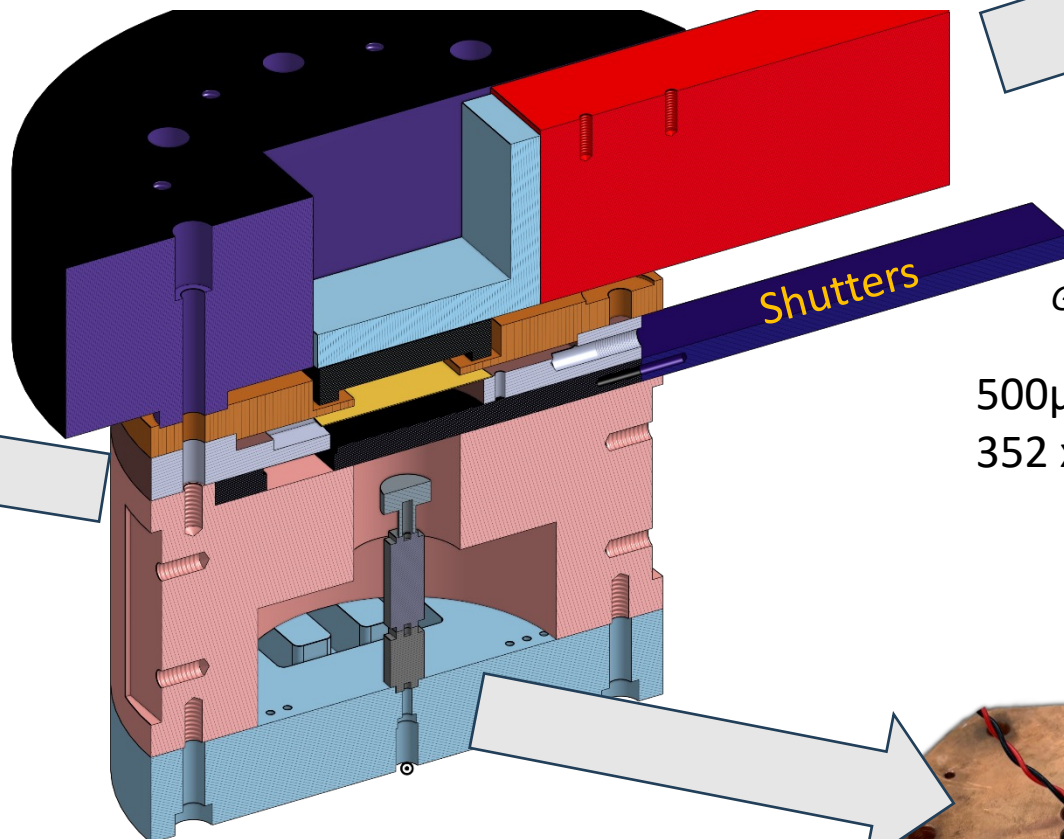
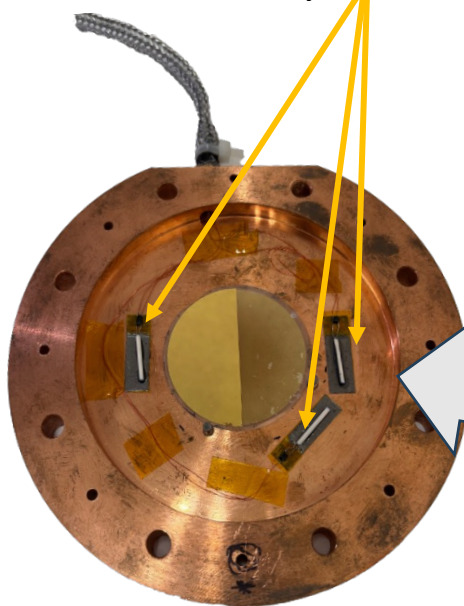
\* The Pound & Rebka initial (final) experiment had a  $8.4 \times 10^{-12} \text{ eV}$  ( $3.5 \times 10^{-12} \text{ eV}$ ) sensitivity

## Known systematics: $\delta E = 10^{-15}$ eV is equivalent to:

- **Temperature**  $\Delta E \sim 10^{-11} \left(\frac{T}{300K}\right)^3 eV/K$  0.1 mK
- **Casimir**  $\Delta E \sim 3 \cdot 10^{-15} \left(\frac{10 \text{ nm}}{d}\right)^4 eV$  20 nm
- **Magnetic**  $\Delta E \ll 10^{-8} eV/T$  50 nT
- **Pressure**  $\Delta E \sim 10^{-10} eV/GPa$  10 kPa

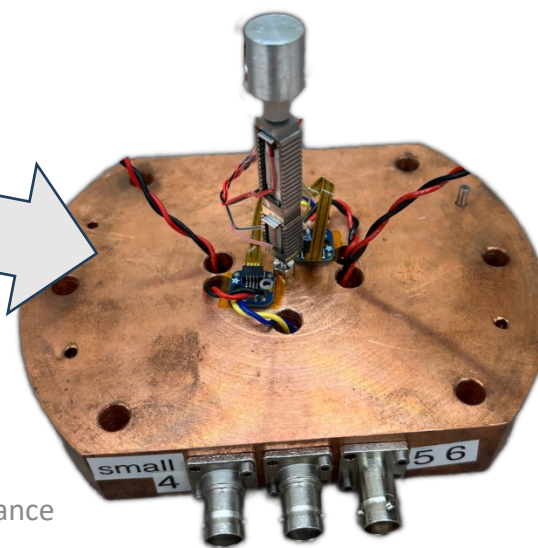
**Across the size of absorber (~5cm diameter)**

Resonant absorber with Au 1/2 coating anchored to a copper plate with absolute and differential thermometry



