

Hiding Single Photons with Spread Spectrum Technology

Chinmay Belthangady,^{1,*} Chih-Sung Chuu,^{1,†} Ite A. Yu,² G. Y. Yin,¹ J. M. Kahn,¹ and S. E. Harris¹

¹Edward L. Ginzton Laboratory, Stanford University, Stanford, California 94305, USA

²Department of Physics, National Tsing Hua University, Hsinchu 300, Taiwan

(Received 1 April 2010; published 1 June 2010)

We describe a proof-of-principal experiment demonstrating the use of spread spectrum technology at the single photon level. We show how single photons with a prescribed temporal shape, in the presence of interfering noise, may be hidden and recovered.

DOI: 10.1103/PhysRevLett.104.223601

PACS numbers: 42.50.Gy, 32.80.Qk, 42.50.Ex, 42.65.Lm

Over the past several years scientists have learned how to extend techniques that are normally associated with classical coherent pulses of light to the single photon level. This has been motivated, in part, by the possibility of using such techniques for quantum communication and key distribution [1,2]. As an example, it is now possible to generate single photons that have a prescribed temporal waveform [3–5], and to use high-speed electro-optic modulators to arbitrarily modulate the amplitude of single photon [6] and biphoton [7] wave packets. Recently, by establishing a time origin with a laser pulse pumping an atom trapped in a cavity, Specht *et al.* [8] have demonstrated phase shaping and quantum interference of single photons. At the same time it has become clear how a single photon, based on its temporal waveform or time of arrival, may carry quantum information [9], and even, an entire image [10].

It is the intent of this Letter to further extend this classical-quantum interplay by demonstrating the use of spread spectrum technology [11] at the single photon level. We show how single photons with a prescribed temporal waveform may be transmitted through a noisy environment that is created by either narrow band thermal photons, or as in this work, an interfering laser beam that has an average power that is about a thousand times larger than the average power of the beam of single photons. We do this by using two synchronously driven electro-optic phase modulators. As shown in Fig. 1(a), the modulator at the transmitter M_1 broadens the spectrum of the incident photon beam from about 1 MHz to about 10 GHz and thereby reduces the spectral power density by a factor of 10^4 . The receiver modulator M_2 is run in antiphase to the transmitter modulator so as to demodulate the photon beam and reduce its bandwidth to the original 1 MHz, thereby allowing it to be transmitted through a narrow bandpass filter F_2 .

Now suppose that one wishes to transmit the single photon beam through an environment of thermal photons that have a linewidth that is comparable to that of the original photon beam. If this noise, or instead, a narrow-band interfering laser beam is injected after the first modulator [Fig. 1(a)], it is spectrally broadened by the second modulator and only a fraction of its power is transmitted

through the narrow band filter at the receiver. As a result, the signal-to-noise is increased approximately in the ratio of the modulation bandwidth to the bandwidth of the final filter, and in the language of spread spectrum technology the system experiences a “processing gain.” In this work we demonstrate a processing gain of a factor of 50.

The spread spectrum technique requires that the bandwidth of the modulated photons be much larger than the bandwidth of the unmodulated photons. Here, we use narrow-band time-energy entangled pairs of photons (biphotons) produced in a high optical density two dimensional magneto-optic trap (MOT) composed of laser cooled ^{87}Rb atoms. On application of counterpropagating, continuous-wave pump, ω_p and coupling, ω_c lasers to

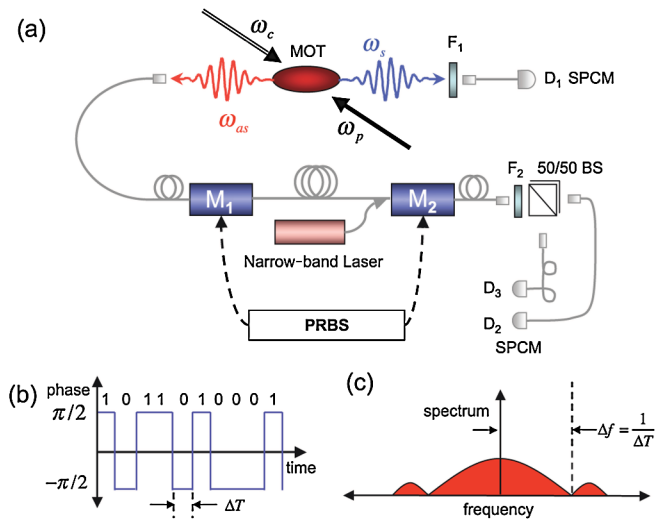


FIG. 1 (color online). Schematic of paired photon generation and spread spectrum experiment. (a) The spectrum of anti-Stokes photons produced in a magneto-optic trap is spread and narrowed by phase modulators M_1 and M_2 driven by a pseudorandom bit sequence (PRBS) generator. A weak narrow-band laser is used to hide the photons. F_1 and F_2 are narrow-band filters. A 50/50 beam splitter (BS) is used to verify the presence of single photons. (b) Pseudorandom phase modulation in the time domain. The labels 1, 0 refer to the outputs of the PRBS generator. (c) Schematic of the optical spectrum after phase modulation by the PRBS generator.

