

Fermions in Quantum Complexity Theory

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Occupation Number Formalism

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States of this form (*Fock states*) make up a subspace H_n , of states describing n particles in m modes. The full *Fock space* is given by

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To go back to $H = H^{\otimes n}$, (anti-)symmetrize over all possible combinations giving x_j particles in mode j .

Moving Between Subspaces

Start with the vacuum state $|0, \dots, 0\rangle$. To add a particle to mode j , apply the *creation operator* a_j^\dagger . By Hermitian adjoint, a_j removes a particle from mode j (*destruction operator*).

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- ▶ A state of n fermions in modes j_1, \dots, j_n is created by

$$\epsilon_{j_1 \dots j_n} \cdot a_{j_1}^\dagger \cdots a_{j_n}^\dagger |0, \dots, 0\rangle$$

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The *number of particles* in mode j is the eigenvalue of $N_j = a_j^\dagger a_j$.

Fermionic Quantum Computation

Any unitary operator acting on H can be generated by the set of creation and destruction operators: $U = \exp(iH)$, where $H(a_1, a_1^\dagger, \dots, a_m, a_m^\dagger)$ is Hermitian.

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Terhal and DiVincenzo (2002) showed quantum computation using a system of noninteracting fermions can be simulated classically, even with adaptive von Neumann measurements. Results were equivalent to a subclass of Valiant's matchgates (2001).

Noninteracting Fermions: Complexity

Consider special case where number of fermions is preserved. A computation step acting between modes α and β is a unitary $U = \exp(iH)$ generated by

$$H = h_{\alpha\alpha} a_{\alpha}^{\dagger} a_{\beta} + h_{\beta\beta} a_{\beta}^{\dagger} a_{\beta} + h_{\alpha\beta} a_{\alpha}^{\dagger} a_{\beta} + h_{\alpha\beta}^{*} a_{\beta}^{\dagger} a_{\alpha}$$

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Suppose U acts on a non-vacuum state with one fermion:

$$U a_j^{\dagger} |0\rangle = U a_j^{\dagger} U^{\dagger} U |0\rangle = U a_j^{\dagger} U^{\dagger} |0\rangle$$

where

$$U a_j^{\dagger} U^{\dagger} = \sum_k V_{jk} a_k^{\dagger}$$

Noninteracting Fermions: Complexity

Now calculate $\langle y|U|x\rangle$ for $|x\rangle = a_{j_1}^\dagger \cdots a_{j_l}^\dagger|0\rangle$ ($j_1 < \cdots < j_l$) and $|y\rangle = a_{k_1}^\dagger \cdots a_{k_l}^\dagger|0\rangle$ ($k_1 < \cdots < k_l$).

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$$\begin{aligned} U|x\rangle &= Ua_{j_1}^\dagger U^\dagger \cdots Ua_{j_l}^\dagger U^\dagger |0\rangle \\ &= \sum_{p_1, \dots, p_l} \left(V_{j_1}^{p_1} \cdots V_{j_l}^{p_l} \right) a_{p_1}^\dagger \cdots a_{p_l}^\dagger |0\rangle \end{aligned}$$

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The bracket is zero unless p_1, \dots, p_l is a permutation of k_1, \dots, k_l , and $\text{sgn}(\sigma)$ if the permutation is σ .

Noninteracting Fermions: Complexity

Rewrite sum in terms of permutations:

$$\langle y|U|x\rangle = \sum_{\sigma \in S_I} \text{sgn}(\sigma) V_{j_1}^{\sigma(k_1)} \cdot V_{j_I}^{\sigma(k_I)} = \det \tilde{V}$$

where \tilde{V} is $I \times I$ a submatrix of V with selected rows j_1, \dots, j_I and columns k_1, \dots, k_I .

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Quantum computation using noninteracting fermions and preserving the number of particles can be efficiently simulated classically.

Noninteracting Fermions: General Cases

- ▶ More generally, quantum computation involving noninteracting fermions need not conserve particle number, as long as it conserves particle parity.
- ▶ The more general case uses creation and destruction operators $c_{2j} = a_j + a_j^\dagger$ and $c_{2j+1} = -i(a_j - a_j^\dagger)$, satisfying $\{c_k, c_l\} = 2\delta_{kl}$.
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- ▶ Terhal and DiVincenzo proved this can also be classically efficiently simulated.
- ▶ Furthermore, complete von Neumann measurements also are not sufficient to bring the system to full universality, but it is for bosons—KLM (2007)

Universal Fermionic Quantum Computation

A universal set for quantum computation using interacting fermions was given by Bravyi and Kitaev (2000) for local fermion modes, consisting of the following unitaries:

- ▶ $\exp\left(i\frac{\pi}{4}a_0^\dagger a_0\right)$ (potential term)
- ▶ $\exp\left(i\frac{\pi}{4}(a_0^\dagger a_1 + a_1^\dagger a_0)\right)$ (tunnelling term)
- ▶ $\exp\left(i\frac{\pi}{4}(a_1 a_0 + a_0^\dagger a_1^\dagger)\right)$ (superconductor interaction term)
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Beenakker, DiVincenzo, Emary, and Kindermann in 2004 showed that allowing for charge detection in noninteracting fermions gives full quantum computation. Constructed a CNOT using beamsplitters, charge detectors, and one ancilla.

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