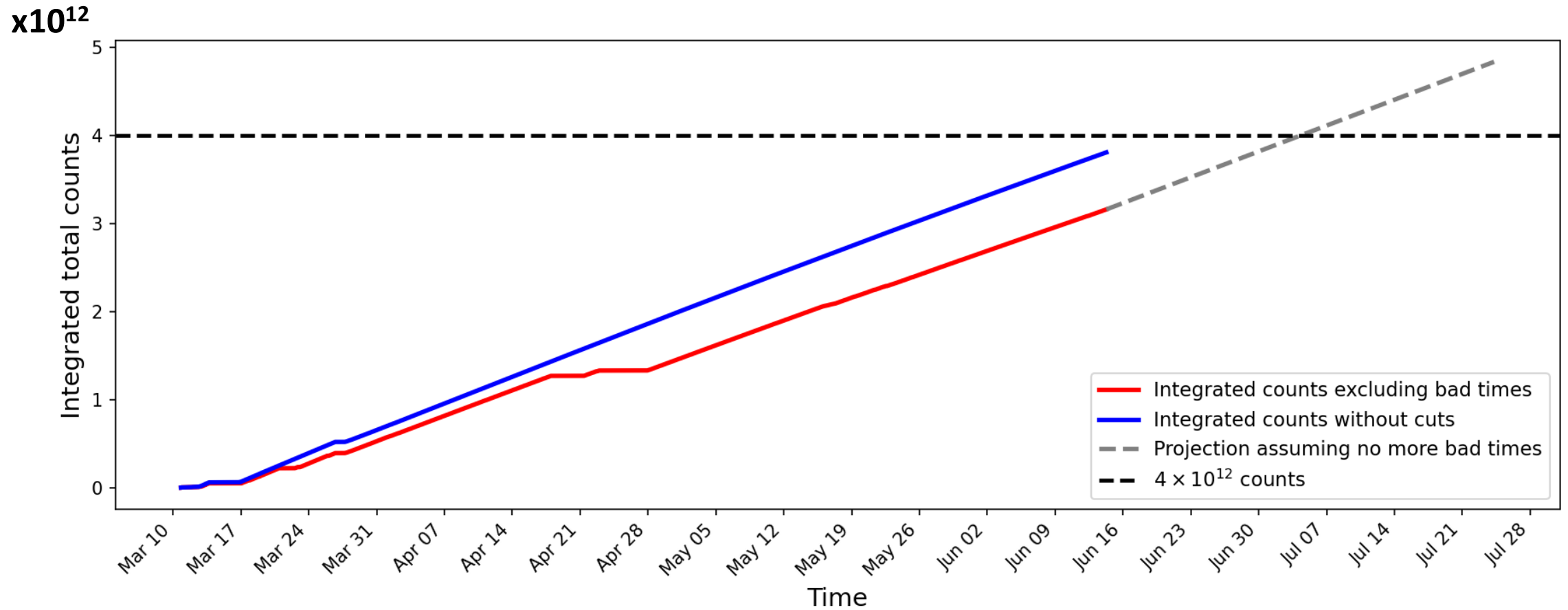
EpiX-10KA X-ray camera (SLAC)  
*G. Blaj et al., Proc. 13th Intl Conf on Synch. Rad. Instrum.*  
<https://doi.org/10.1063/1.5084693>

500µm thick Si (75% QE @ 14.4keV)  
352 x 384 pixels      1 kFrame/s



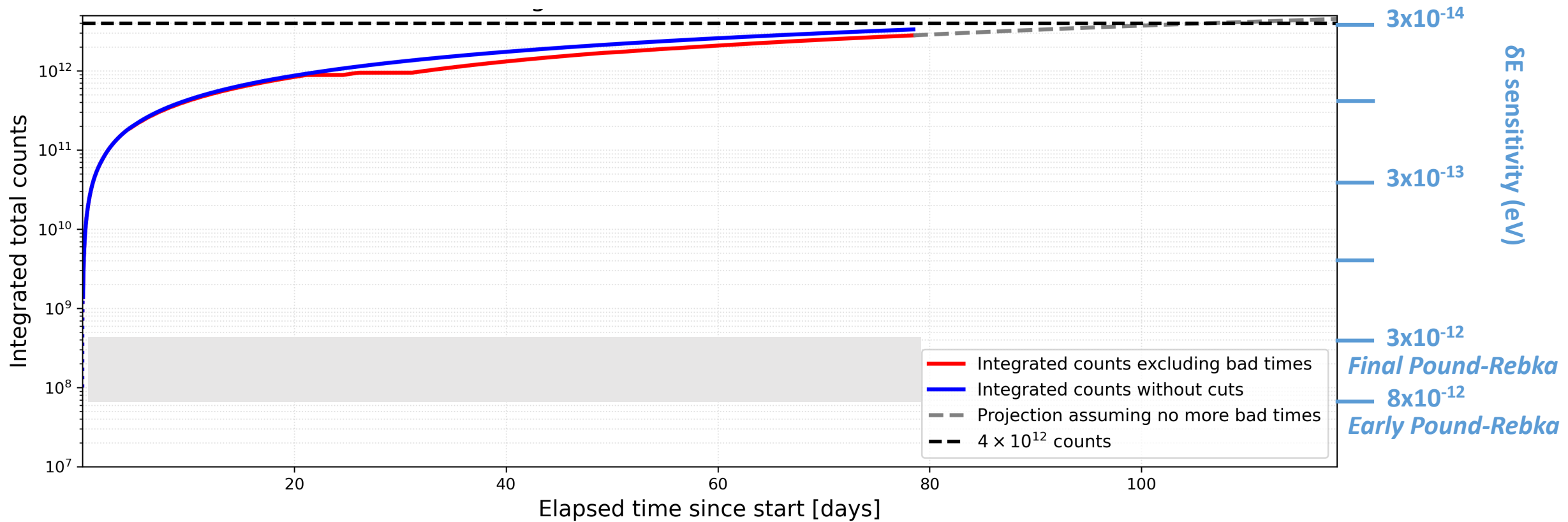
The source moves at 0.1mm/s with 10pm/s precision