the atom cloud, pairs of counterpropagating Stokes, ω_s and anti-Stokes, ω_{as} photons are spontaneously generated in the medium and collected into single-mode fibers. Because of electromagnetically induced transparency (EIT), the coupling laser renders the atom cloud transparent at the anti-Stokes frequency. As a result of the slow light effect associated with the narrow EIT window, the anti-Stokes photon propagates with a group velocity that is about 10^4 times slower than the Stokes photon, and the temporal length of the biphoton wave packet (250 ns) is approximately equal to the dimension of the Rb cloud (1.7 cm) divided by group velocity of the anti-Stokes photon [12]. The physics of the paired photon generation process has been discussed by Balic *et al.* [4], and Kolchin [13], and operation with a subnatural linewidth has been reported by Du *et al.* [14].

We set the optical depth of the MOT at 30. With $2\gamma_{13}$ equal to the spontaneous decay rate out of the $|5P_{1/2}, F=2\rangle$ state of ^{87}Rb , we set the coupling laser Rabi frequency $\Omega_c = 4.01\gamma_{13}$, the pump laser Rabi frequency $\Omega_p = 1.66\gamma_{13}$ and the pump laser detuning, $\Delta\omega_p = 48.67\gamma_{13}$. Under these conditions the source produces narrow-band biphotons with an estimated linewidth of 3.5 MHz at a rate of about 11500 pairs/s.

The detection of the Stokes photon by detector D_1 sets the time origin for measurement of the biphoton wave function at detector D_2 . The anti-Stokes photon is sent through a 20-GHz electro-optic phase modulator (Eospace, Inc.) driven by one of two differential data outputs of a pseudorandom bit sequence (PRBS) generator (Centellax TG2P1A) with a bit rate tunable between 0.05–10 Gb/s. In each time slot ΔT , equal to the inverse bit rate, the PRBS outputs a random voltage equal to $V_\pi/2$ (corresponding to logic 1) or $-V_\pi/2$ (logic 0), where V_π is the half-wave voltage of the phase modulators. We choose a bit pattern consisting of $32\,767 (2^{15} - 1)$ random bits. At a bit rate of 10 Gb/s, the phase therefore switches randomly [Fig. 1(b)] between $-\pi/2$ and $\pi/2$ radians every 0.1 ns. This process of random phase modulation in the time domain spreads the spectrum in the frequency domain from about 3.5 MHz for the unmodulated photon to about 10 GHz for the modulated photon. The envelope of the spread spectrum is a sinc-squared function with the first null, as shown in Fig. 1(c), at a frequency equal to the bit rate.

The phase-modulated anti-Stokes photon is transmitted over a channel consisting of a long single-mode optical fiber [Fig. 1(a)]. At the receiver end, the photon is sent through a second phase modulator, identical to the first and driven by the complementary output of the PRBS generator. The two phase modulators are synchronized by adjusting the lengths of the electric cables connecting them to the PRBS generator such that the relative time delay, Δt_{RF} of the radio frequency signal sent to the two modulators equals the travel time, Δt_p , of the anti-Stokes photon

between the modulators. Fine adjustment of Δt_{RF} is achieved by means of a radio frequency waveguide of variable length connected to one of the two rf cables. When the phase modulators are run in synchronism such that the phase imposed on the photon by the first modulator is exactly undone by the second modulator, the spectral width of the photon is restored to its original value of 3.5 MHz. When the phase modulators are not run in exact phase opposition, the uncompensated residual phase results in a spectrum that is wider than 3.5 MHz. After passing through the second phase modulator the anti-Stokes photon is sent through a 65-MHz fiber based Fabry-Perot filter (Micron Optics) with a free spectral range of 13.6 GHz. If the spectral width of the photon after the second phase modulator is less than the filter bandwidth, the photon passes through with little loss of probability amplitude. Conversely, if the spectral width of the photon is larger, then it is attenuated by the filter. The spectrum of narrow-band noise entering the channel after the first modulator is spread by the second modulator and is attenuated by the filter. To hide the photon we use an attenuated beam from an external-cavity diode laser with a linewidth of about 300 kHz set at the same frequency as the anti-Stokes photon. A 50/50 beam splitter (BS) and a third detector D_3 is placed after the filter to verify the single photon nature of the experiment.

With both phase modulators and the (noise simulating) laser turned off, the shape of the anti-Stokes photon wave packet is shown in Fig. 2(a). This shape is measured by recording coincidence counts between detectors D_1 and

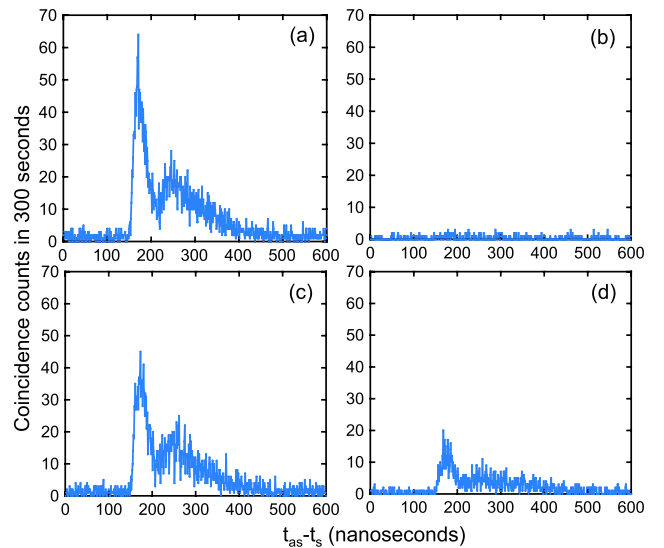


FIG. 2 (color online). Shape of the anti-Stokes photon wave packet measured by recording coincidence counts between detectors D_1 and D_2 with the beam splitter removed. (a) Anti-Stokes photon wave packet measured when modulators M_1 and M_2 are off. (b) Modulator M_1 is on and modulator M_2 is off. (c) Both modulators are on and run synchronously. (d) Both modulators are on with $\Delta t_{\text{RF}} > \Delta t_p$ (see text).

D_2 , with the beam splitter removed, as a function of the difference in arrival times of the anti-Stokes, t_{as} , and Stokes, t_s , photons at the two detectors. Coincidence measurement is done by means of a time-to-digital converter (FAST Comtec TDC 7886s) with counts binned into 1-ns bins. The waveform consists of a sharp leading edge spike called a Sommerfield-Brillouin precursor [15], and a main waveform whose width is approximately equal to the group delay at the anti-Stokes frequency [14]. Conditioned on the detection of the Stokes photon, and neglecting the flat uncorrelated background, this shape represents the squared modulus of the wave function of the anti-Stokes photon. When the first phase modulator is switched on, the spectrum of this photon is spread to about 10 GHz. If the second phase modulator is off, the anti-Stokes photon is severely attenuated by the narrow-band filter F_2 and very little photon probability amplitude leaks through to detector D_2 . This is shown in Fig. 2(b). When the second phase modulator is switched on and driven so as to cancel the phase imposed by the first modulator, the photon spectrum is narrowed and passes through filter F_2 with little loss of probability amplitude. This is shown in Fig. 2(c). The small residual loss in photon probability amplitude probably arises from a slight mismatch between the two modulators and their complementary driving signals. In Fig. 2(d), the difference in lengths of the electric cables connecting the PRBS to the two modulators is chosen such that Δt_{RF} is greater than Δt_p by about 40 ps. As a result, the phase imposed on the photon by the first modulator is not exactly negated by the second modulator. Now, after transmission through filter F_2 we observe a reduced, but not zero, coincidence count rate. A similar result is obtained when $\Delta t_{RF} < \Delta t_p$.