**We already have ~100 days of data on disk, using a ~30mCi  $^{57}\text{Co}$  source.**



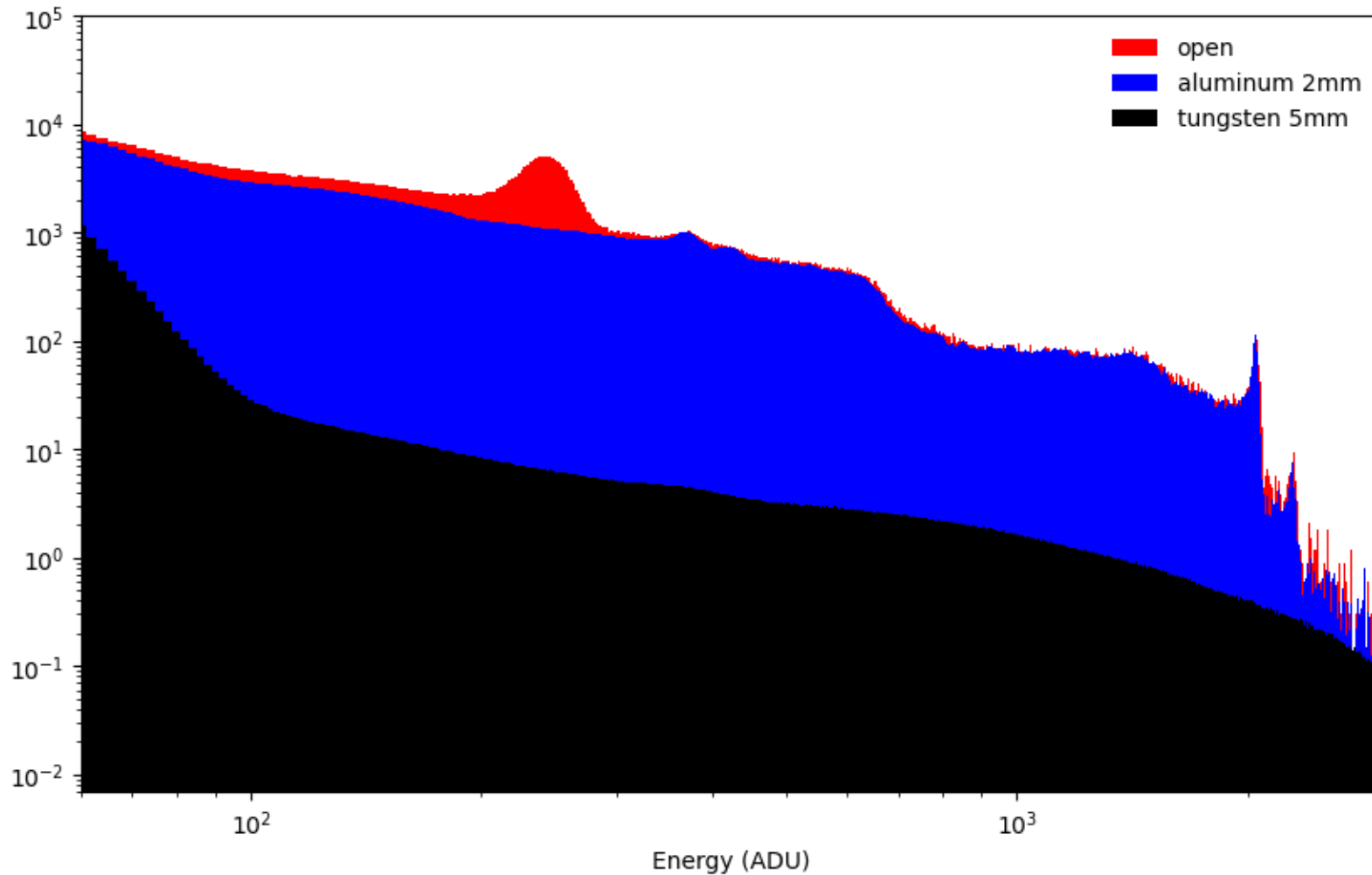
# Probably the most sensitive Mössbauer experiment ever done!

...and we are likely to be the first imaging Mössbauer experiment



- 1) Being compact makes it easier
- 2) Being compact means a large solid angle, which makes it harder  
... as we will see in a minute

# Single pixel energy spectrum with different shutter configurations for 1 day of data



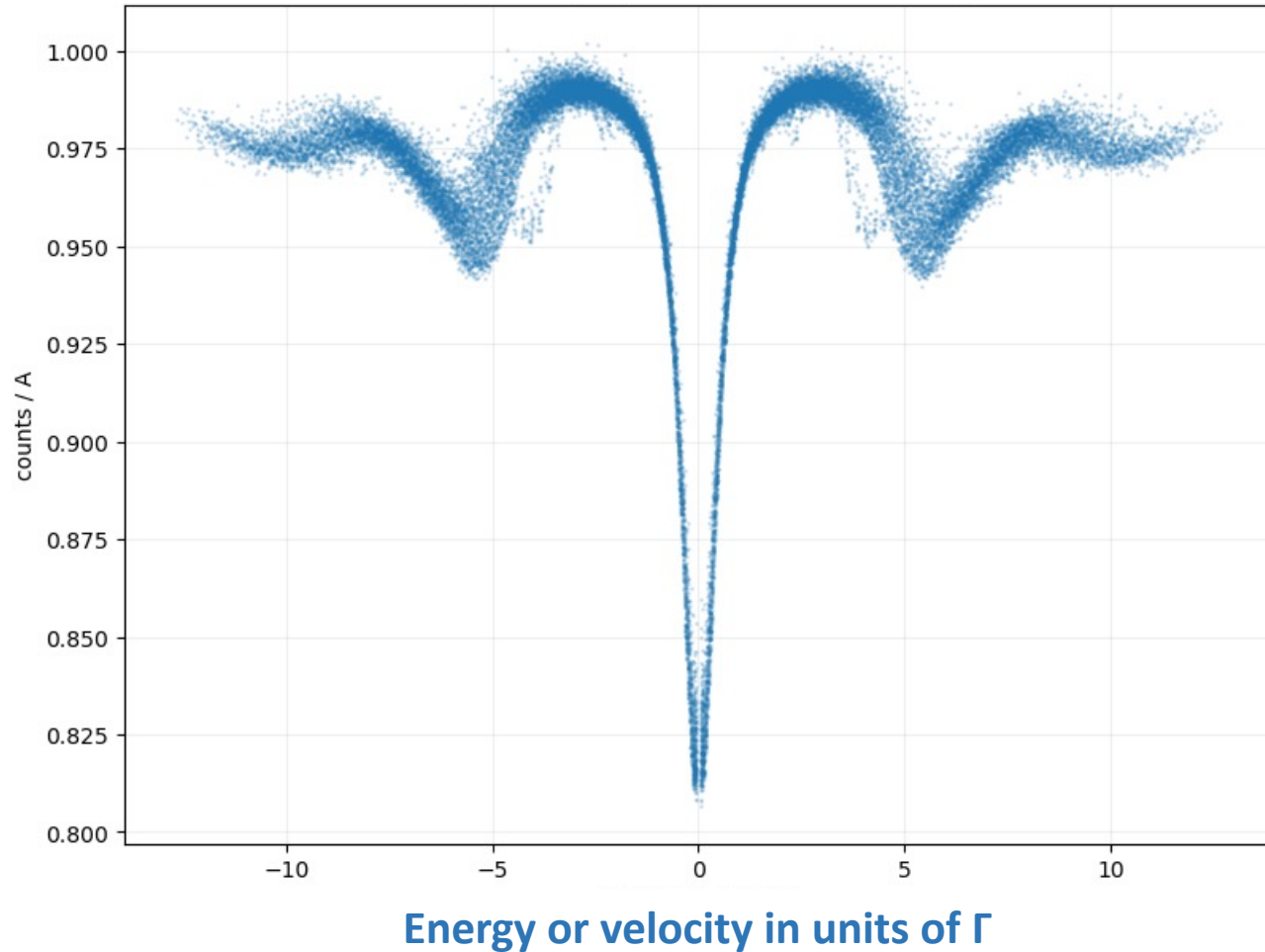
Closing the tungsten shutter eliminates all signals other than electronic noise and generic radioactive/cosmic ray background.

An aluminum shutter removes 99.2% of the 14.4keV gammas, but transmits 92% of the 122.1keV gammas.

With no shutters the 14.4keV peak is clearly visible with a rate of  $\sim 5 \text{ evt s}^{-1} \text{ pixel}^{-1}$  or  $\sim 0.7 \text{ Mevt s}^{-1}$  for the entire camera.

The rate on the camera is higher because of the other lines.

Now a broader velocity scan for the entire camera:  
250 velocities, 5 min each, 20 hr data total @  $5 \times 10^5$  ev/s

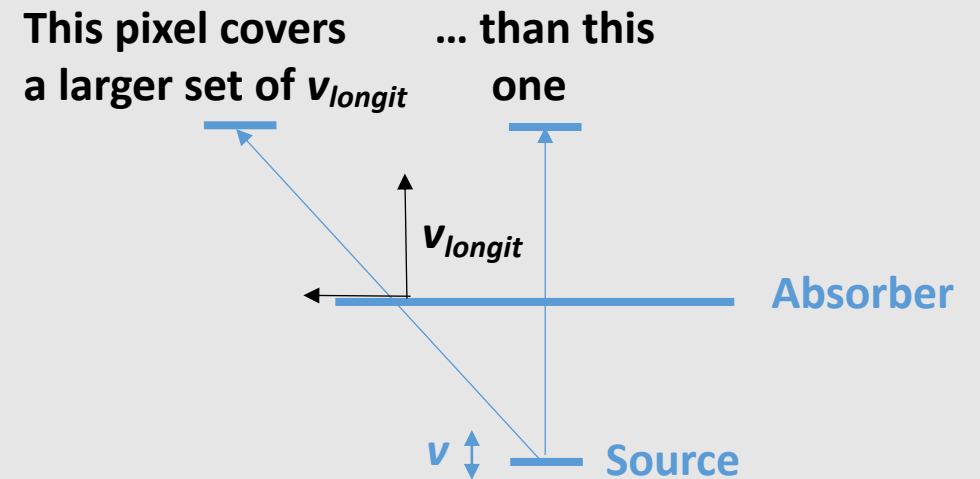


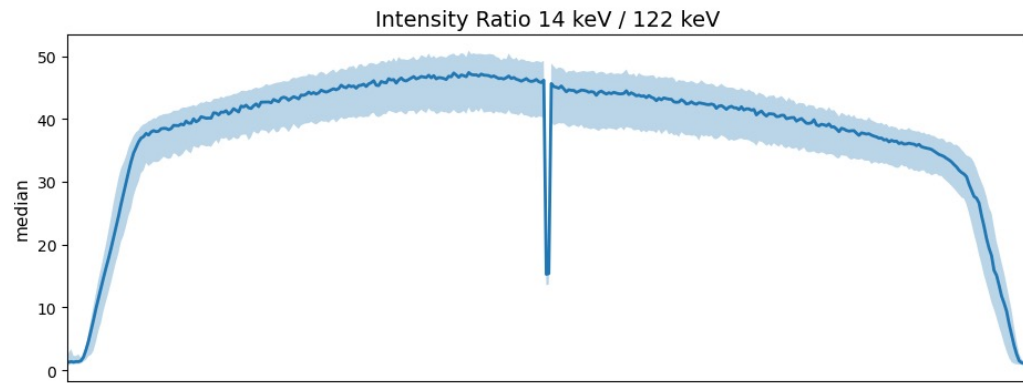
This data is used for calibration and to check that things are well-understood.

The “science data” only covers the central line.

The side-lines are asymmetric, as the longitudinal component the velocity depends on the velocity.

In addition, they are “fuzzy” because pixels at different azimuthal angles cover different ranges of (longitudinal) velocity.