We next deliberately inject a weak diode laser beam to hide the anti-Stokes photon. With additional laser photons at a rate of $40\,000\text{ s}^{-1}$ and anti-Stokes signal photons at a rate of 30 s^{-1} [Fig. 3(a)], anti-Stokes photons can no longer be seen. If the two modulators are now turned on and run in synchronism, the spectrum of the anti-Stokes

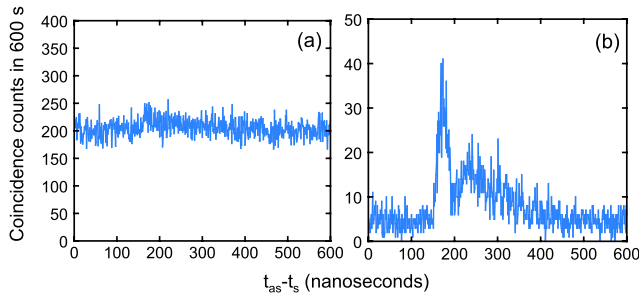


FIG. 3 (color online). Retrieval of the anti-Stokes photon from injected noise. (a) Anti-Stokes photon wave packet as hidden by an additional laser beam that is injected into the channel. Here, both modulators are off. (b) Switching on both modulators and running them synchronously recovers the anti-Stokes photon (see text).

photon is spread by the first modulator and compressed by the second modulator so that the photon passes through the filter and arrives at the detector with very little attenuation. The spectrum of the laser photons, on the other hand, is spread by the second phase modulator and attenuated by the filter and very little of the laser power reaches the detector. When this is the case, the shape of the anti-Stokes photon wave packet can be recovered from the noise [Fig. 3(b)]. The decrease in the level of the flat uncorrelated background by a factor of about 50 corresponds to the spread spectrum processing gain.

We next test to insure that the experiment results are dominated by single photon events. To the extent that the down conversion process produces only biphotons, and when the square of the rate of biphoton generation multiplied by the gate width is much less than unity, then no threefold coincidences will be observed between the detectors D_1 , D_2 , and D_3 . The quality of a single photon source can be quantified by means of the conditional Glauber correlation function [6] defined as

$$g_{\text{cond}}^{(2)}(0) = \frac{N_{123}N_1}{N_{12}N_{13}}, \quad (1)$$

where N_{123} is the number of threefold coincidences between detectors D_1 , D_2 , and D_3 , N_{12} , N_{13} are the number of twofold coincidences between detectors D_1 , D_2 and D_1 , D_3 , respectively, and N_1 is the number of counts detected by detector D_1 . Here all coincidence counts are measured within a time window equal to the nominal length of the anti-Stokes photon wave packet. We turn off the weak narrow-band laser beam and measure $g_{\text{cond}}^{(2)}(0)$ as a function of the rate of detected Stokes photons with both modulators turned off [black line (blue line) in Fig. 4] and with both modulators turned on and running synchronously [gray line (red line) in Fig. 4]. At a Stokes rate of approximately 1200 s^{-1} we obtain $g_{\text{cond}}^{(2)}(0) = 0.22 \pm 0.06$ for the unmodulated photon and $g_{\text{cond}}^{(2)}(0) = 0.14 \pm 0.04$ for the

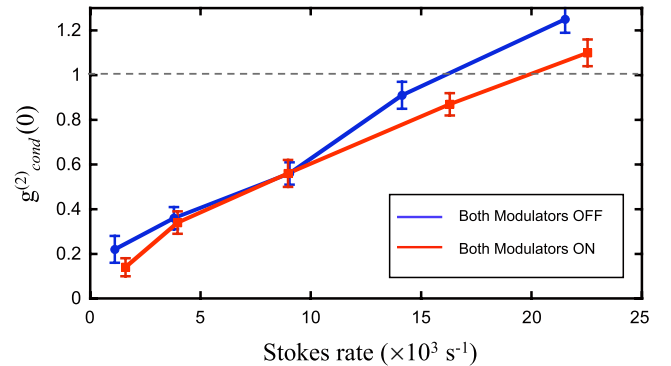


FIG. 4 (color online). Conditional threefold correlation function as a function of Stokes rate with both modulators off [black (blue)] and with both modulators on [gray (red)]. The dashed line is the classical limit. The errors are calculated by assuming a Poisson distribution for the coincidence rates.

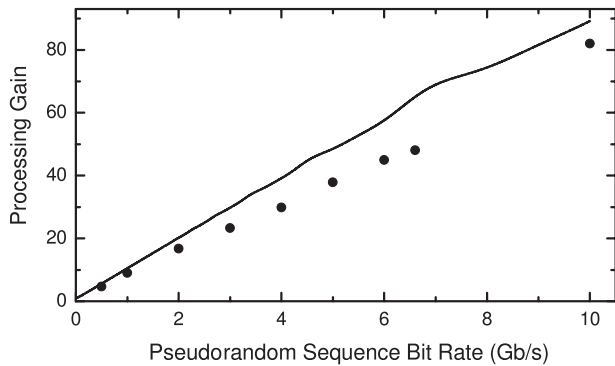


FIG. 5. Spread spectrum processing gain as a function of the PRBS bit rate. The measurement is made with a cw diode laser. Dots represent experimental data and the solid curve is calculated (see text).

phase-modulated-demodulated photon. These values of $g_{\text{cond}}^{(2)}(0)$ are less than the classical limit of 1 (dashed line in Fig. 4) and reasonably less than the limiting value of 0.5, i.e., the conditional Glauber correlation function for a two-photon Fock state.