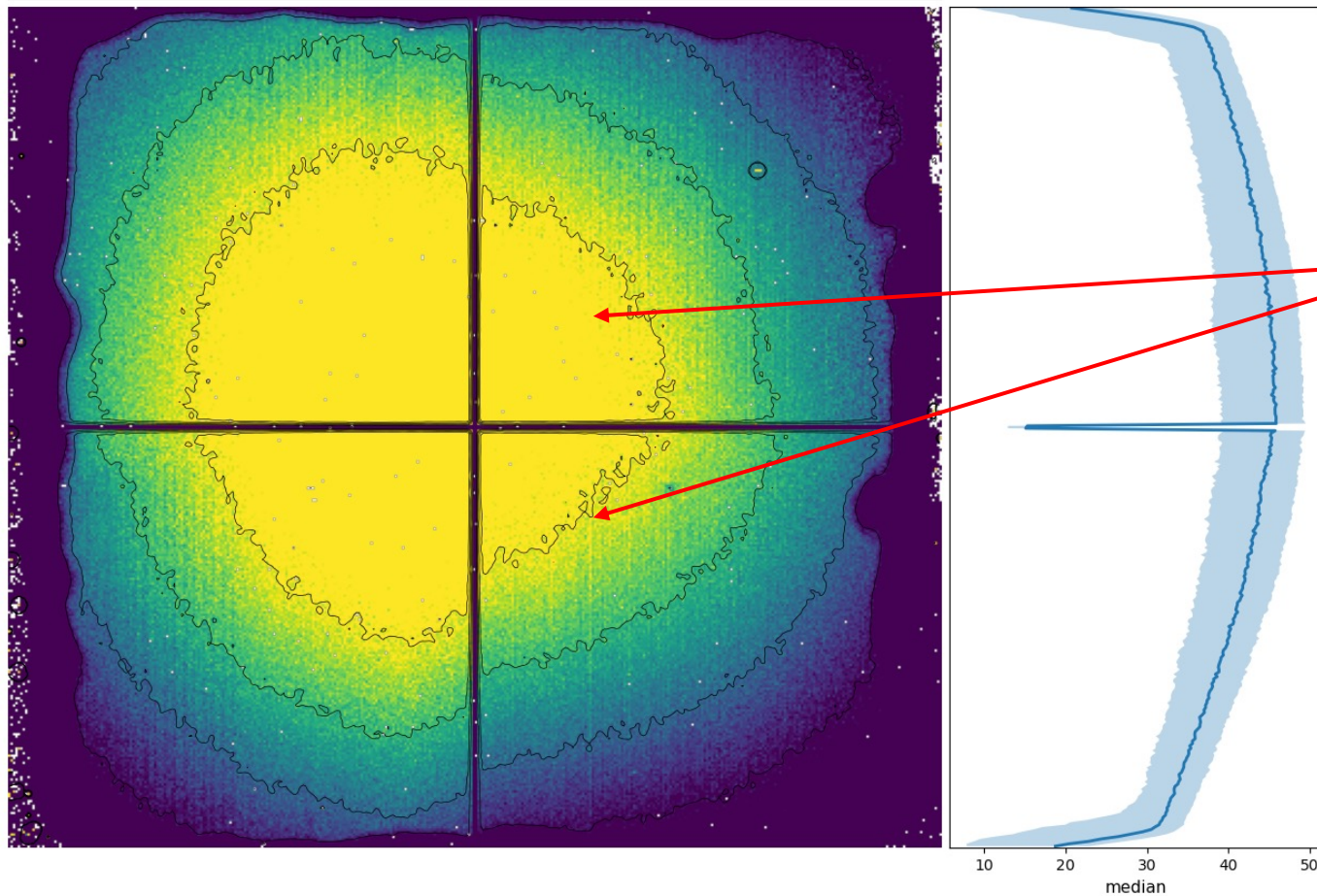


1 day at  $\sim 0.5$  Mcts/s (for 14.4 keV events)

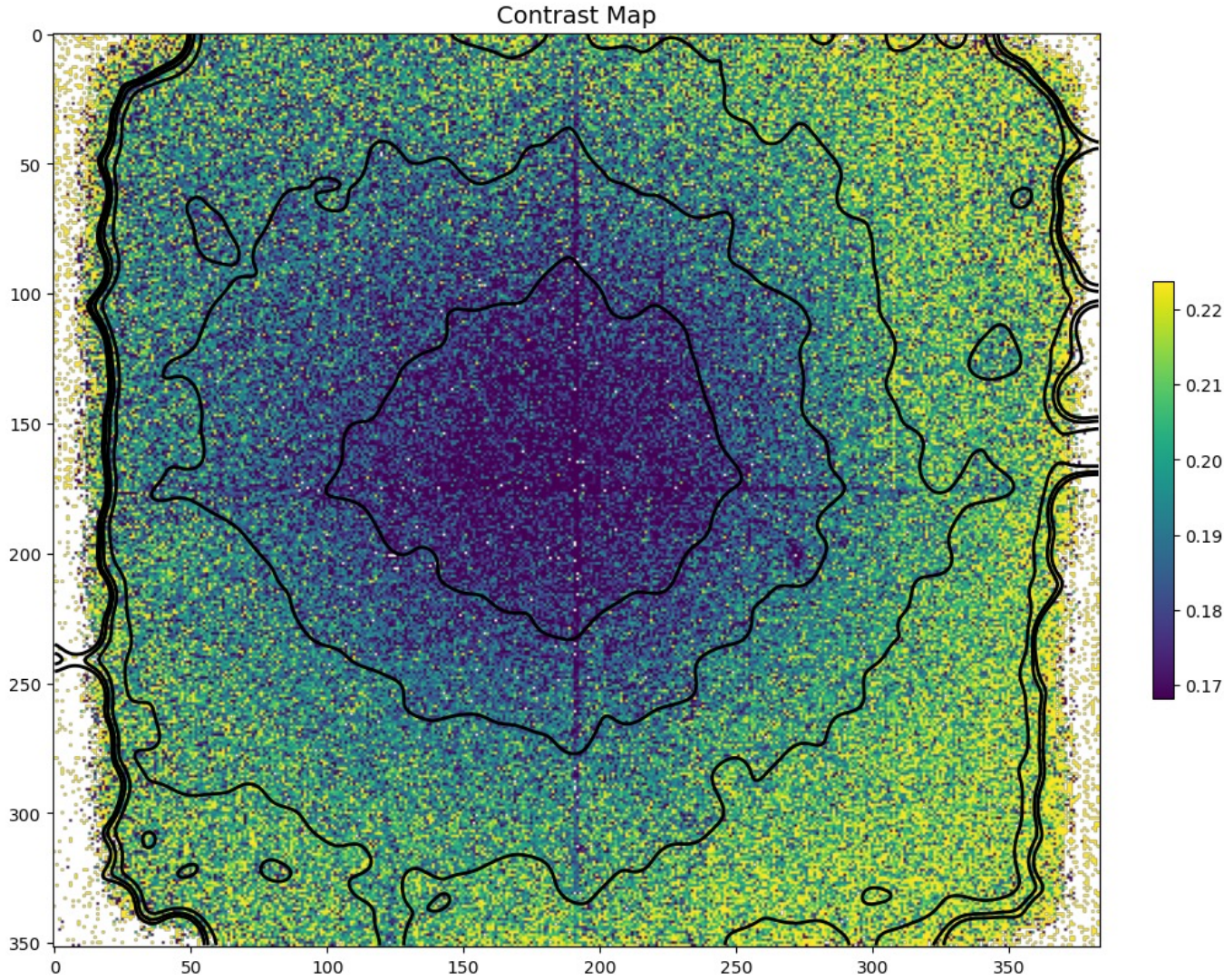
Note that the source is not perfectly centered over the camera.

Since the 122.1 keV go thru the Au with little losses, the 14.4/122.1 ratio readily shows where the Au is.

*One may ask: why don't you show a Geant simulation of this. Answer: because Geant does not have a Mössbauer package. We are fixing this now.*

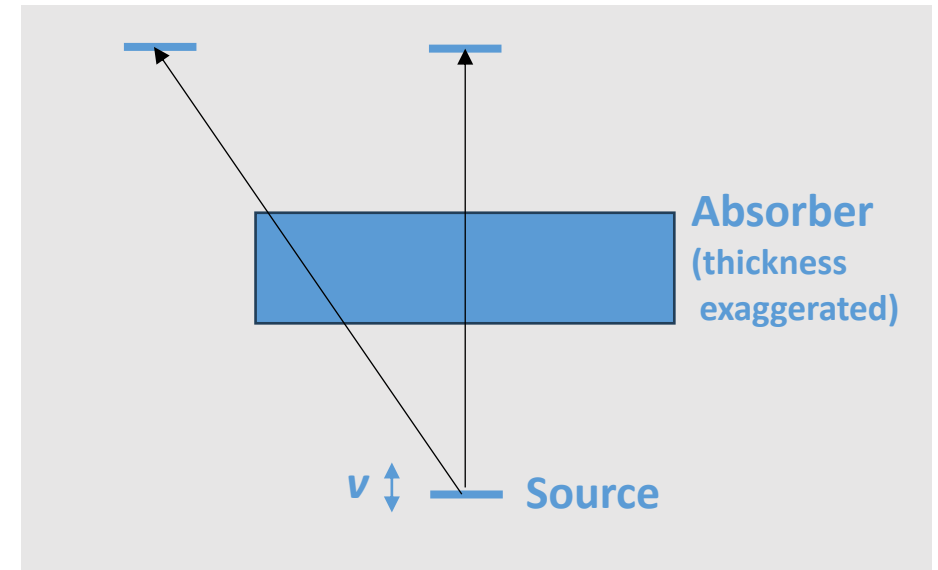


# Resonant absorption contrast map



The contrast only depends  
on the  $^{57}\text{Fe}$  thickness,

so, it increases at large angles because the slant  
depth increases with the azimuthal angle.



## Next on the Mössbauer front

A Geant Mössbauer package is in preparation

New experiment with a 100 mCi source (3x statistics, we may change the absorber/attractor parameters)

Developing a natural linewidth  $^{181}\text{Ta}$  source –should be ~150x narrower, although no one has ever gotten close to the natural width.

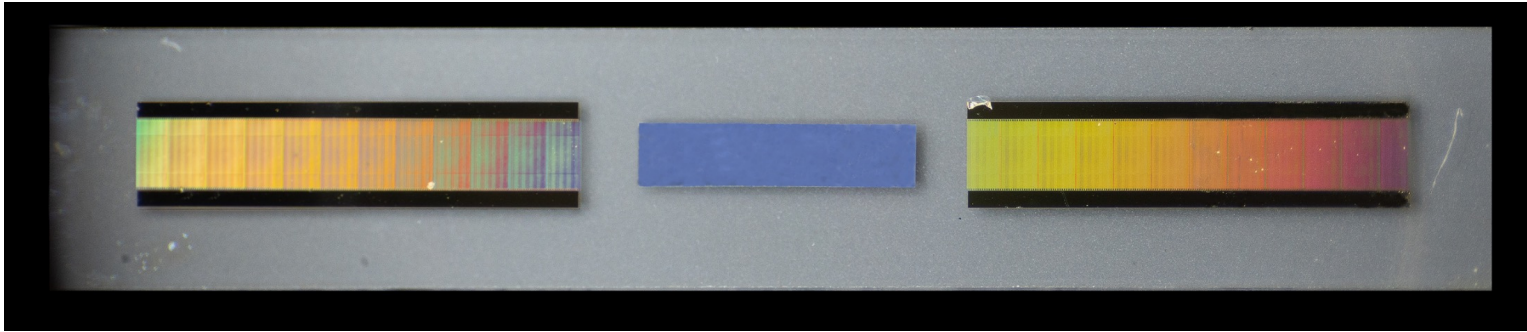
Rather tricky:

- procure enriched  $^{180}\text{W}$
- produce a  $^{180}\text{W}$  film on MgO substrate
- neutron activate to  $^{181}\text{W}$  in a reactor
- absorber is a  $^{181}\text{Ta}$  on a Be substrate