As an additional check on the modulators that are used in this work, as shown in Fig. 5, we measure the transmission of the final filter versus the bit rate of the PRBS source. We send a laser beam with a linewidth of about 300 kHz through phase modulator M_2 and narrow-band filter F_2 . The ratio of the laser power detected after the filter with the modulator off to the power measured with the modulator on gives the value of what is termed as the spread spectrum processing gain. The dots are experimental data and the solid line is calculated by evaluating the overlap integral between the transfer function of the Fabry-Perot filter and the sinc-squared spectrum produced by pseudorandom phase modulation. The fact that the experimental points lie below the calculated curve indicates that the spectrum of the laser beam is not spread to an ideal sinc-squared spectrum with the first null at a frequency equal to the bit rate. This deviation tends to be larger for higher bit rates, suggesting that it arises from the finite bandwidth of the modulators and the drive electronics.

This work has shown how narrow-band single photons may be phase modulated so as to increase the width of their spectrum by several orders of magnitude while at the same time retaining the information that characterizes their waveform. By compensation with an antiphased modula-

tor, the biphoton waveform may be reconstructed at a distant location. Though antiphased operation of the modulators may be a problem at large distances, intermittent transmission of synchronizing pulses could be used to overcome it. Since the spread spectrum technique leaves the time-frequency entanglement properties of the photons intact, dispersion in the anti-Stokes channel may be canceled nonlocally by dispersion in the Stokes channel [16].

We have shown how spread spectrum technology may be used to hide a single photon in the presence of laser photons of the same frequency and similar linewidth. Applications may include an additional level of classical security for quantum key distribution, and someday, multiplexing of single photon communication channels. This work may also be the first use of spread spectrum techniques combined with single photon detection.

This work was supported by the U.S. Air Force Office of Scientific Research, the U.S. Army Research Office, and the Defense Advanced Research Projects Agency.

*chinmayb@stanford.edu

†cschuu@stanford.edu

- [1] N. Gisin, G. Ribordy, W. Tittel, and H. Zbinden, *Rev. Mod. Phys.* **74**, 145 (2002).
- [2] K. Inoue, E. Waks, and Y. Yamamoto, *Phys. Rev. Lett.* **89**, 037902 (2002).
- [3] M. Keller *et al.*, *Nature (London)* **431**, 1075 (2004).
- [4] Vlatko Balic, Danielle A. Braje, Pavel Kolchin, G. Y. Yin, and S. E. Harris, *Phys. Rev. Lett.* **94**, 183601 (2005).
- [5] A. Kuhn, M. Hennrich, and G. Rempe, *Phys. Rev. Lett.* **89**, 067901 (2002).
- [6] P. Kolchin, C. Belthangady, S. Du, G. Y. Yin, and S. E. Harris, *Phys. Rev. Lett.* **101**, 103601 (2008).
- [7] C. Belthangady, S. Du, C.-S. Chu, G. Y. Yin, and S. E. Harris, *Phys. Rev. A* **80**, 031803(R) (2009).
- [8] H. P. Specht *et al.*, *Nat. Photon.* **3**, 469 (2009).
- [9] Y. Yamamoto and H. A. Haus, *Rev. Mod. Phys.* **58**, 1001 (1986).
- [10] C. J. Broadbent *et al.*, *Phys. Rev. A* **79**, 033802 (2009).
- [11] R. Scholtz, *IEEE Trans. Commun.* **25**, 748 (1977).
- [12] M. H. Rubin *et al.*, *Phys. Rev. A* **50**, 5122 (1994).
- [13] P. Kolchin, *Phys. Rev. A* **75**, 033814 (2007).
- [14] S. Du, P. Kolchin, C. Belthangady, G. Y. Yin, and S. E. Harris, *Phys. Rev. Lett.* **100**, 183603 (2008).
- [15] S. Du, C. Belthangady, P. Kolchin, G. Y. Yin, and S. E. Harris, *Opt. Lett.* **33**, 2149 (2008).
- [16] J. D. Franson, *Phys. Rev. A* **45**, 3126 (1992).