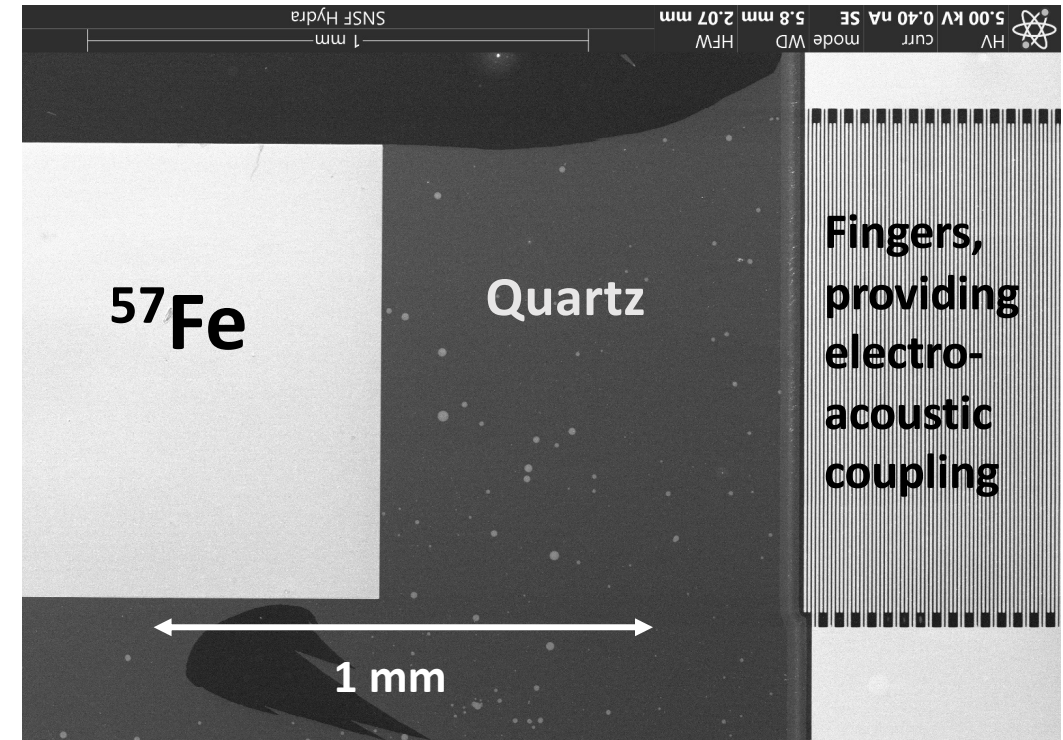
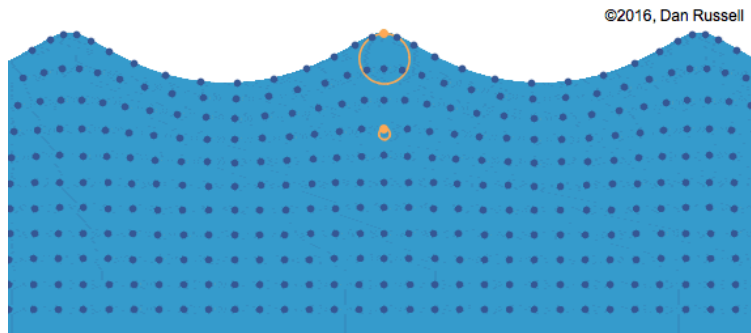
Internal conversion is larger for Ta (than for Fe) and, in addition, there is a velocity shift between source and absorber.

# Fun side-project: a $\gamma$ -ray AOM

*Nazeeri et al. PRL 136 18 (2026)*



The resonant absorber is now deposited over a region of quartz where a surface acoustic wave propagates.

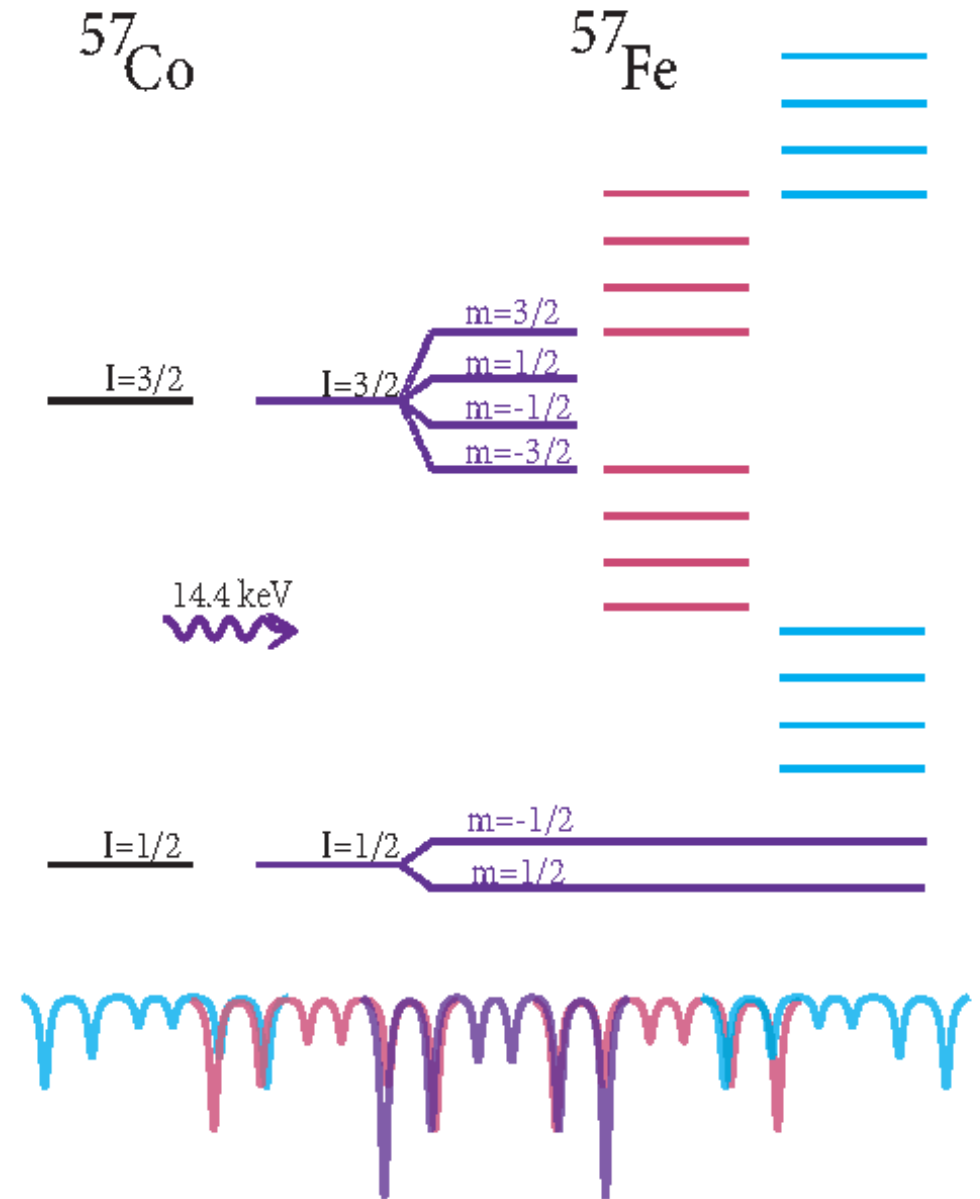


# Expected effect

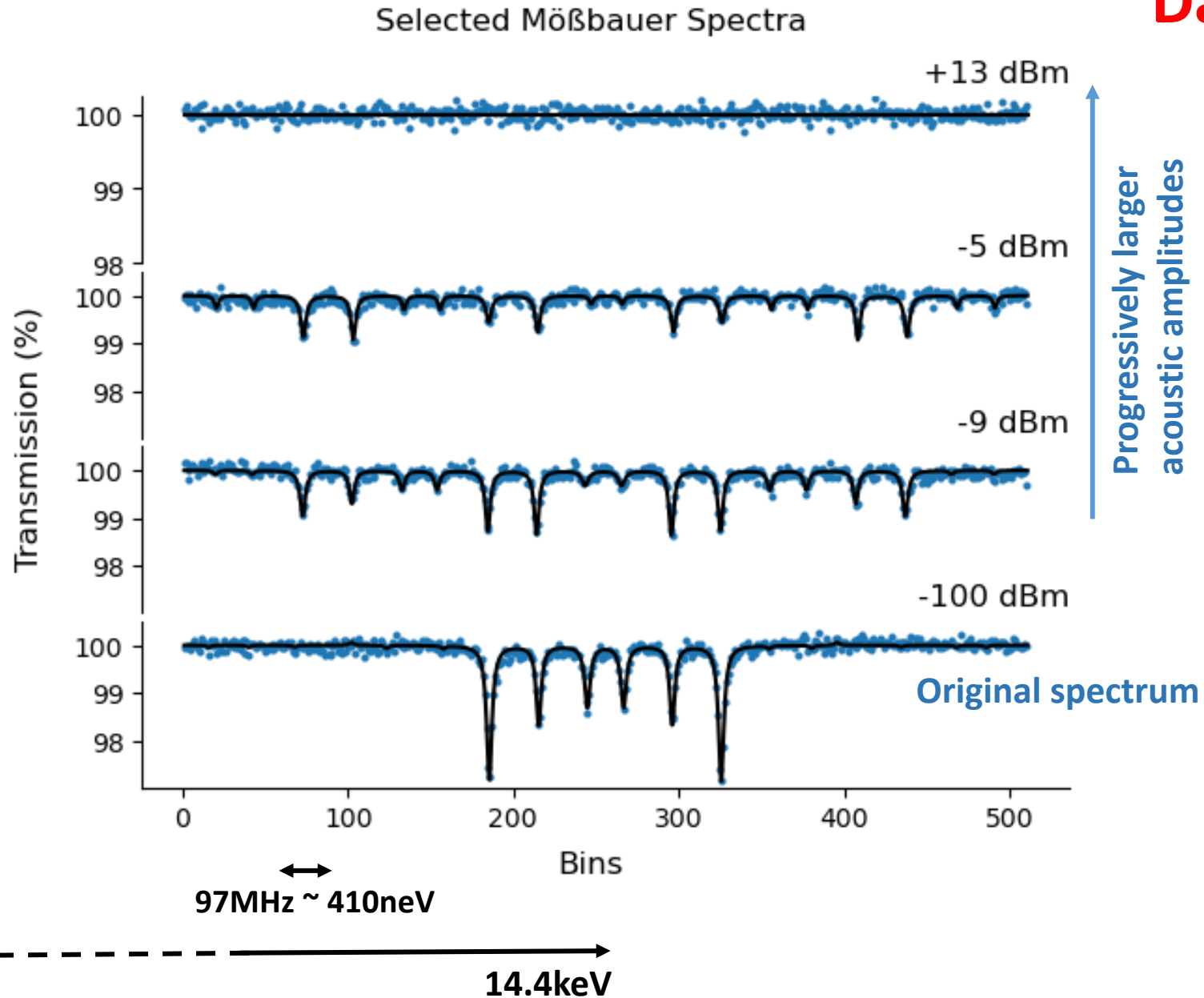
Like in an AOM, phonons in the quartz transfer energy to the  $\gamma$ -rays, producing higher order modes.

This is happening in the regime where  $\tau_{\text{SAW}} \sim 10\text{ns} < \tau_{\text{Fe-57}} \sim T_{1/2}/\sqrt{2} \sim 70\text{ns}$

In our device, the acoustic frequency, 97 MHz, is almost exactly the same as the energy difference between different  $m$  states, producing various degeneracies.



# Data



Note that the integral of counts is constant in the different spectra.

As the acoustic amplitude increases, a larger fraction of the integral is contributed by sidebands away from the center.

Note that at max amplitude, the counts are all in far away sidebands (out of range for the spectrometer).

# Concluding, the current cast



Chiara Brandenstein  
Lorenzo Magrini

Weiran Xu  
Kenneth Kohs

Saul Barcañel-Salazar

Chengjie Jia

Malayne Perry

Lin Si

Zhengruilong Wang

Meimei Liu

Jaqueline Huang

Ralph deVoe

